

チャームで見る エキゾチックハドロン・原子核

KEK

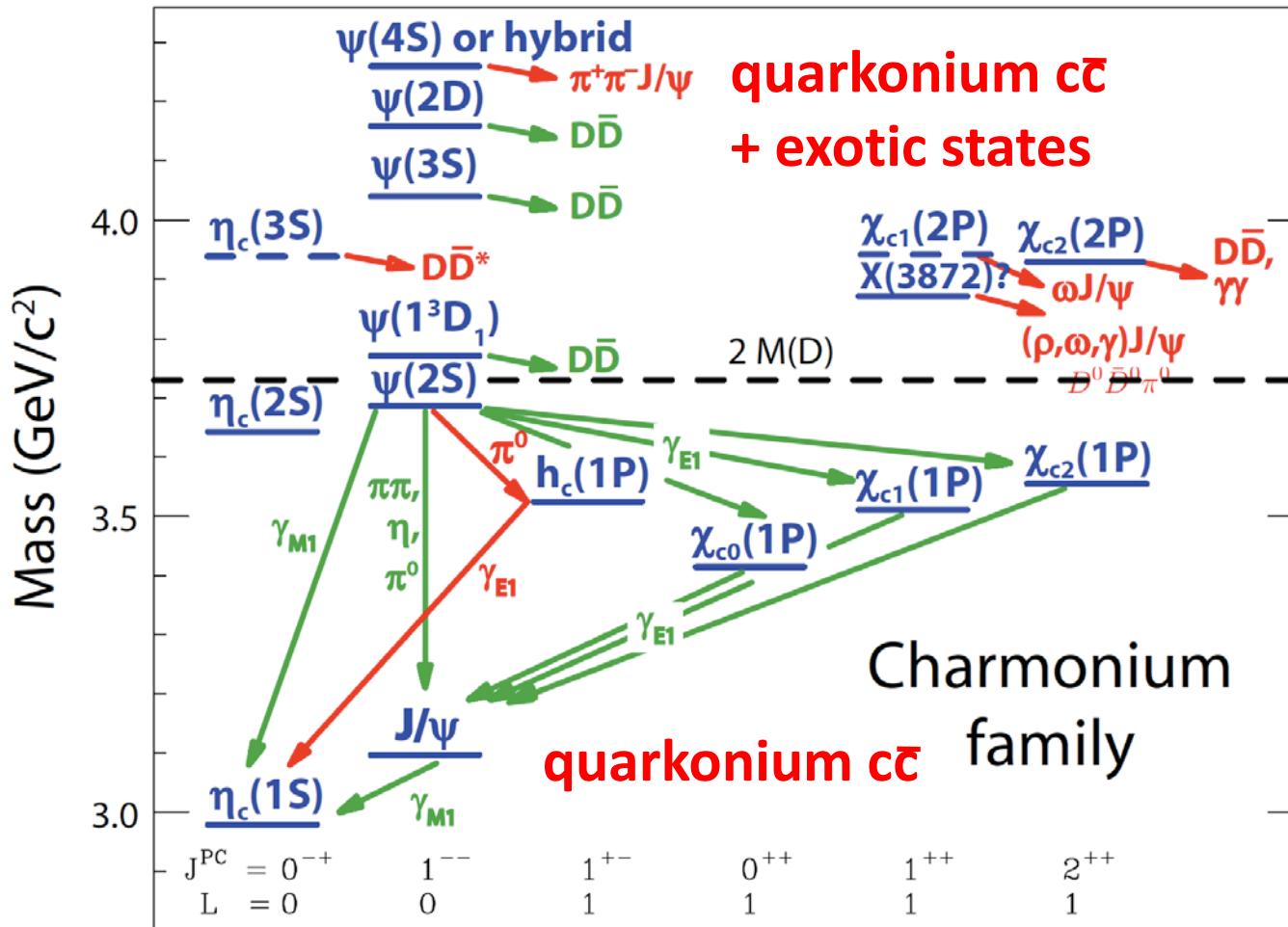
安井 繁宏

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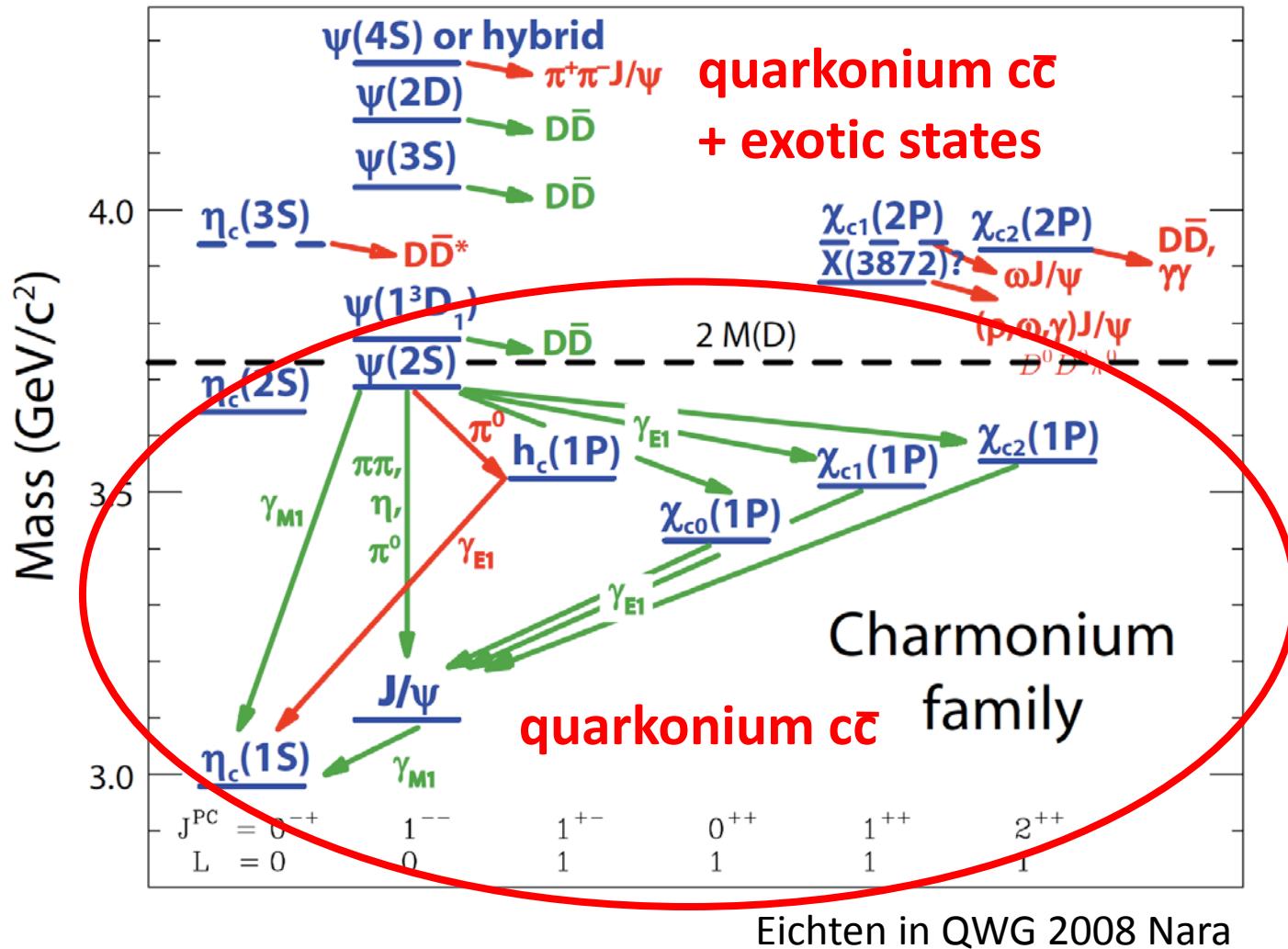
1. Introduction

Many new $c\bar{c}$ states above threshold



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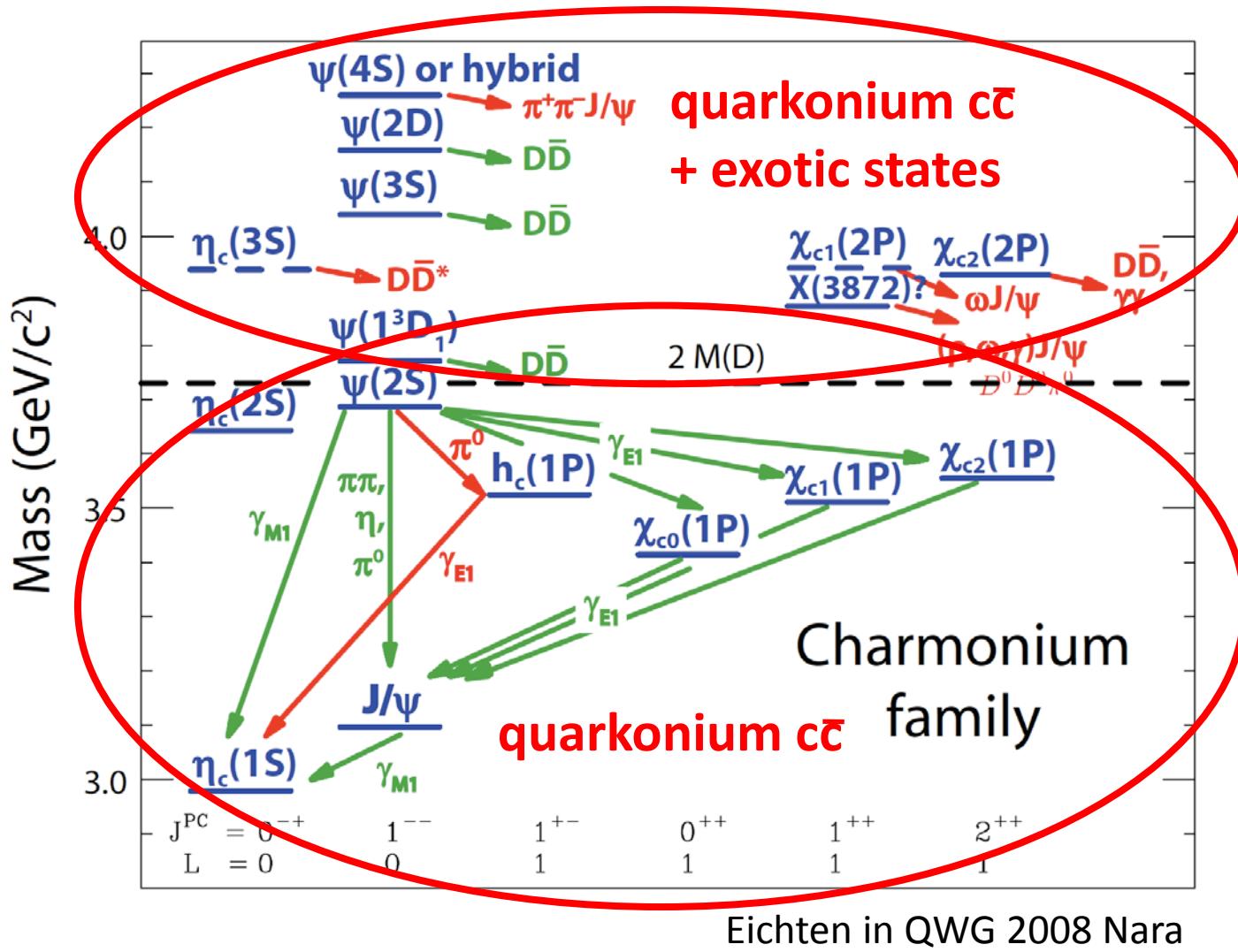
Many new $c\bar{c}$ states above threshold



Eichten in QWG 2008 Nara

1. Introduction

Many new $c\bar{c}$ states above threshold



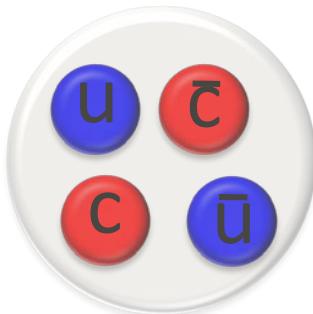
1. Introduction

$X(3872)$, $X(3940)$, $Y(4260)$, $Z(4430)^+$, ... Y_b , Z_b , ...

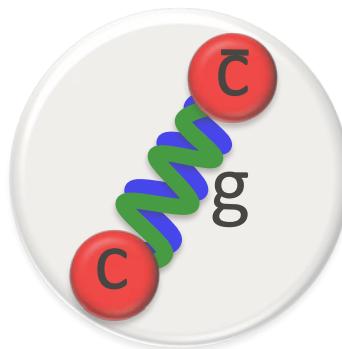
Why are they exotic?

- unusual decay widths and branching ratios
- mass difference from $c\bar{c}$ in quark model (~ 100 MeV)

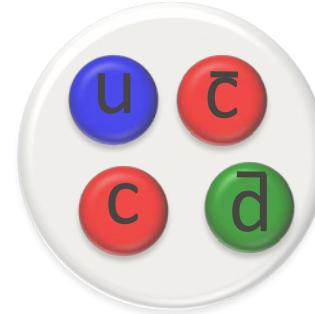
$X(3872)$



$Y(4260)$



$Z(4430)^+ (Z_b)$



T_{cc}^{-1}



$D\bar{D}^*$ ($c\bar{c}u\bar{u}$) molecule
Tornqvist (1991)



$D\bar{D}^*$ threshold ?

$c\bar{c}g$ hybrid
S. L. Zhu (2005)



gluon dynamics ?

$c\bar{c}u\bar{d}$ ($b\bar{b}u\bar{d}$)
really tetraquark



new correlation ? diquark correlation ?



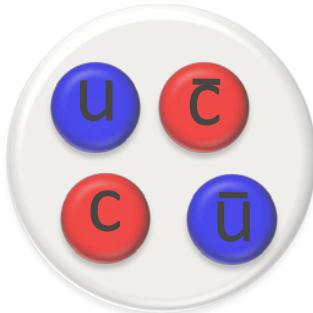
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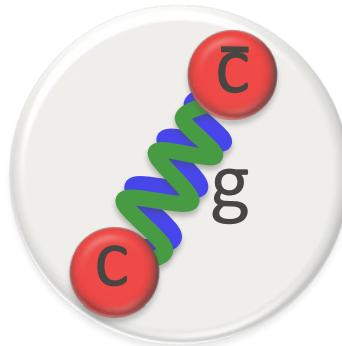
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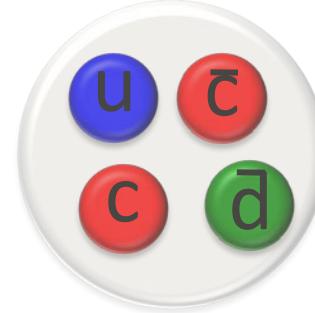
X(3872)



Y(4260)



Z(4430)⁺ (Z_b)



T_{cc}¹



D̄D^{*} (c̄cū) molecule
Tornqvist (1991)

c̄cg hybrid
S. L. Zhu (2005)

c̄cūd̄ (b̄b̄ud̄)
really tetraquark

c̄cud̄
really tetraquark

What states are allowed as solutions of QCD?

Can we construct effective theories from QCD?

2. Symmetry and dynamics of heavy quark hadrons



Chiral Symmetry

$SU(3)_L \times SU(3)_R$

π , K as Nambu-Goldstone bosons

Chiral multiplets (ρ - a_1 , N - N^* , ...)

Nucleus formed by tensor force

Hypernuclei as g.s. of QCD matter

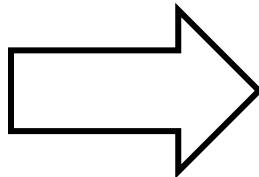
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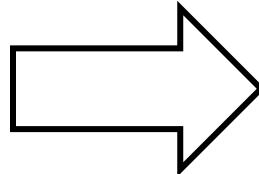
2. Symmetry and dynamics of heavy quark hadrons

3 5 150 200
u d s Λ_{QCD}

1500 $\Sigma\Sigma$ 4700 [MeV]
c b m

Chiral Symmetry

$SU(3)_L \times SU(3)_R$



Chiral Symmetry
 $SU(4)_L \times SU(4)_R$

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Chiral multiplets (ρ - a_1 , N - N^* , ...)

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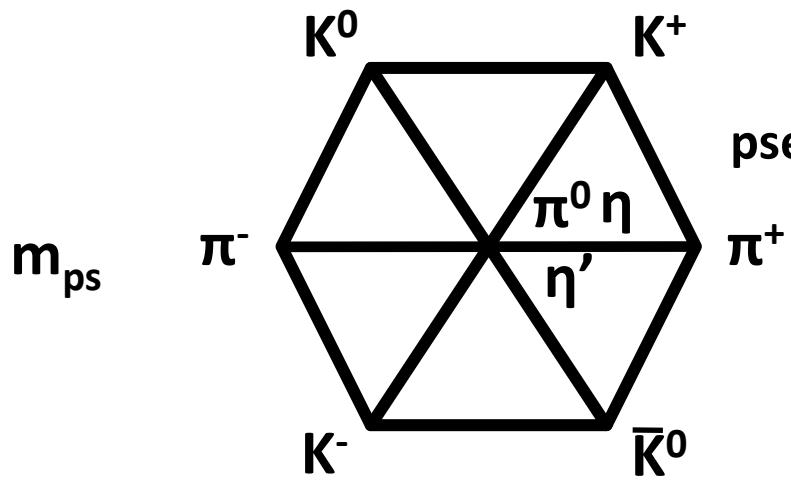
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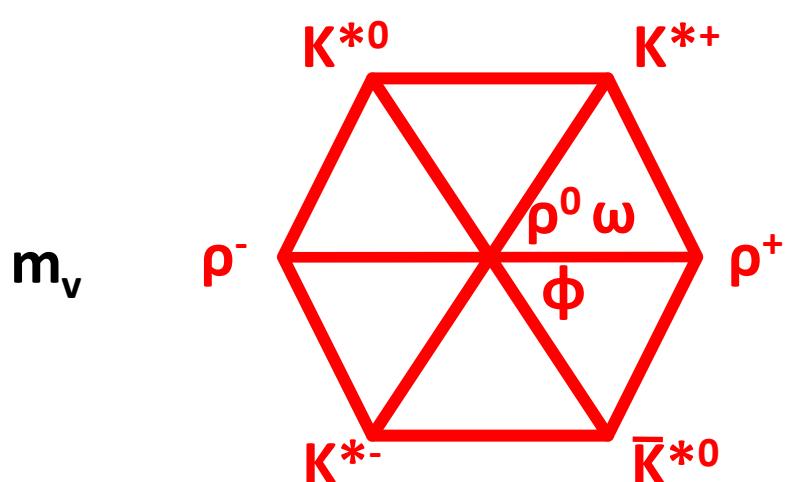
SU(3)_f symmetry



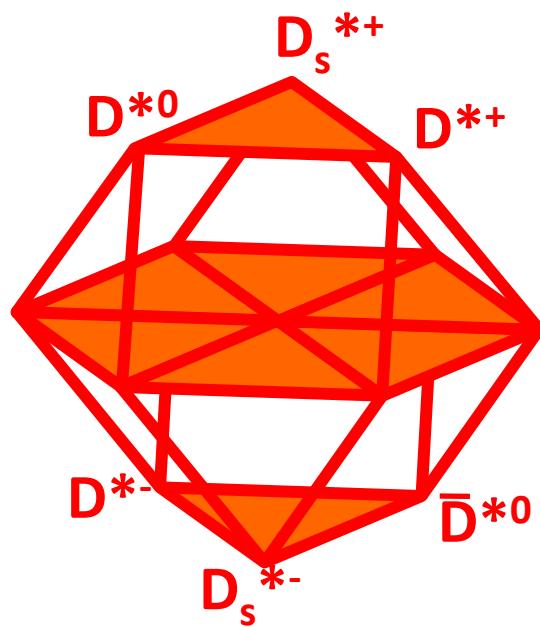
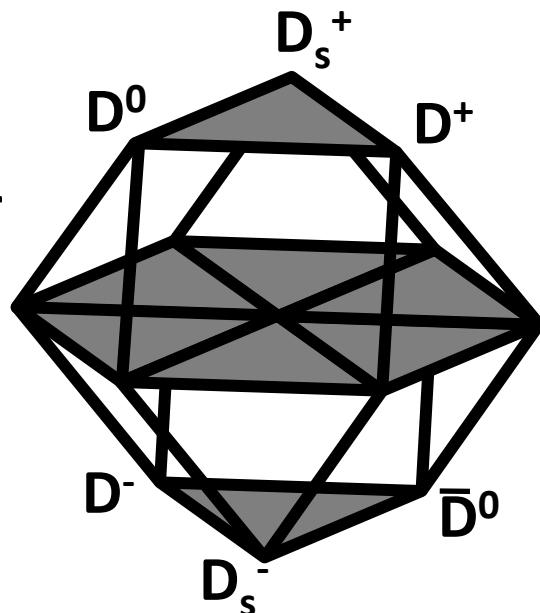
SU(4)_f symmetry ?



pseudo-scalar



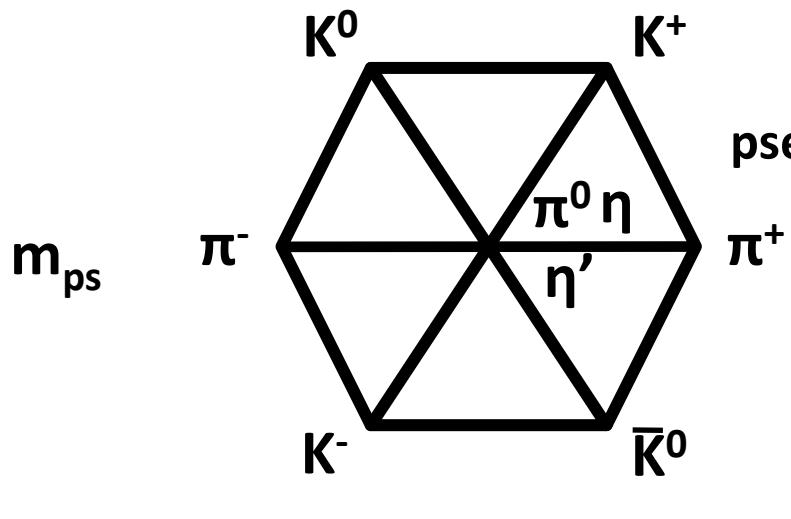
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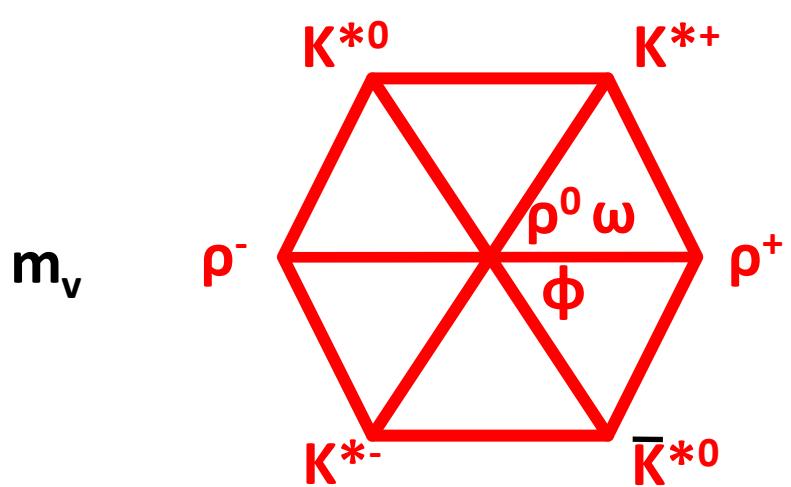
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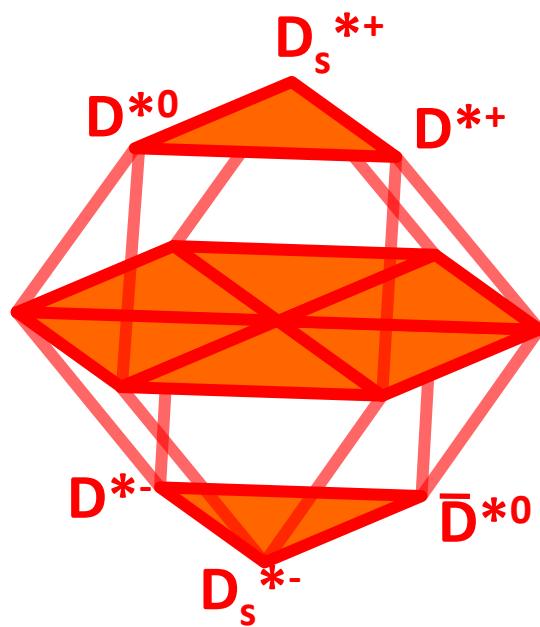
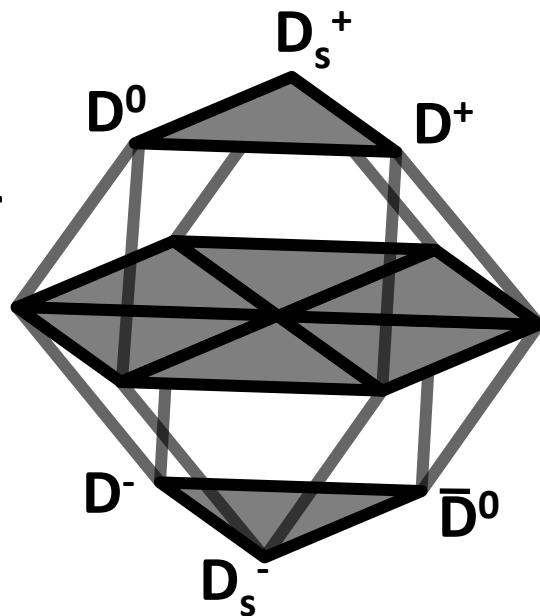
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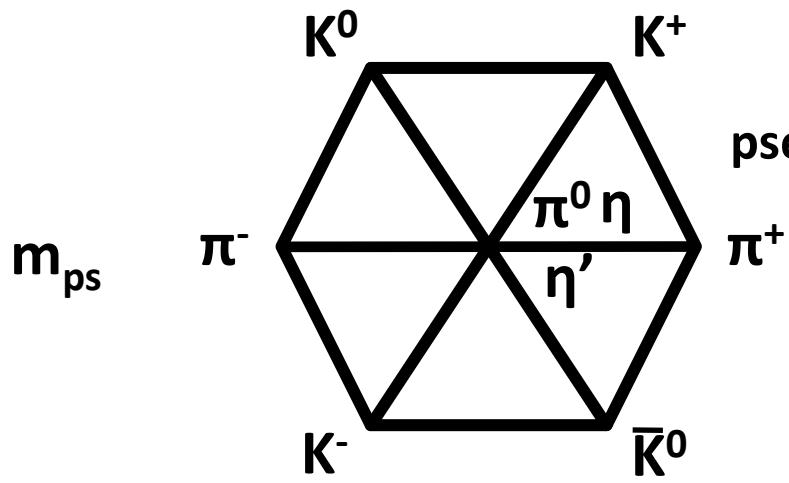
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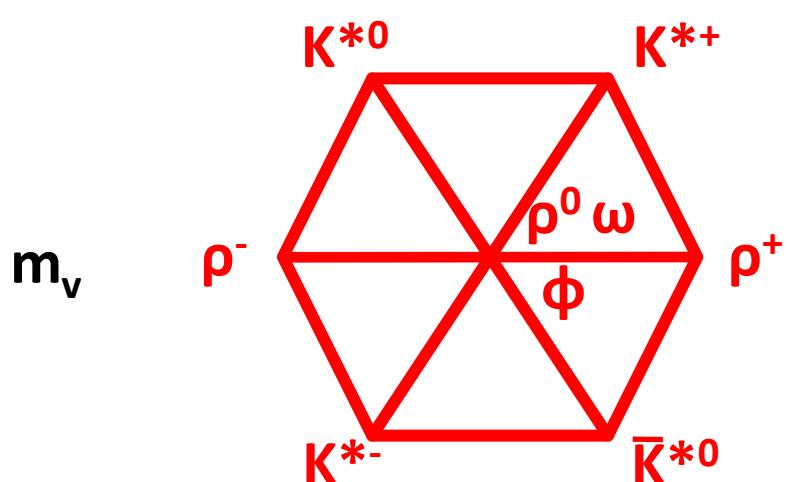
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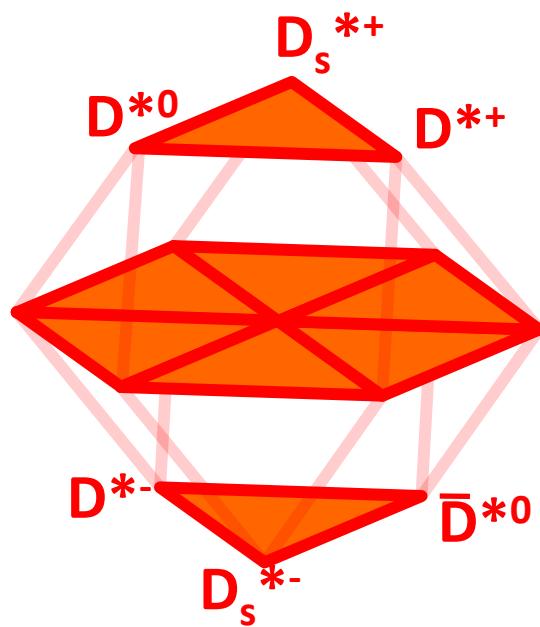
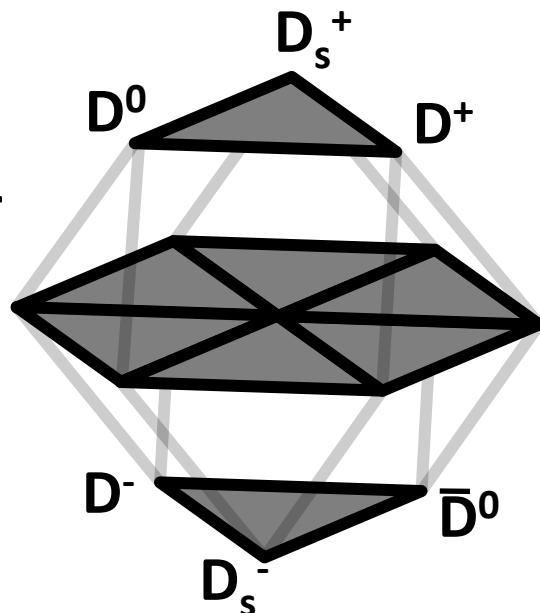
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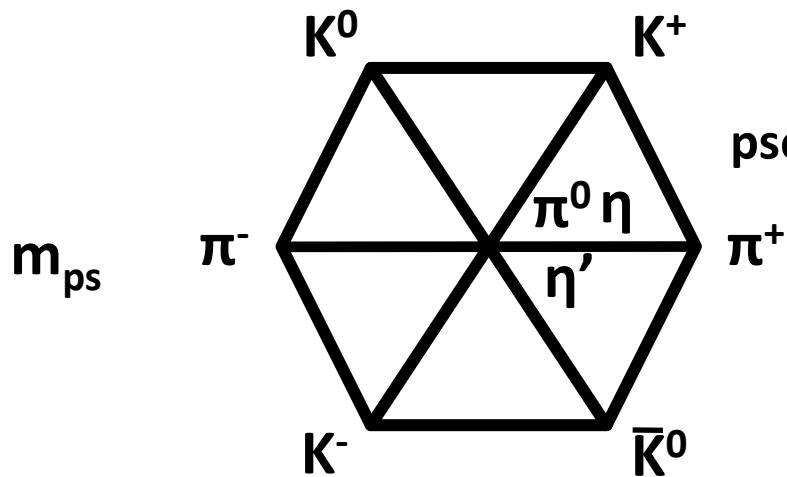
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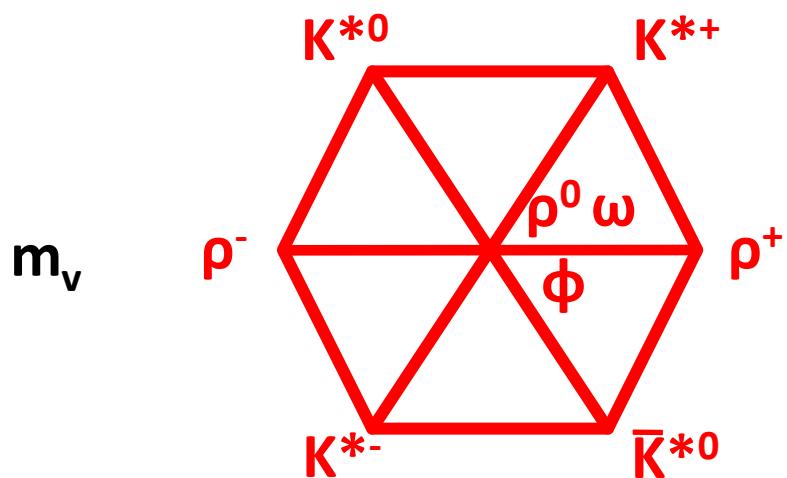
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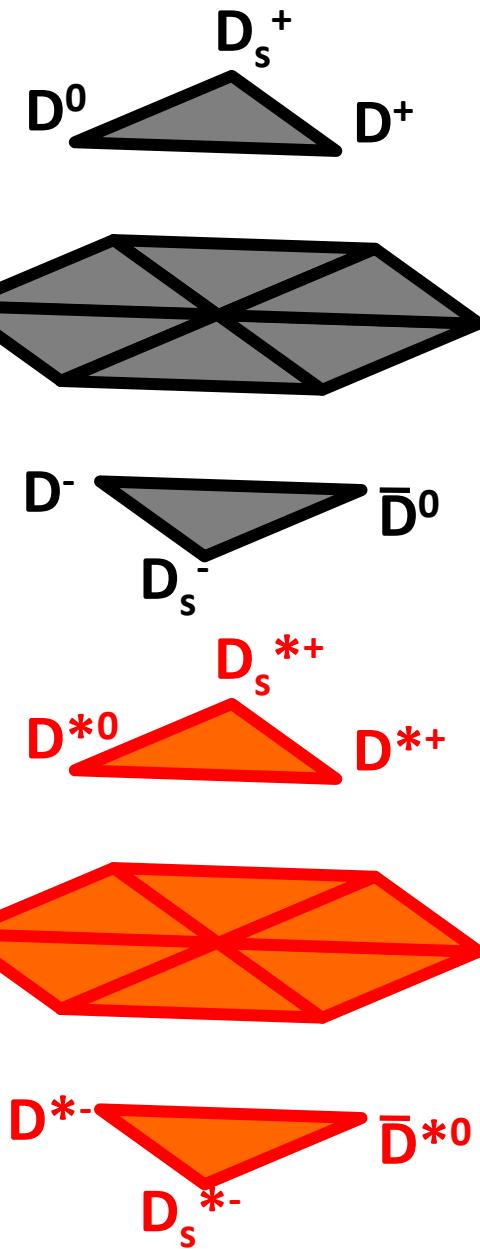
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pseudo-scalar



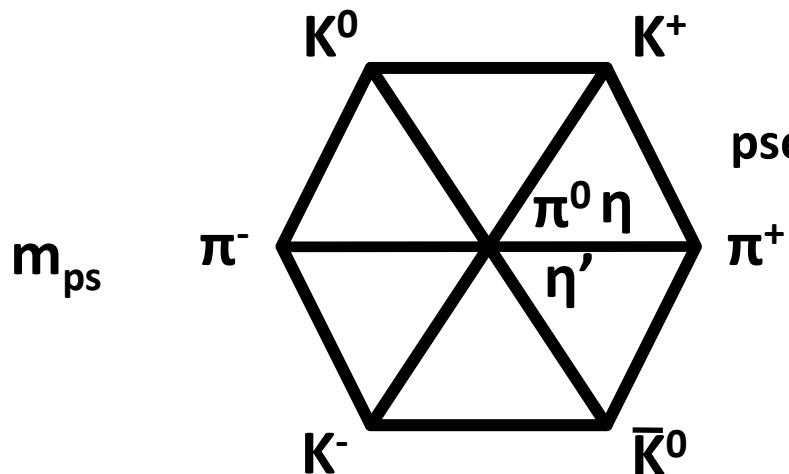
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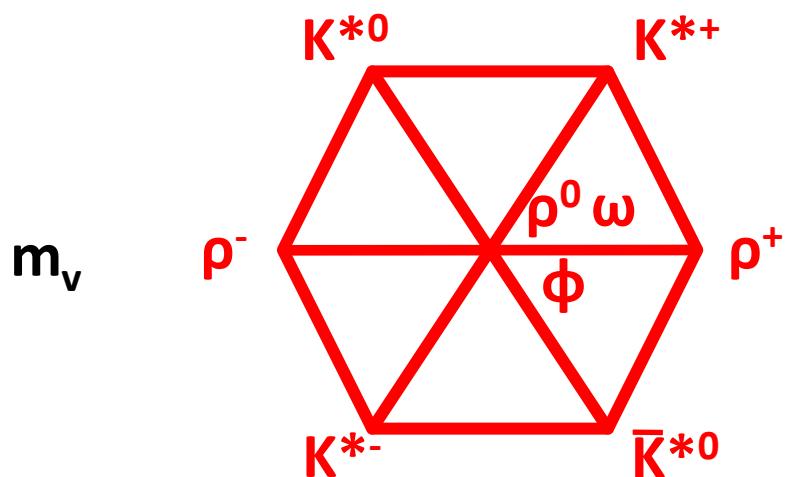
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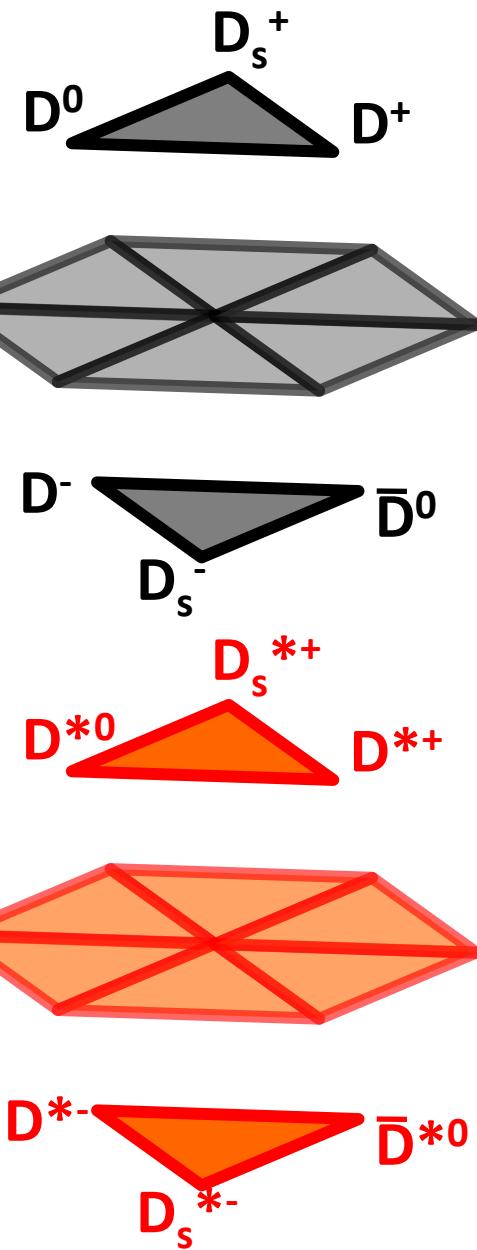
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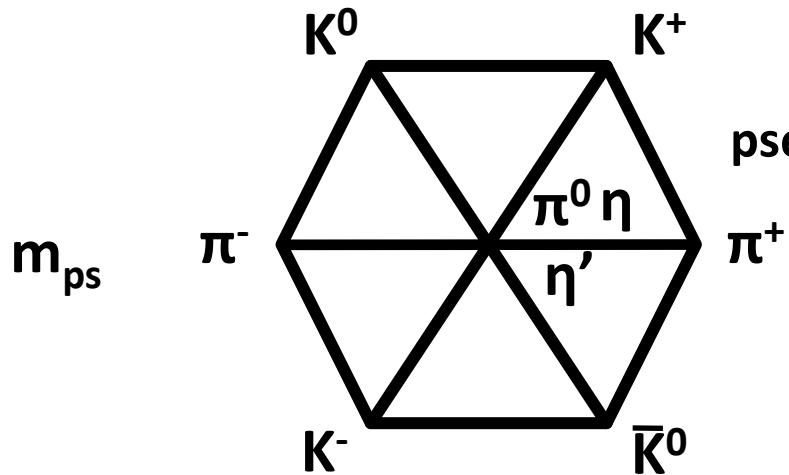
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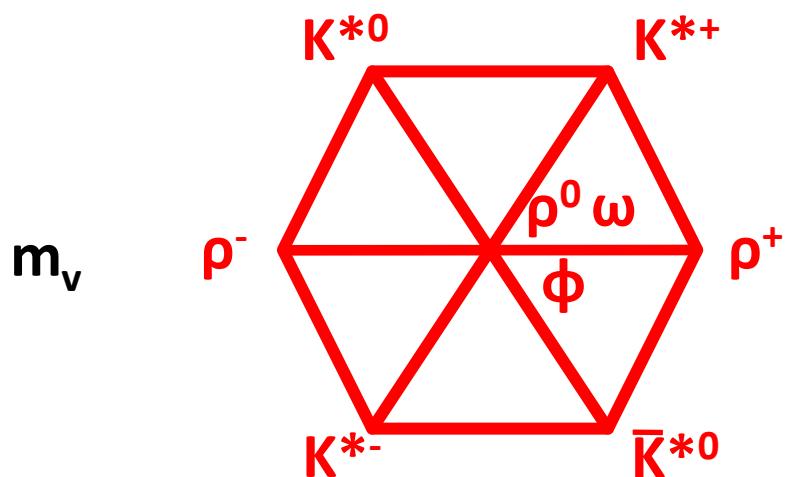
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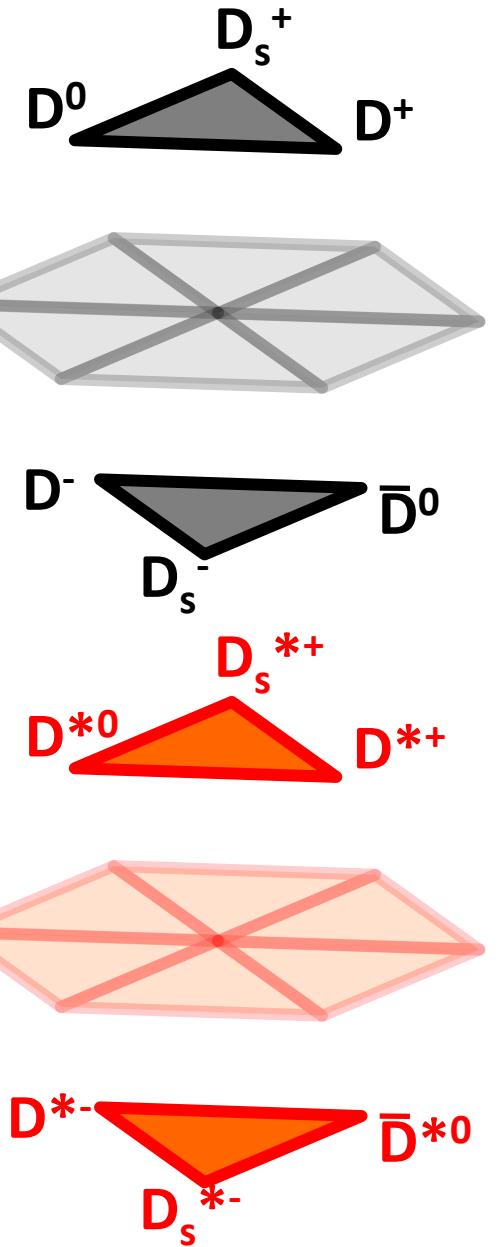
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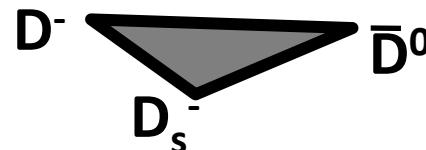
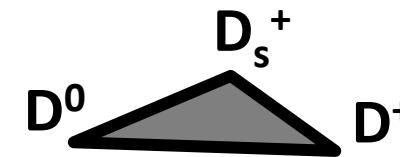
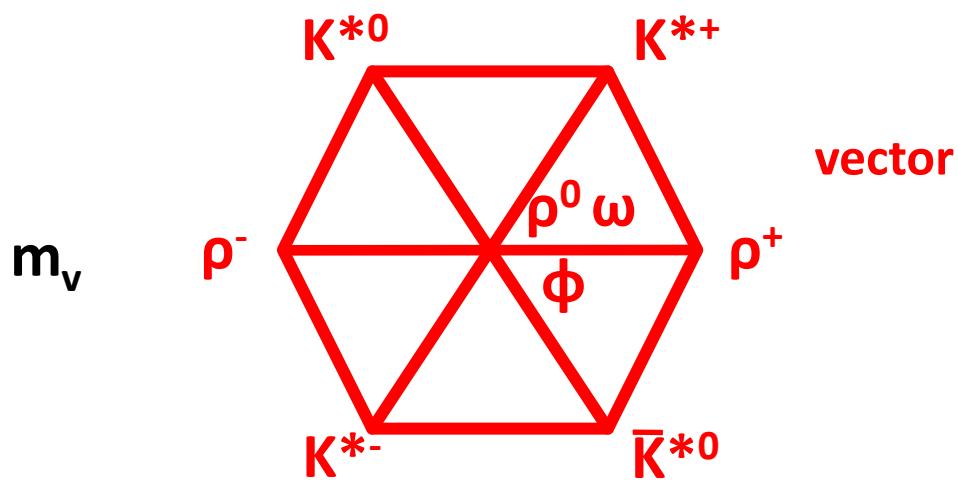
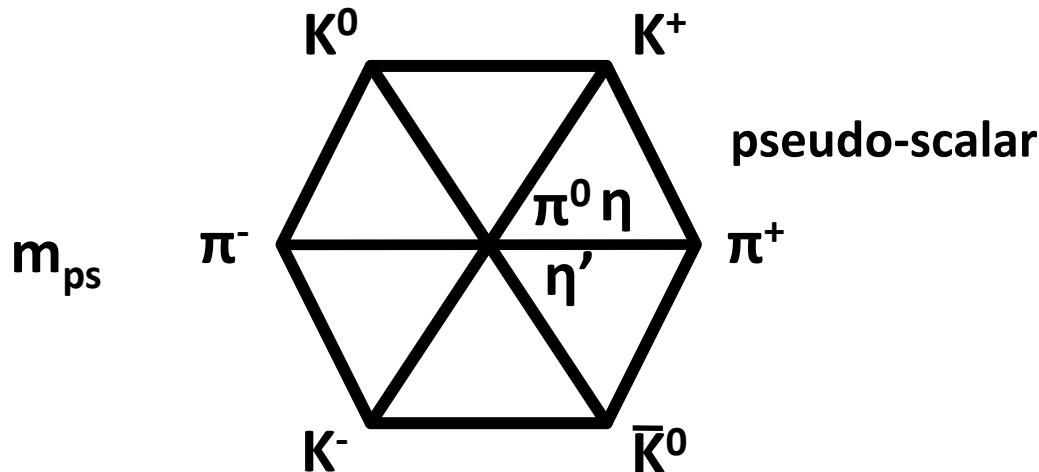
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SU(3)_f symmetry



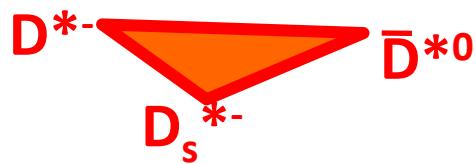
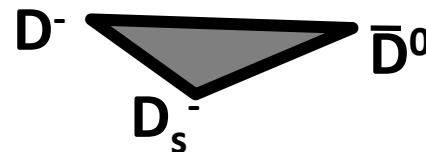
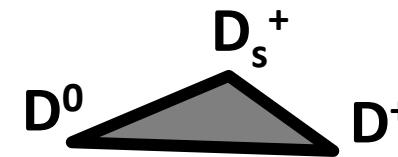
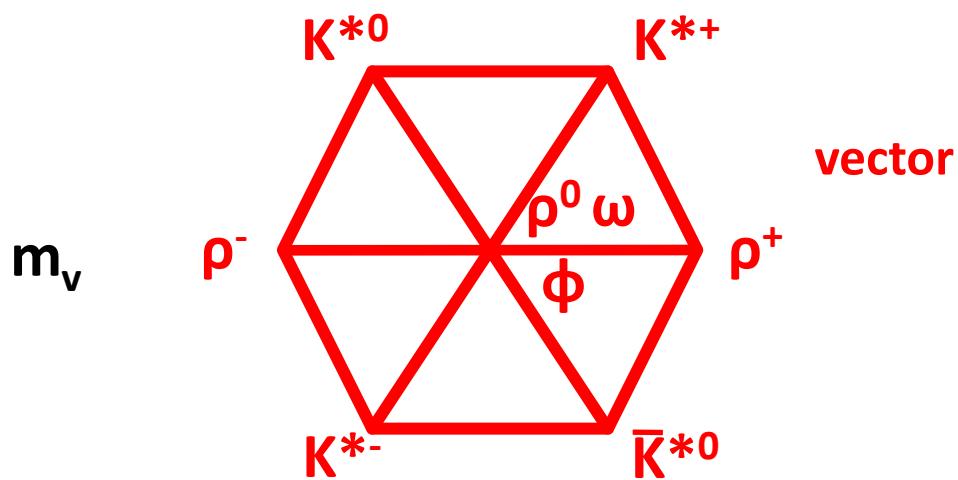
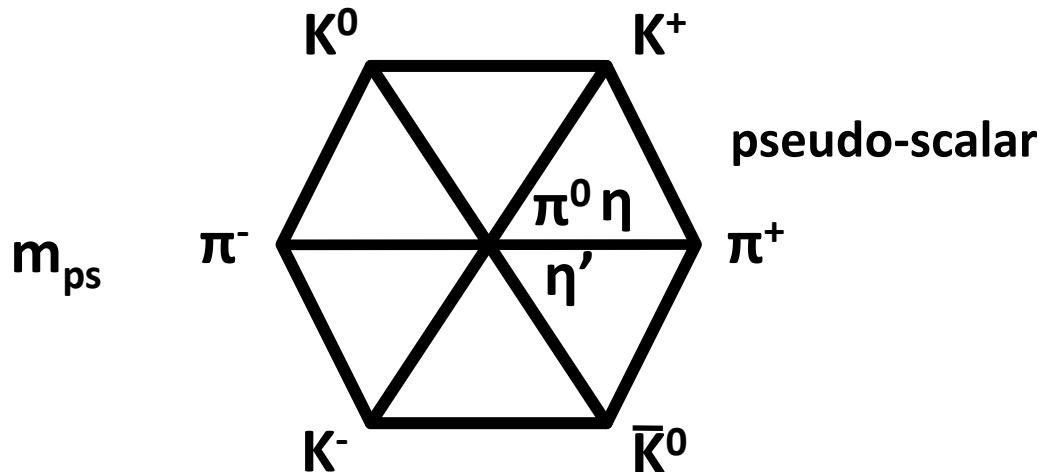
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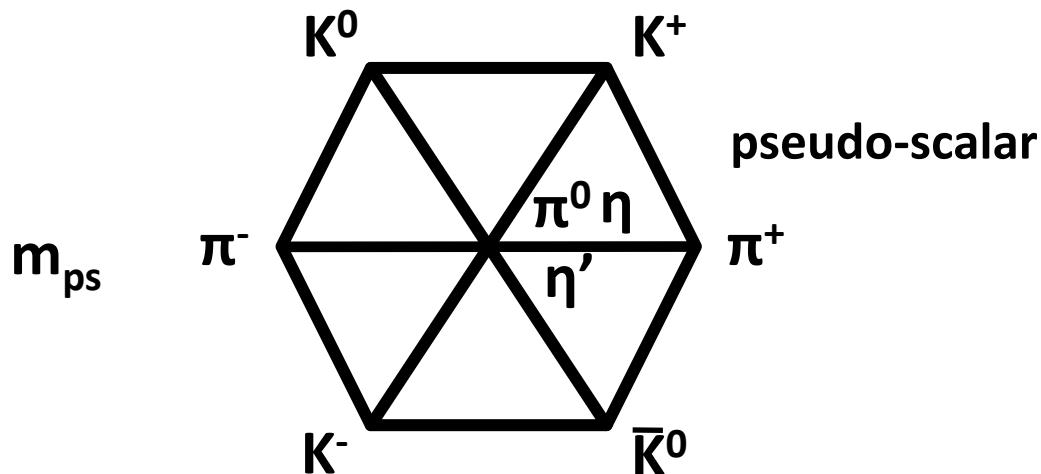
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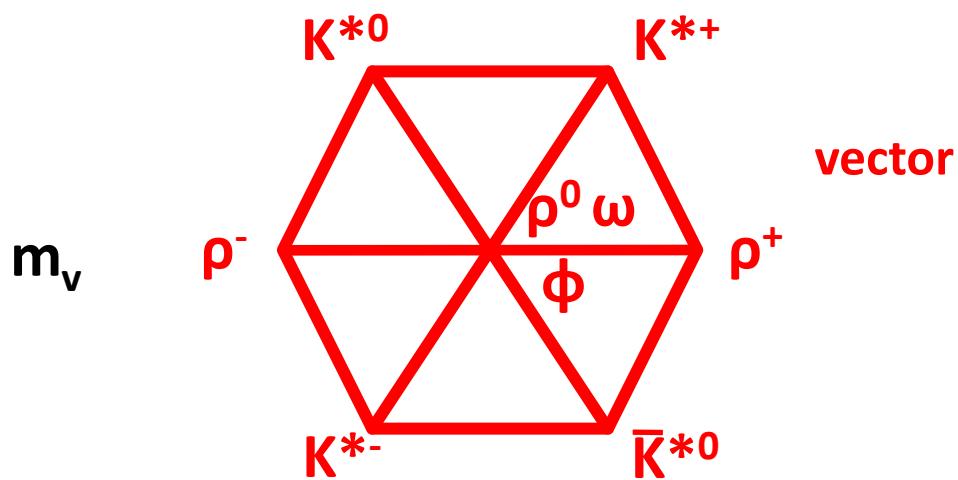
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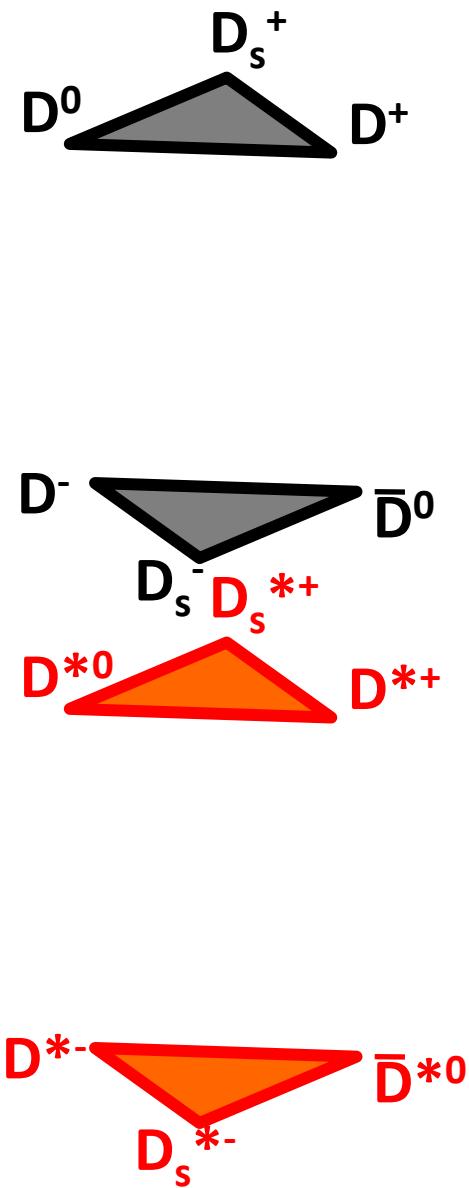
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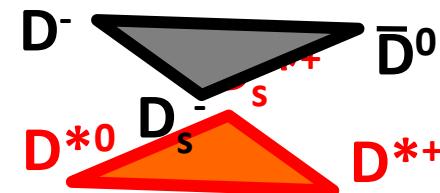
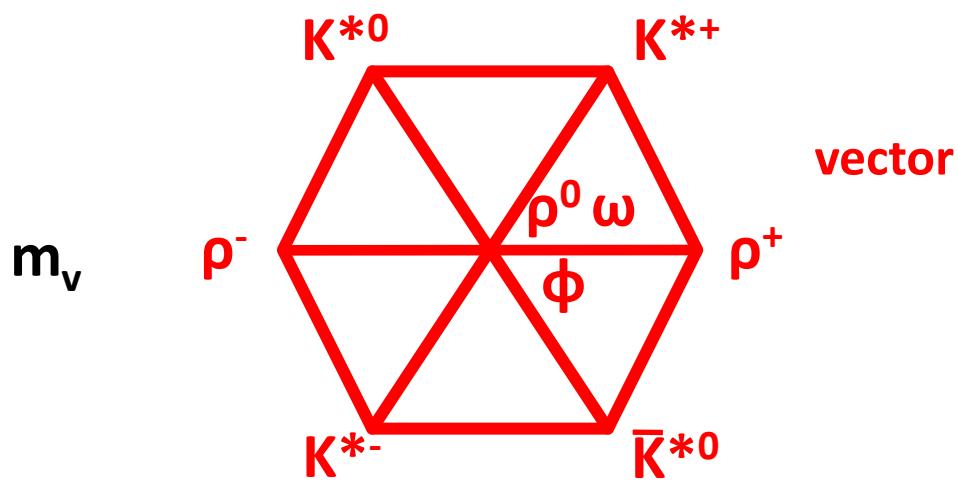
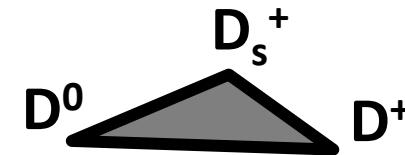
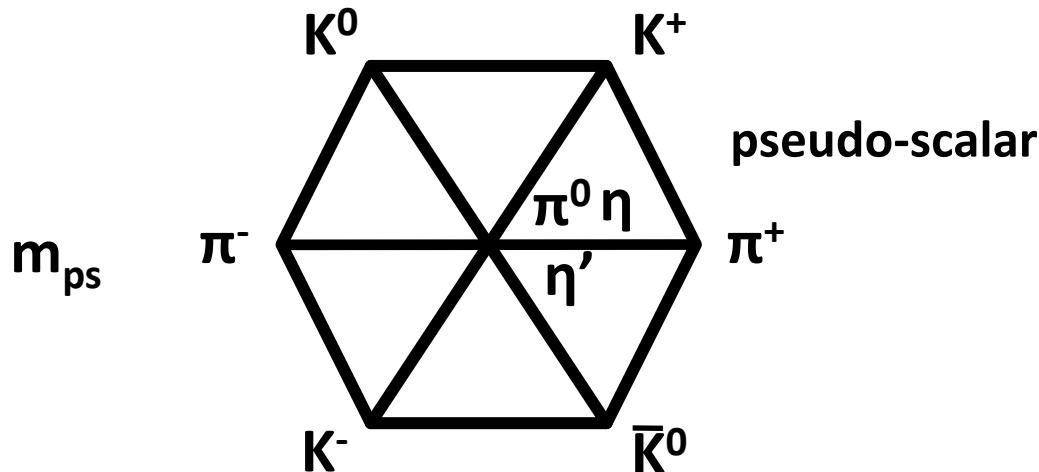
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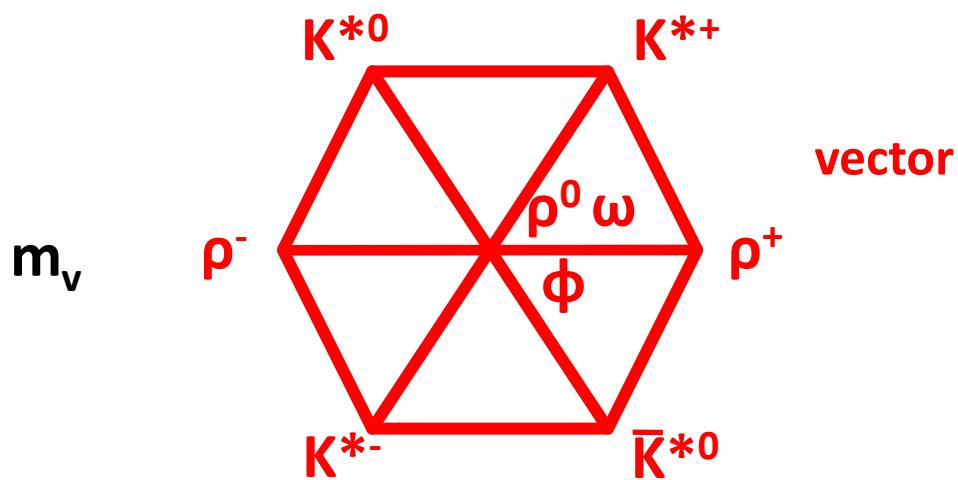
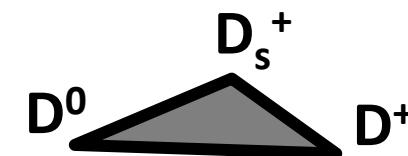
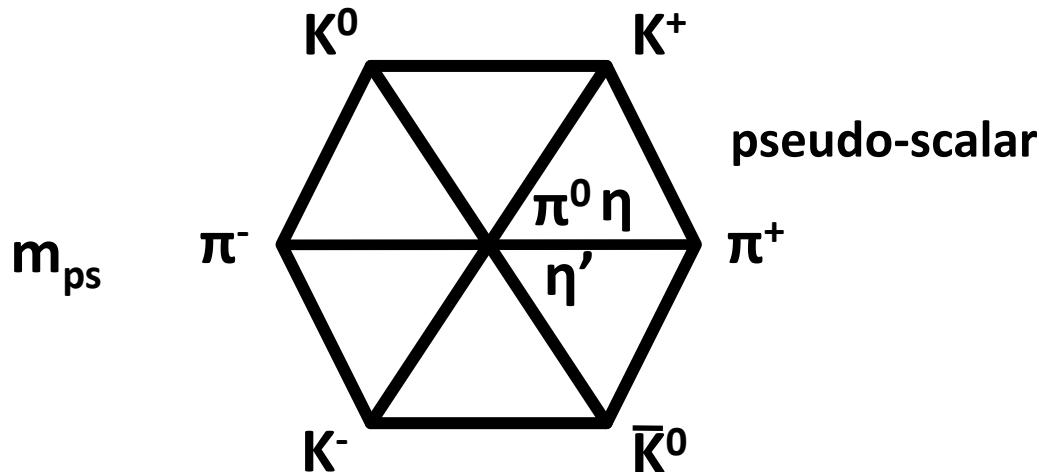
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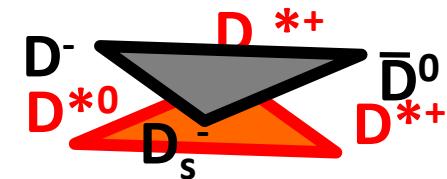
SU(4)_f symmetry ?



SU(3)_f symmetry → **SU(4)_f symmetry ?**



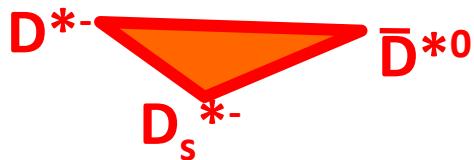
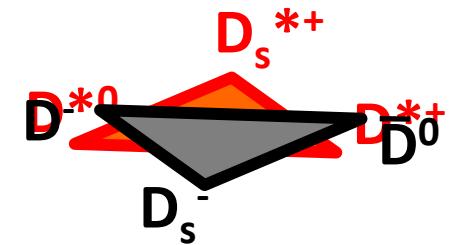
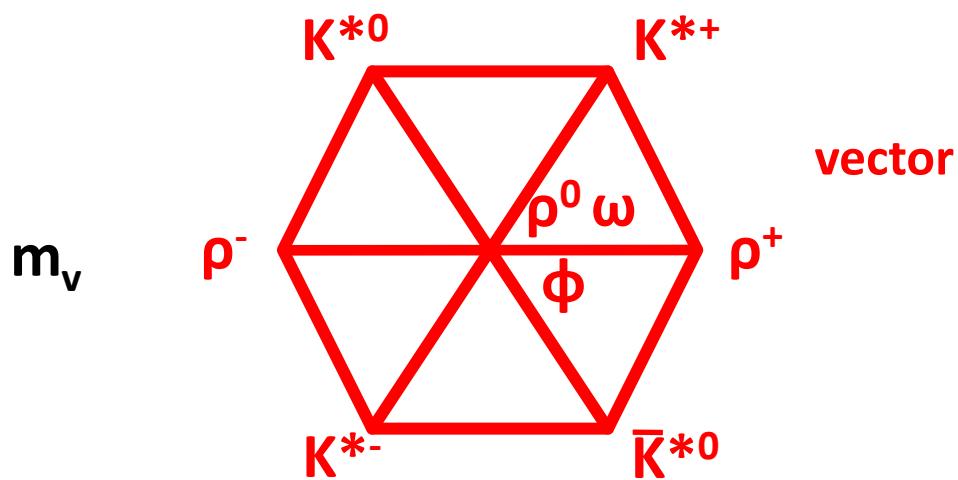
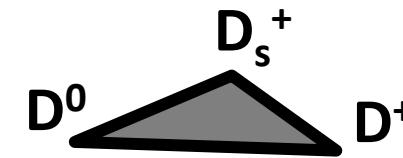
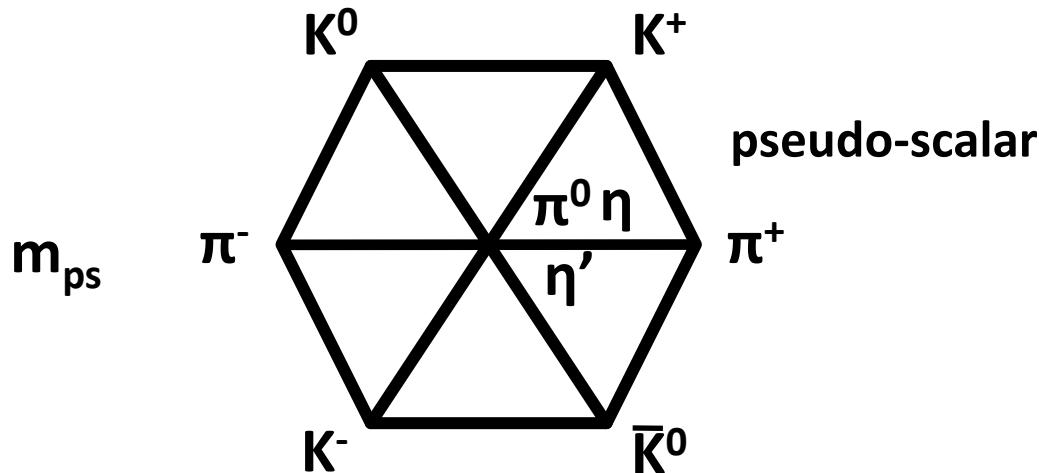
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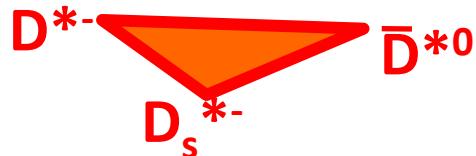
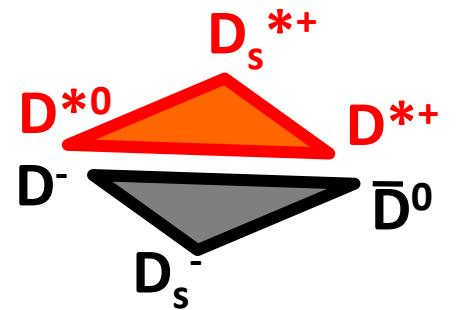
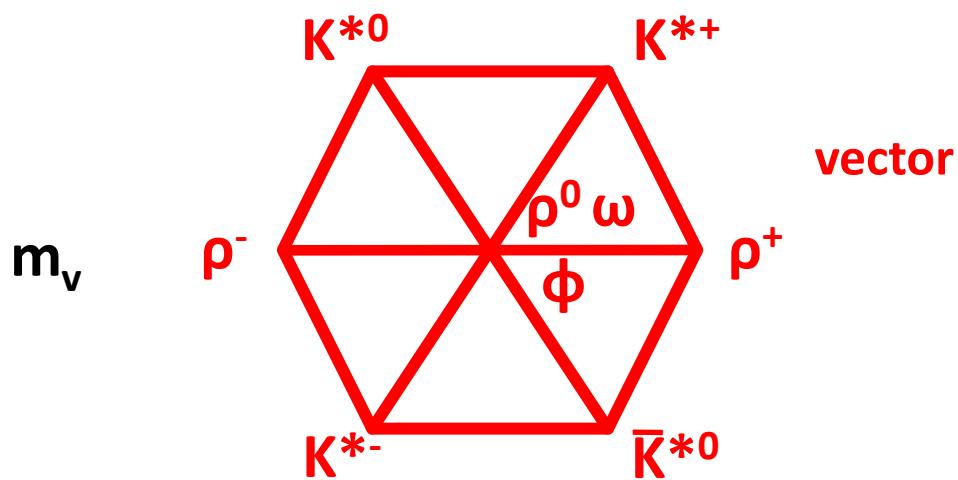
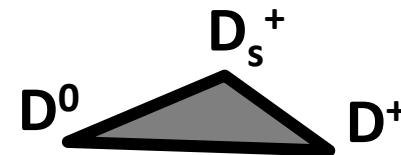
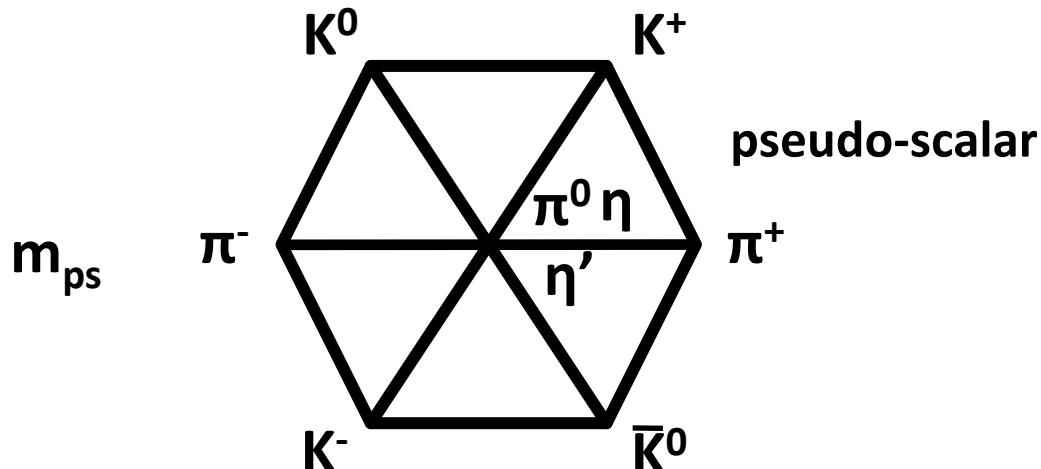
SU(4)_f symmetry ?



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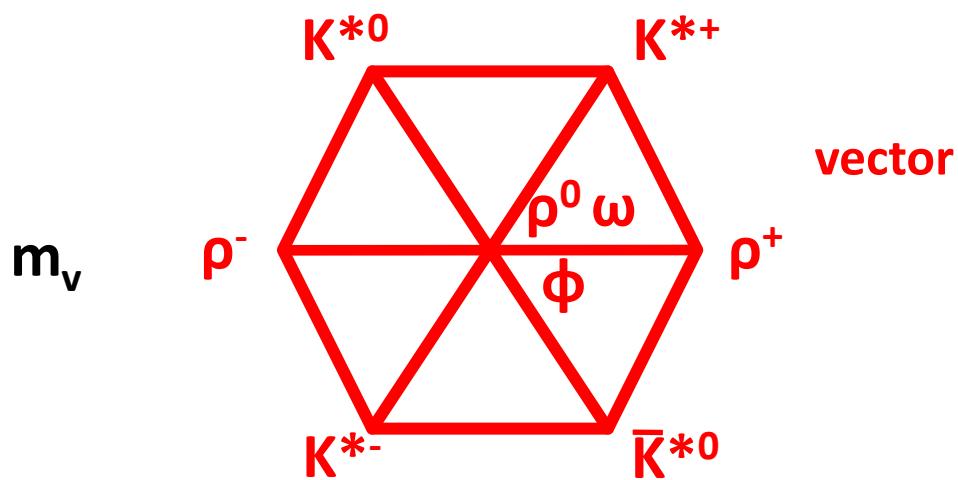
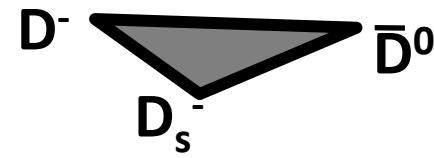
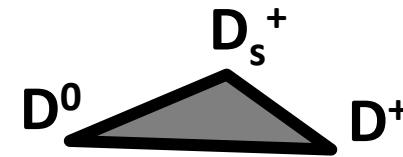
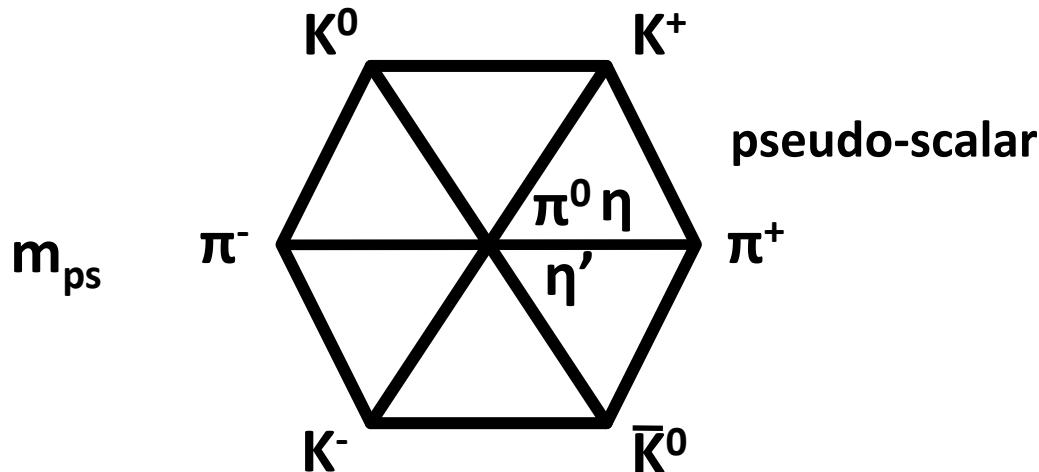
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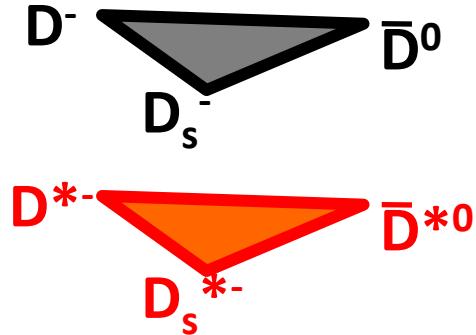
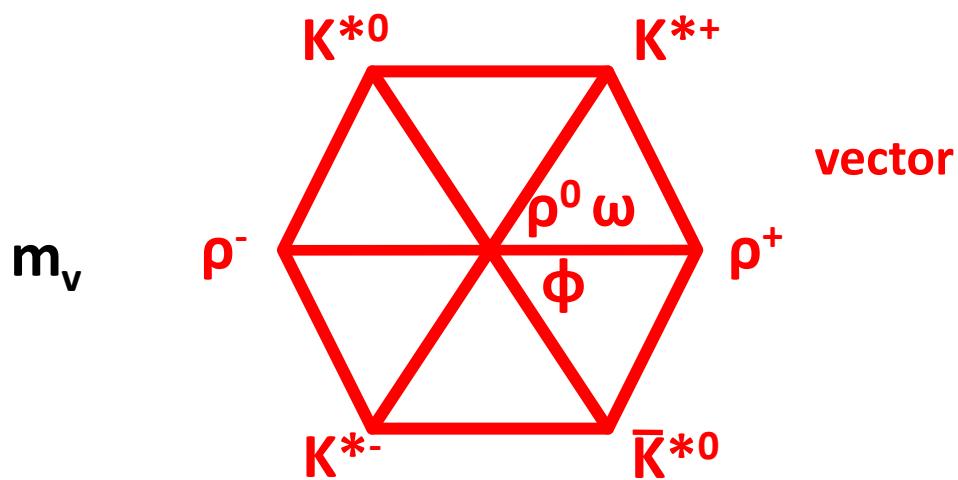
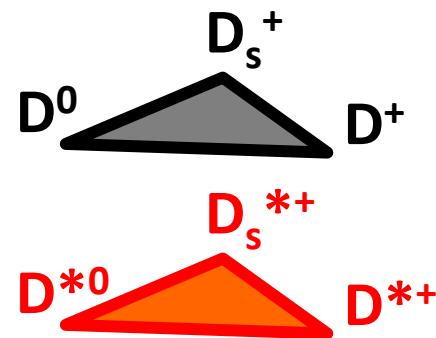
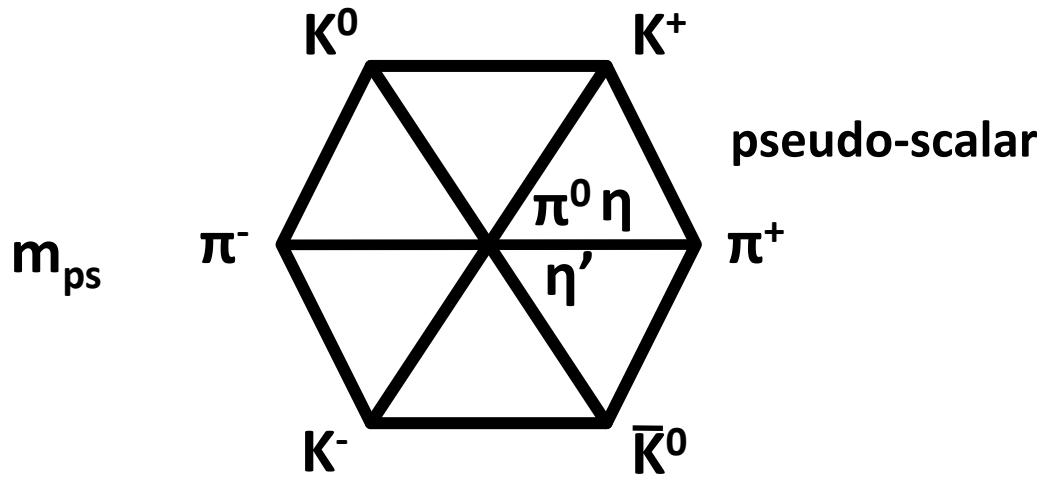
SU(4)_f symmetry ?



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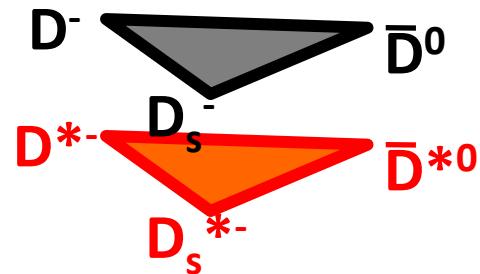
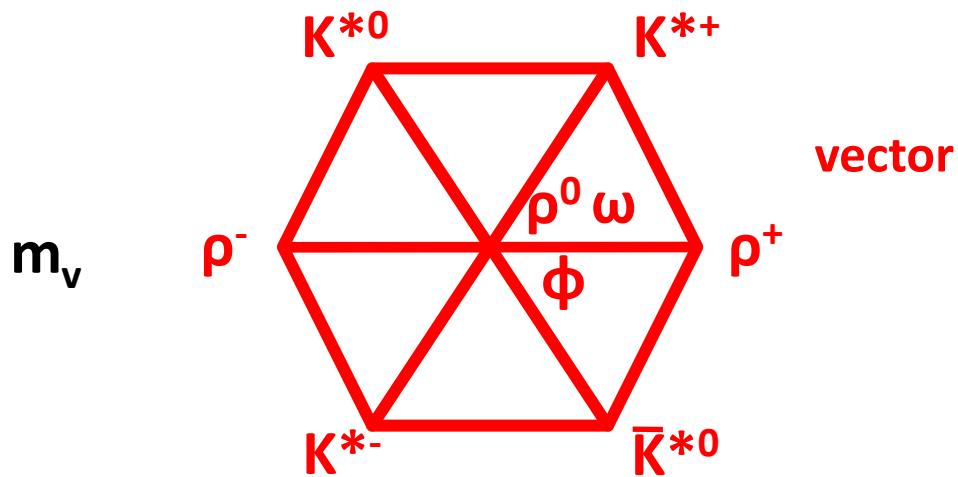
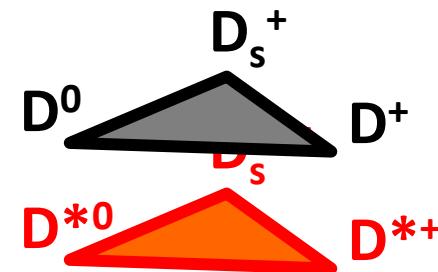
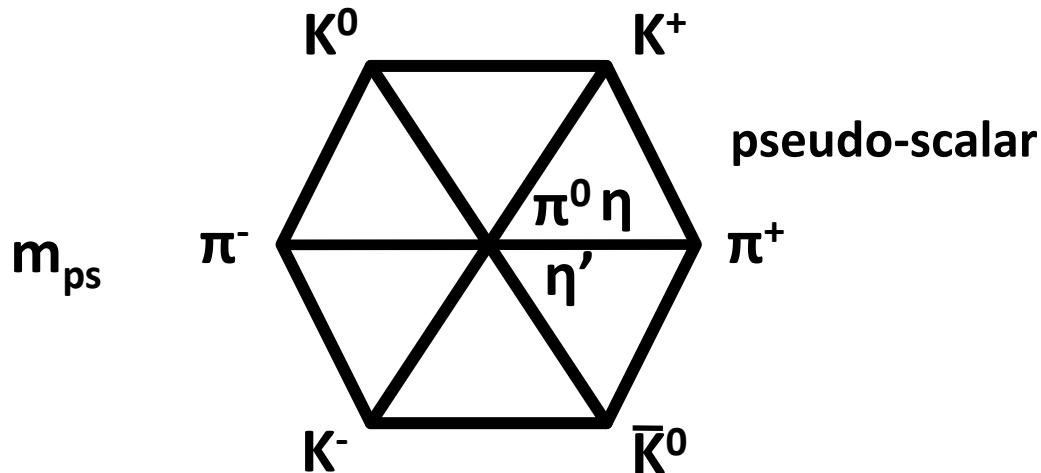
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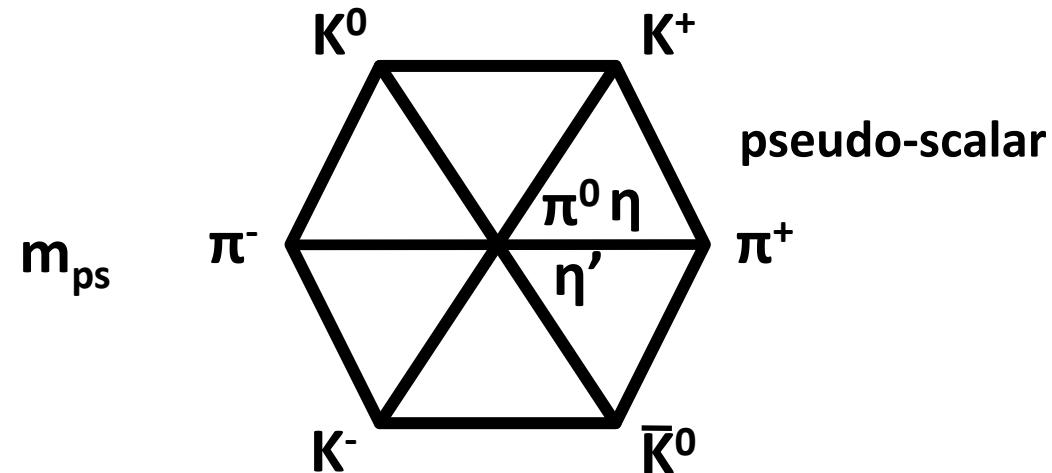
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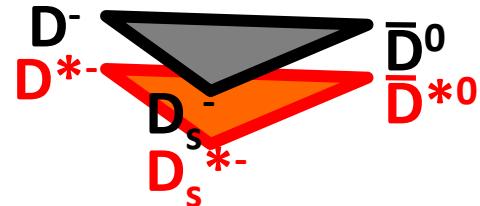
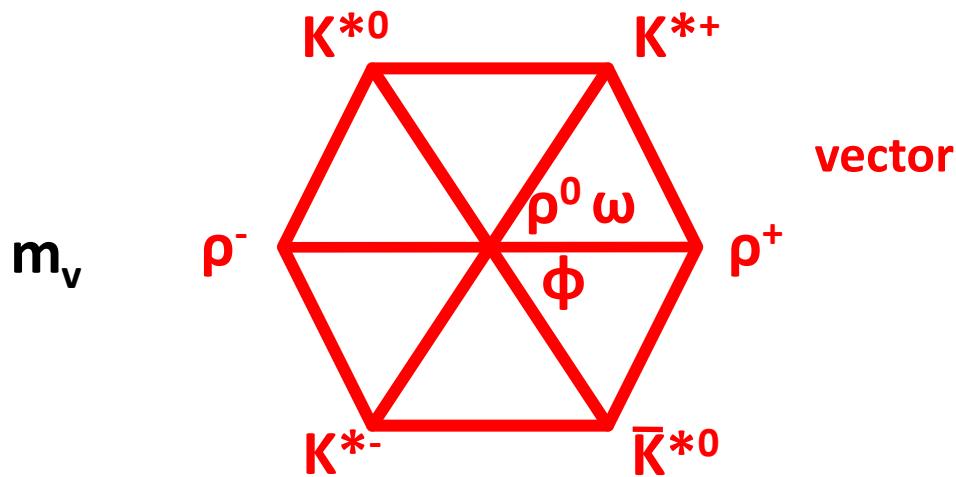
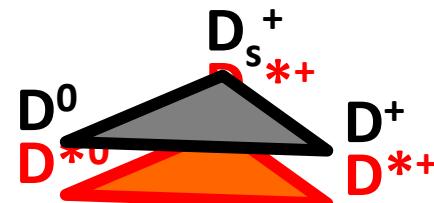
SU(4)_f symmetry ?



SU(3)_f symmetry

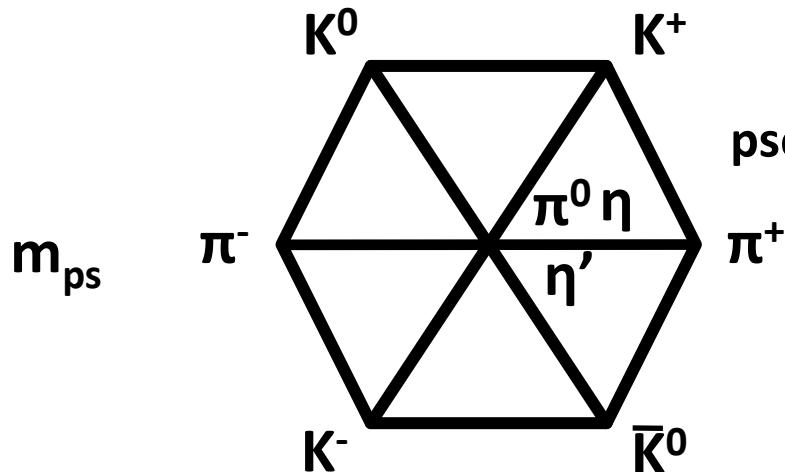


SU(4)_f symmetry ?

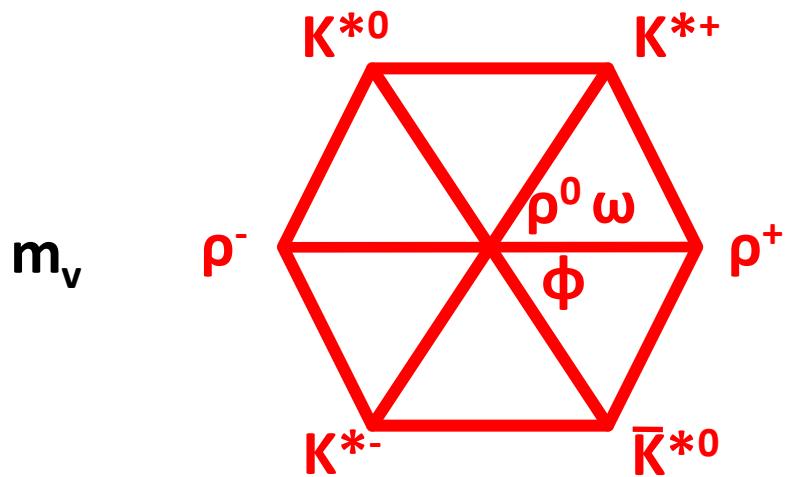


SU(3)_f symmetry

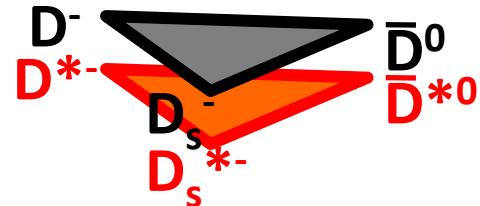
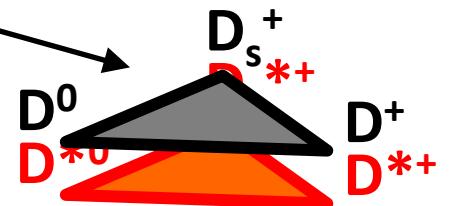
SU(4)_f symmetry ?



pseudo-scalar

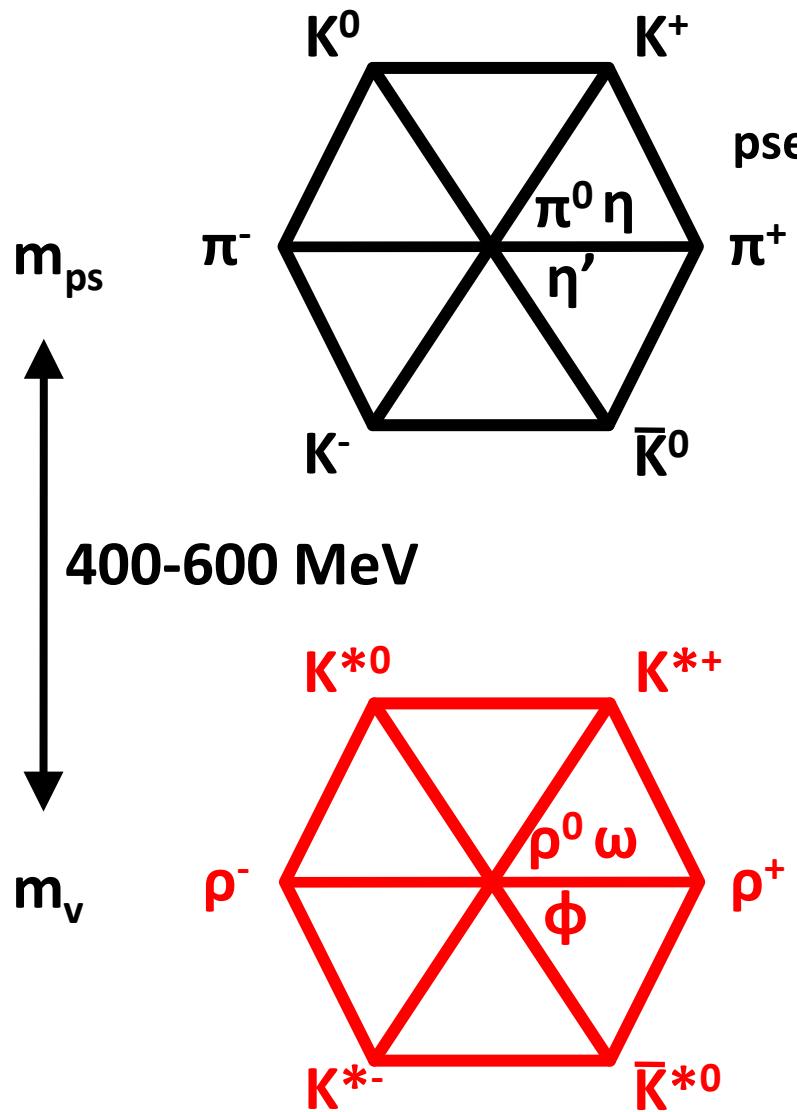


vector



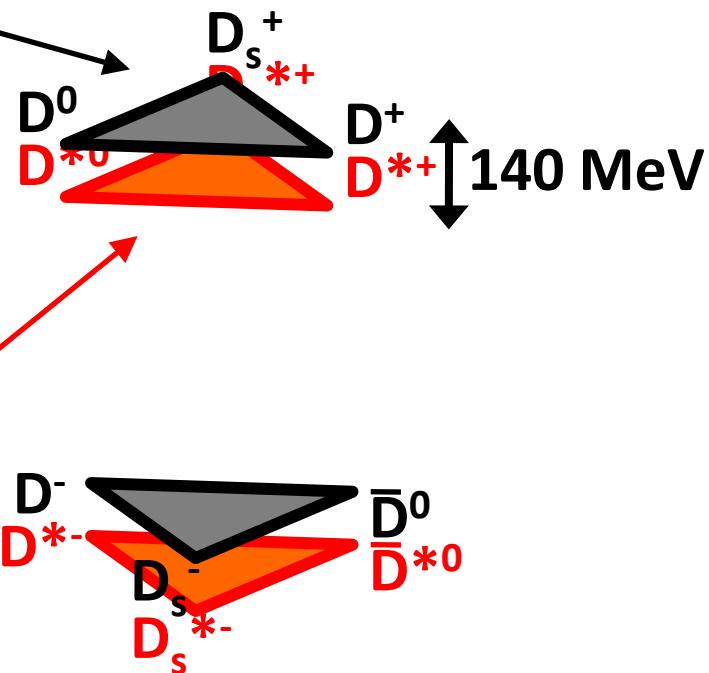
SU(3)_f symmetry

SU(4)_f symmetry ?



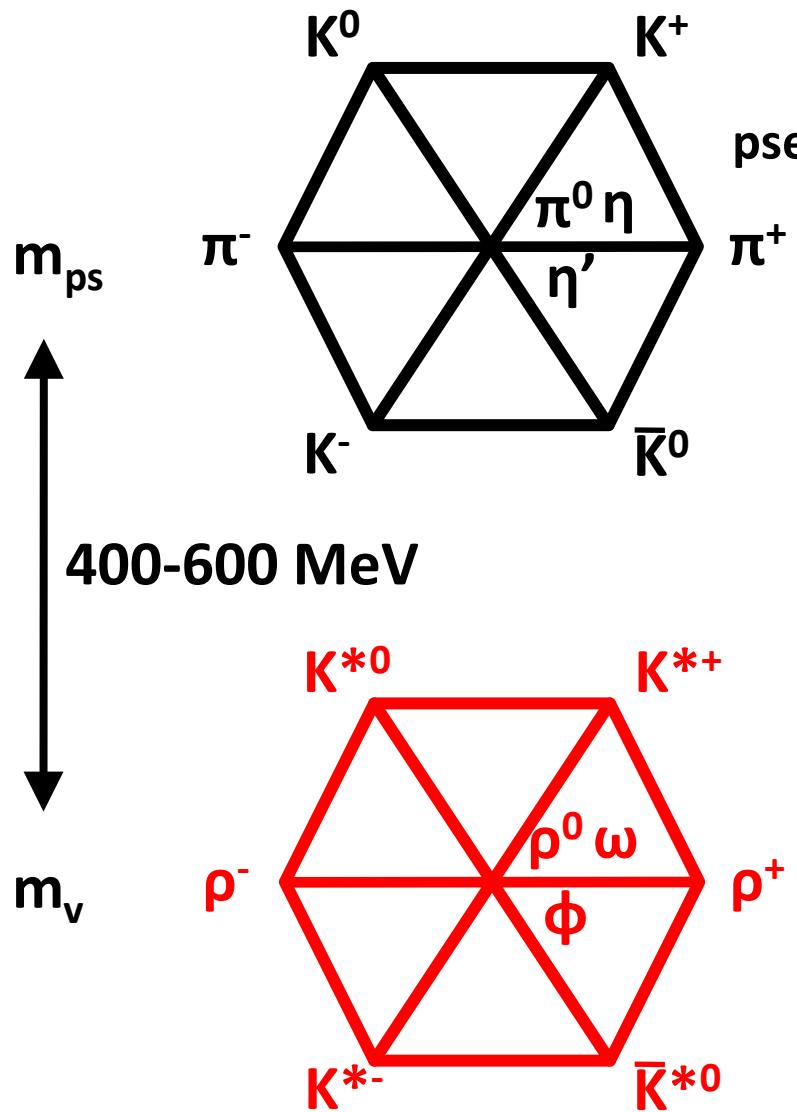
pseudo-scalar

vector



SU(3)_f symmetry

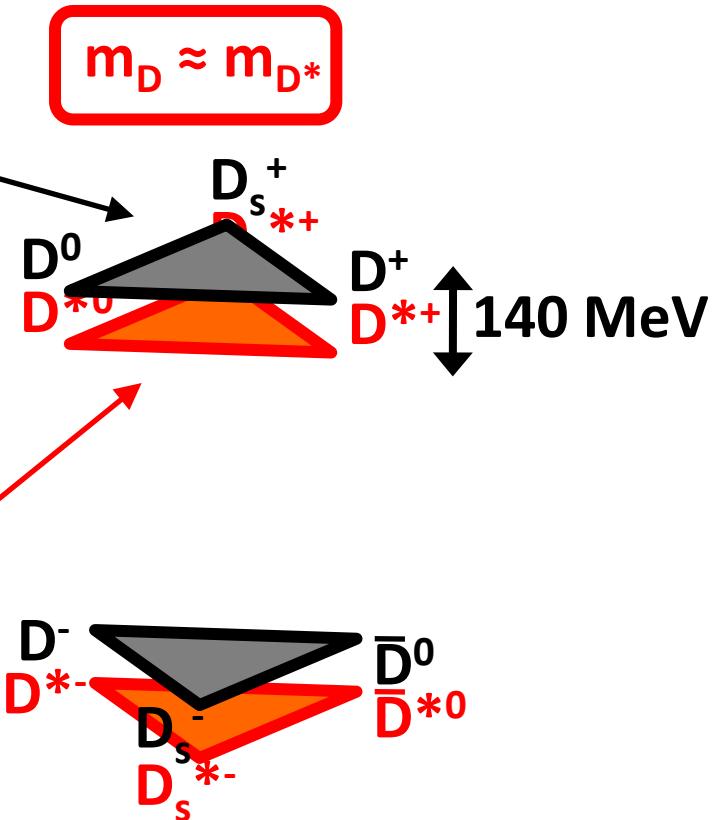
SU(4)_f symmetry ?



pseudo-scalar

$$m_D \approx m_{D^*}$$

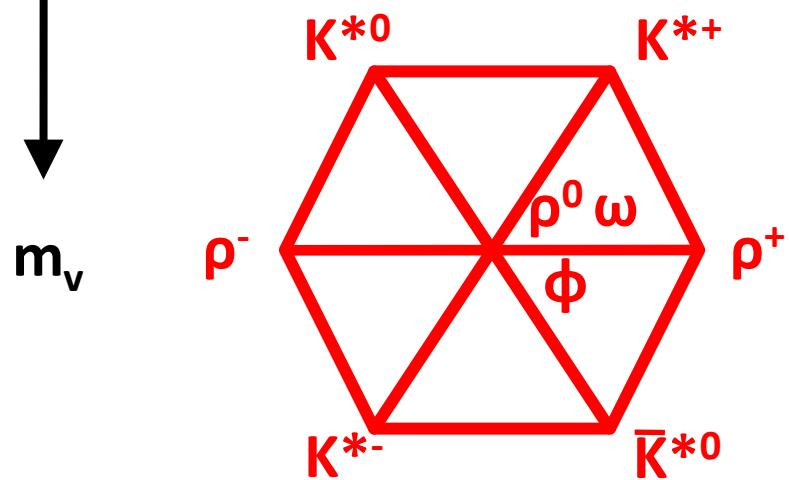
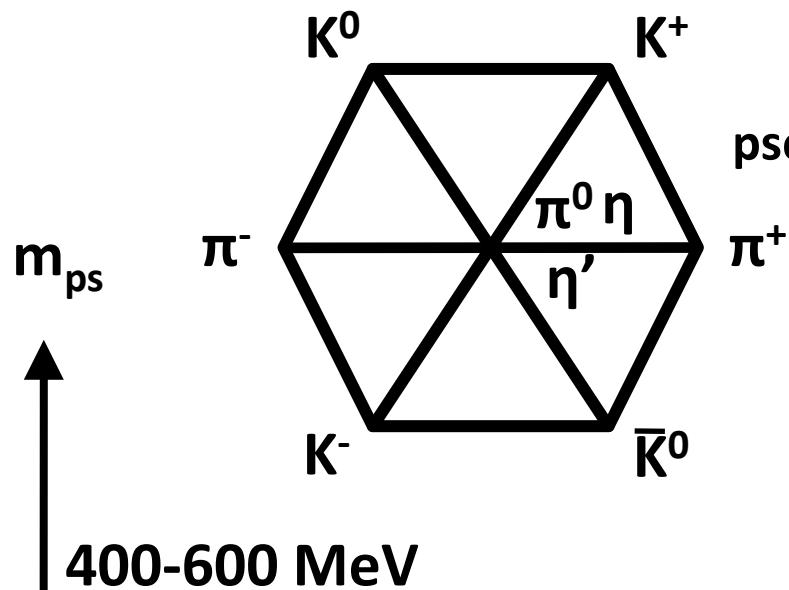
vector



$SU(3)_f$ symmetry

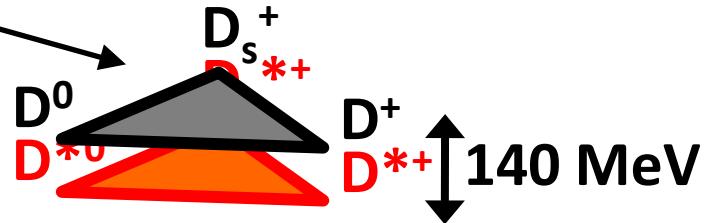
~~$SU(4)_f$ symmetry ?~~

Heavy Quark Symmetry



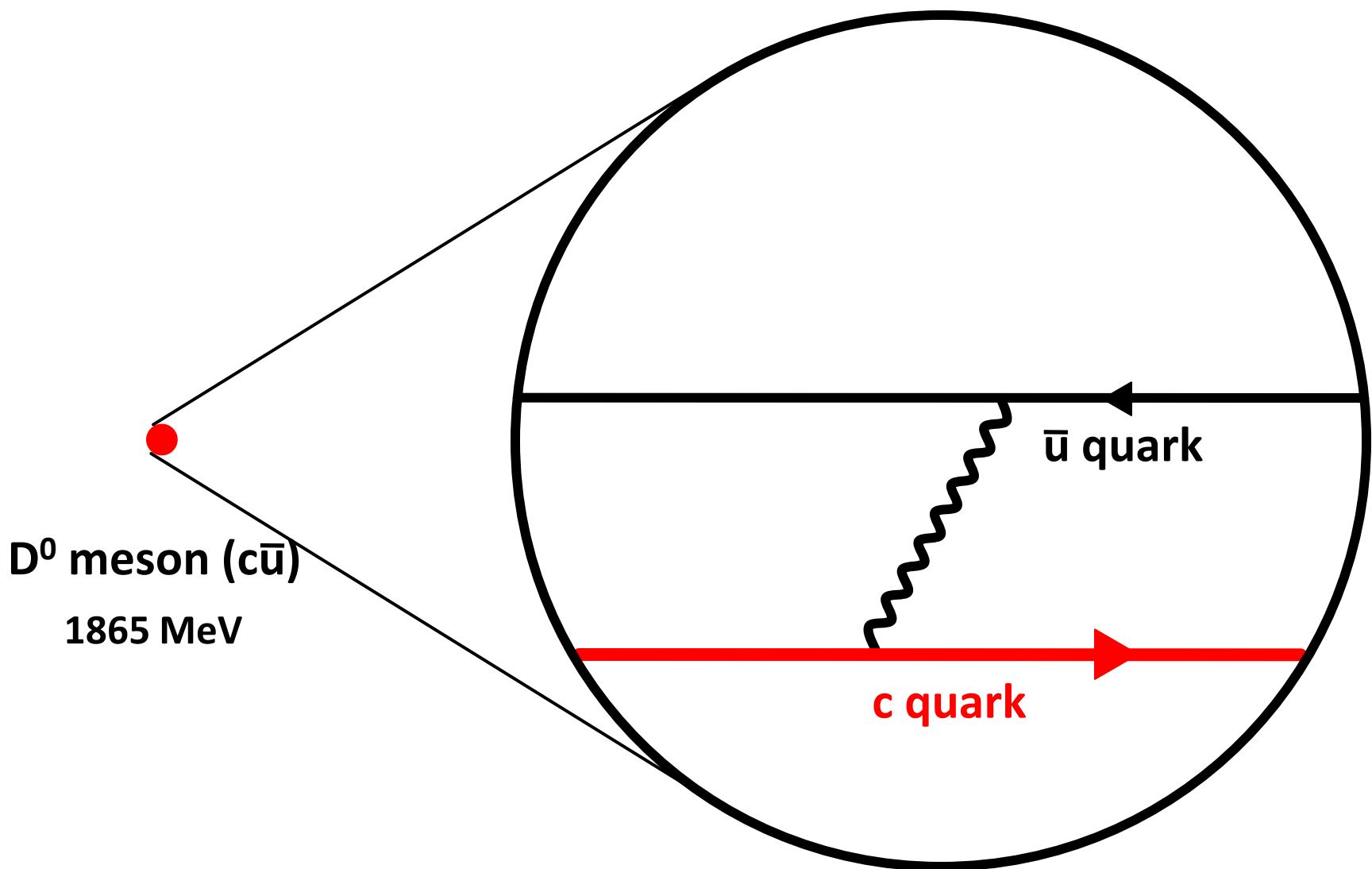
pseudo-scalar

$$m_D \approx m_{D^*}$$

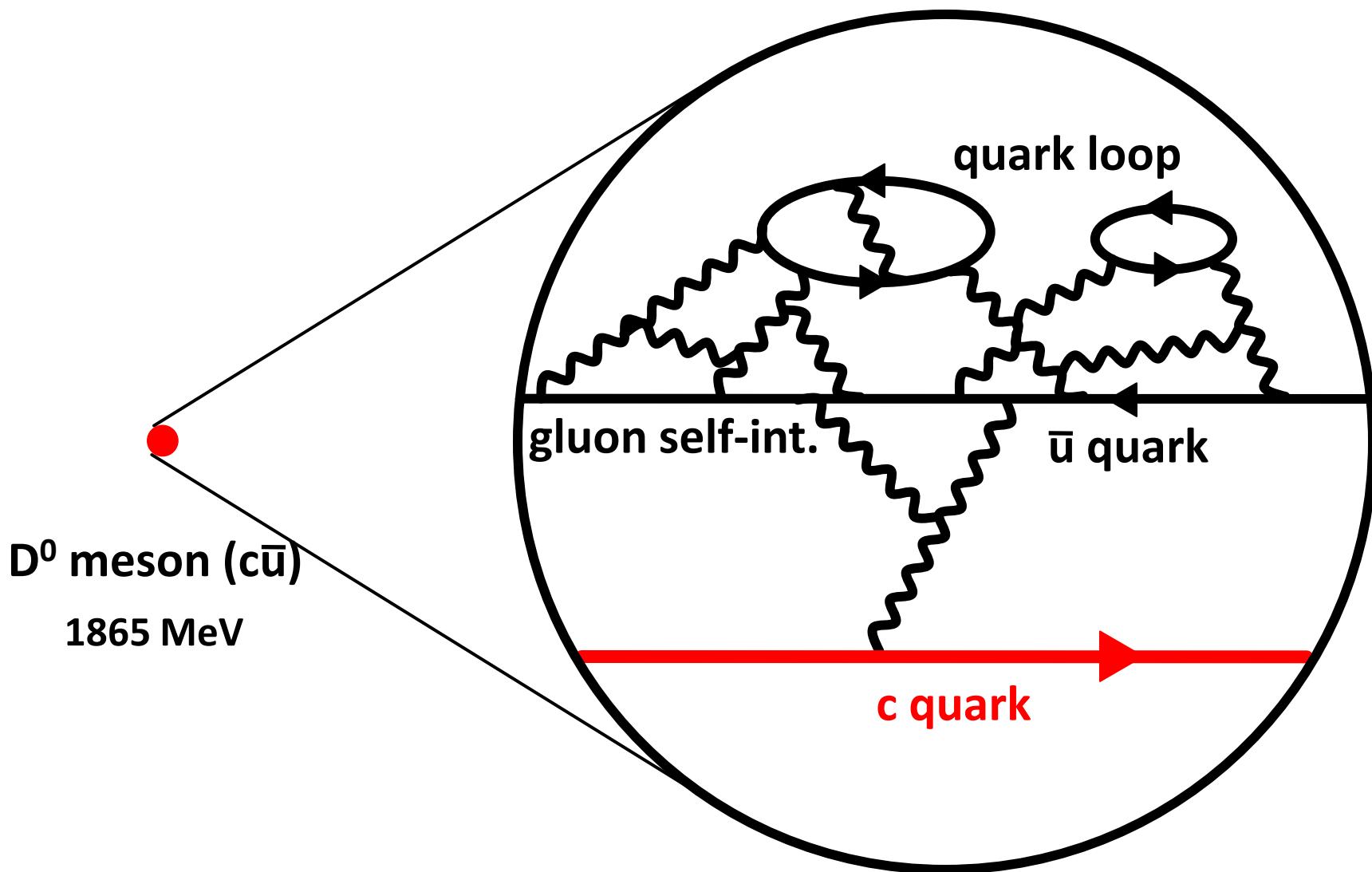


vector

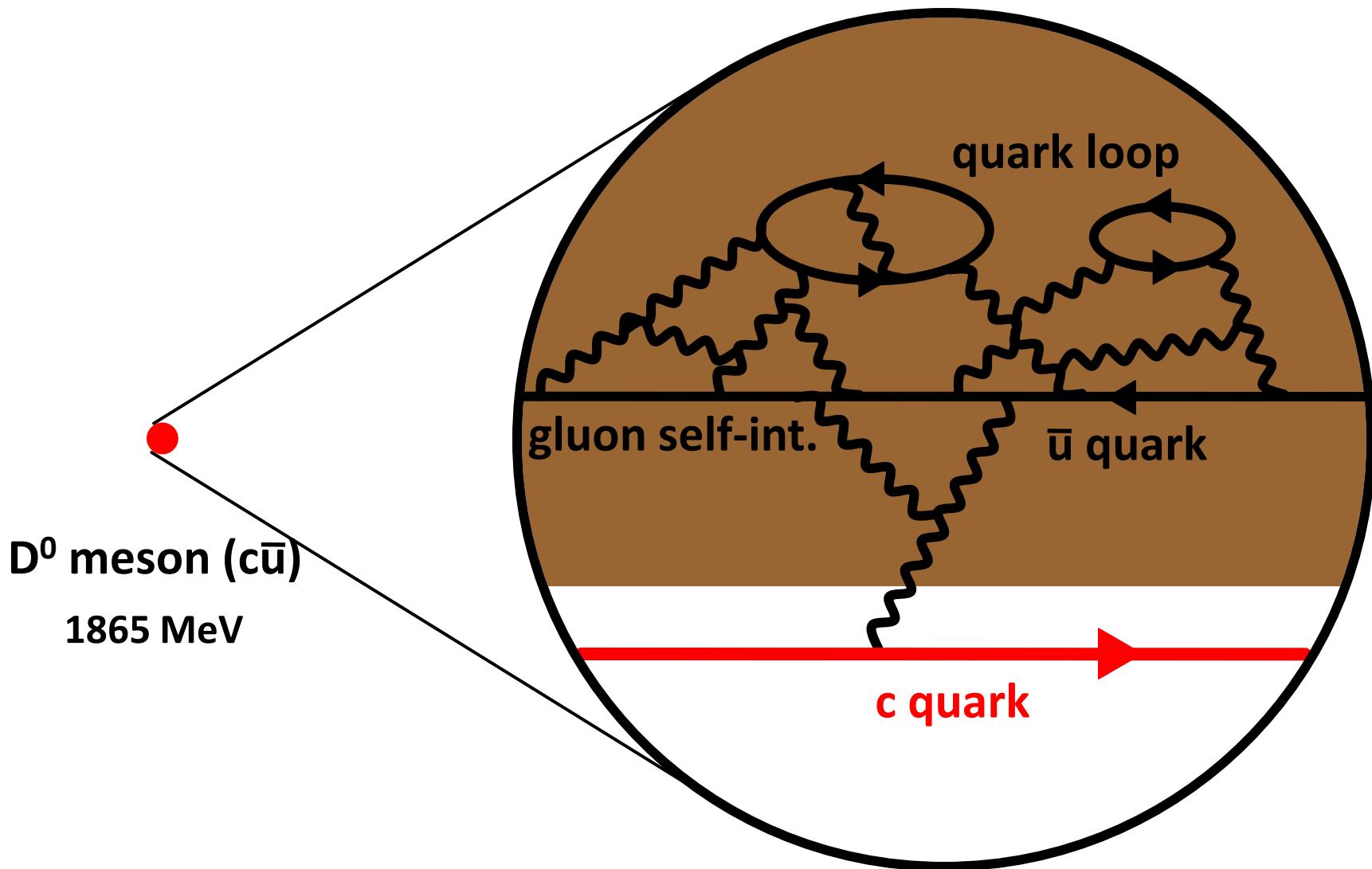
Heavy Quark Symmetry



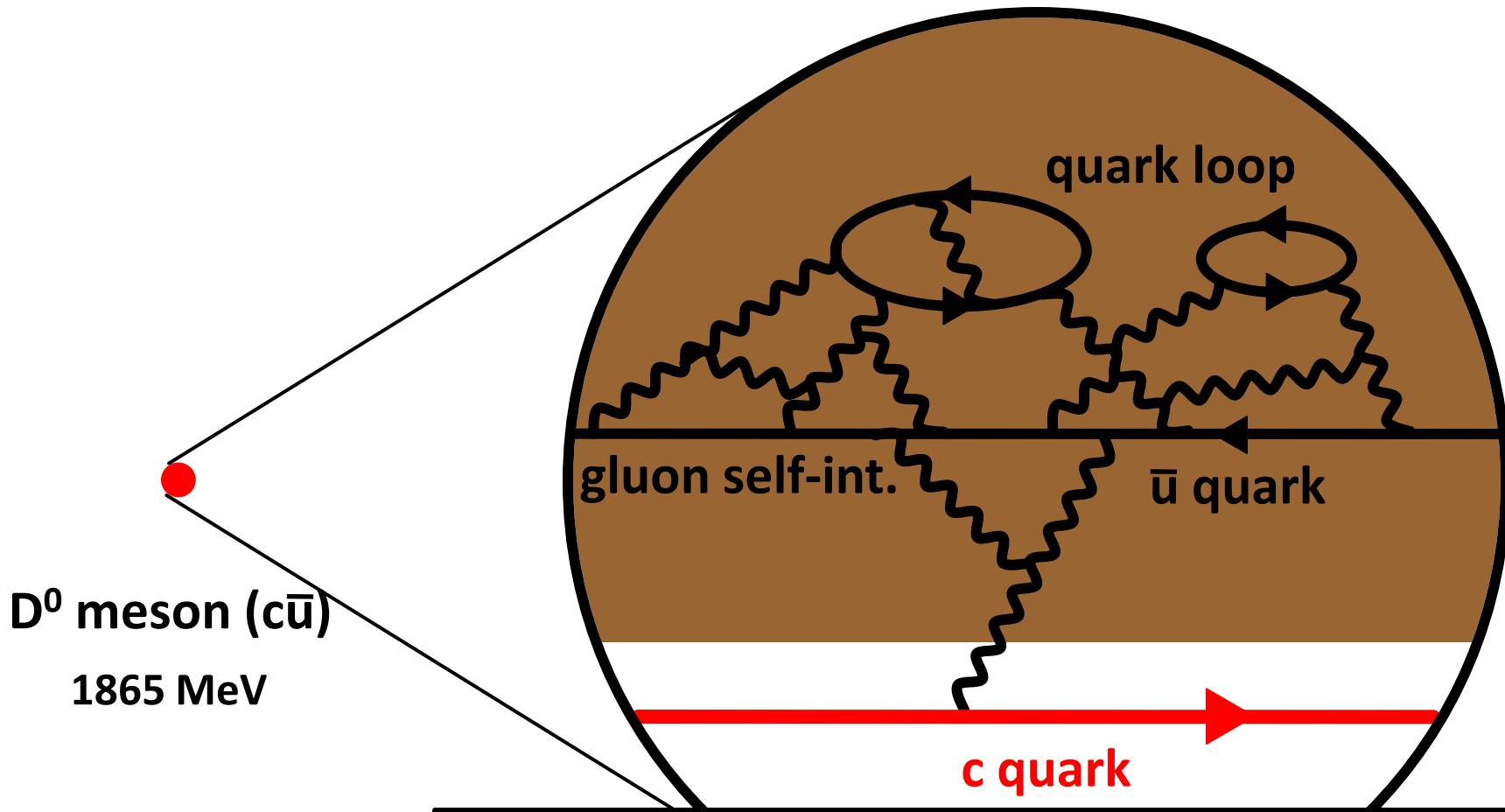
Heavy Quark Symmetry



Heavy Quark Symmetry

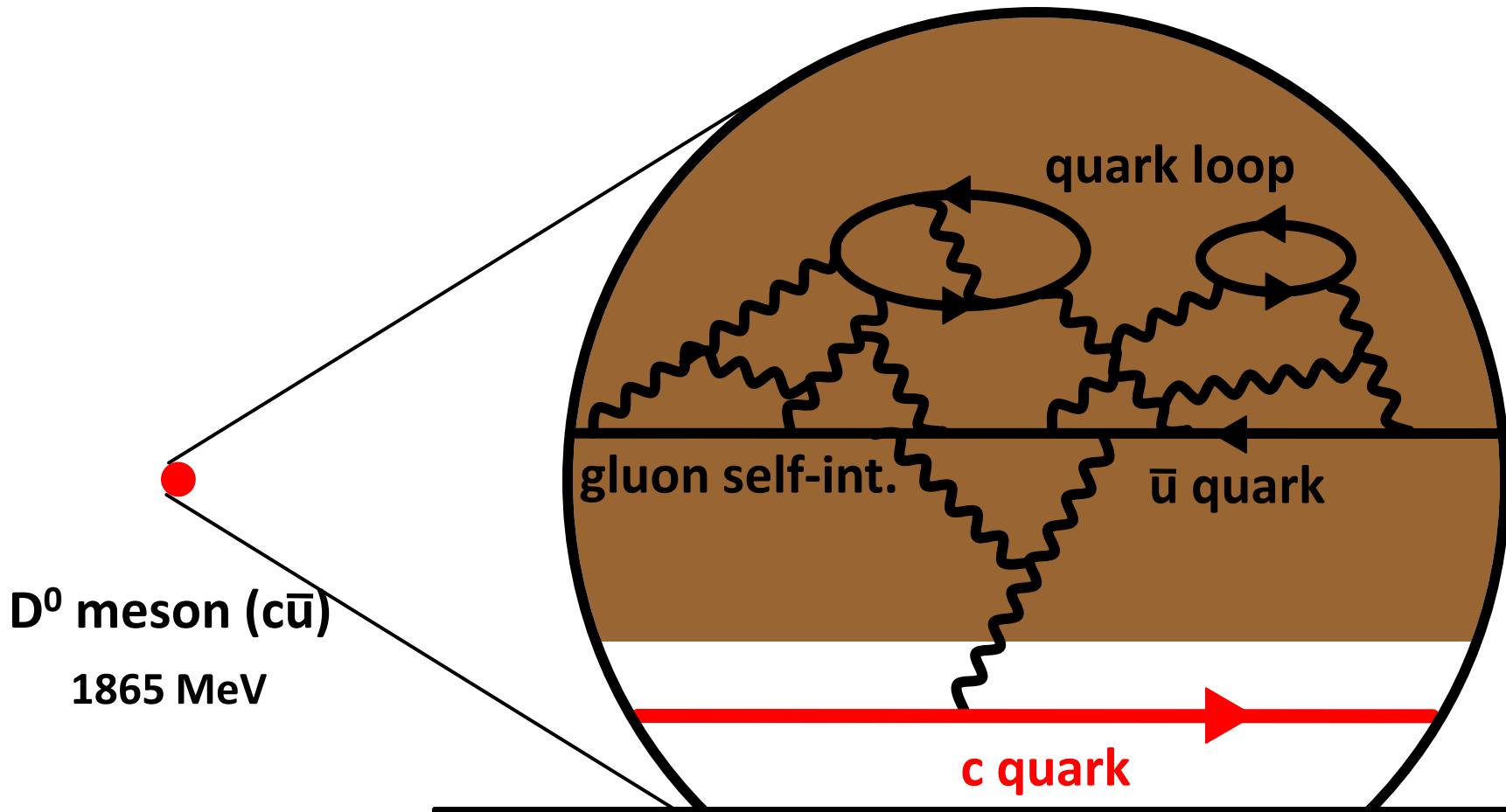


Heavy Quark Symmetry



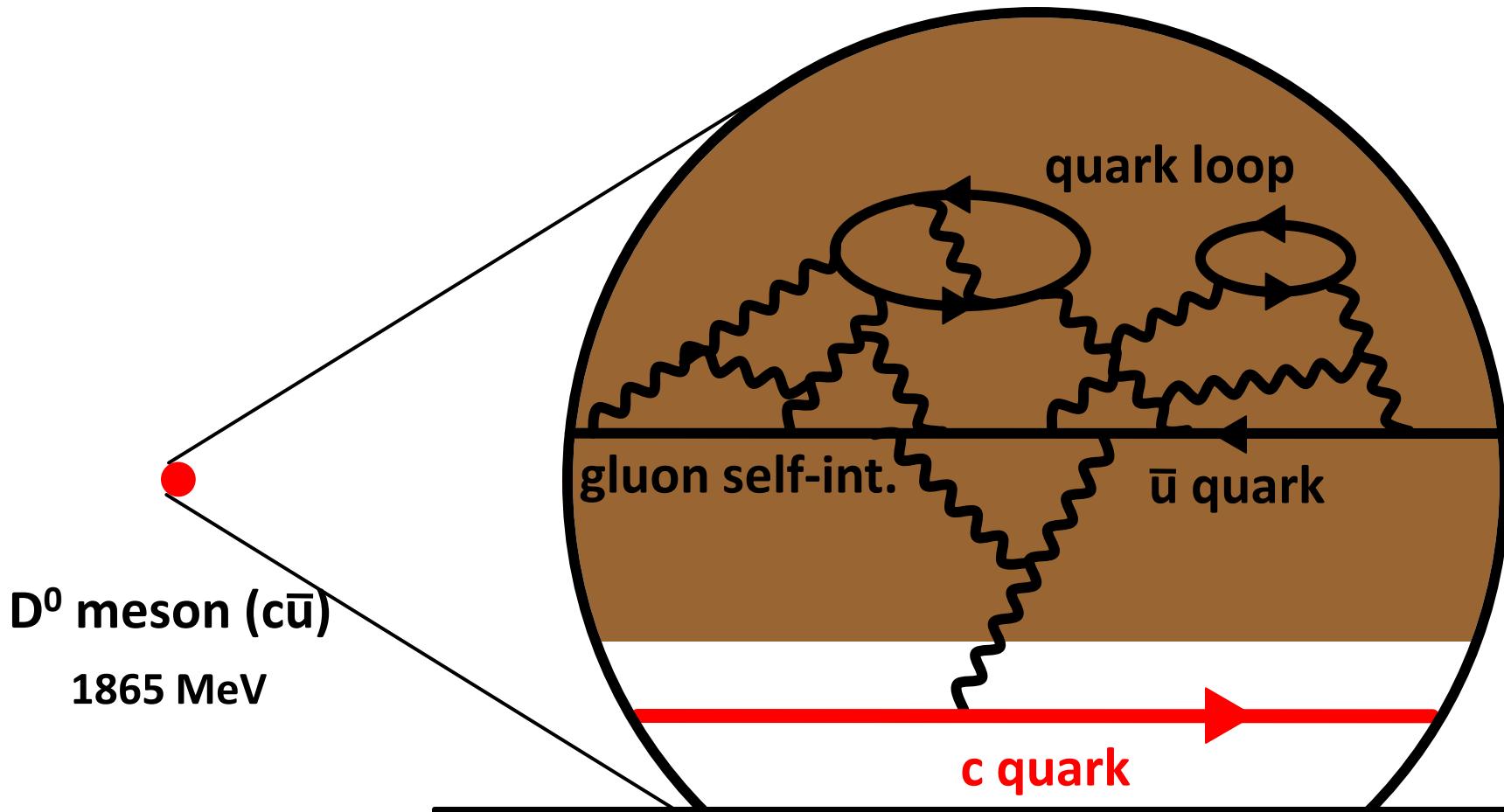
$$\begin{aligned}\mathcal{L}_{\text{HQET}} = & \bar{Q}_v i v \cdot D Q_v \\ & - \bar{Q}_v \frac{D_\perp^2}{2m_Q} Q_v - \alpha(\mu) g \bar{Q}_v \frac{\sigma_{\mu\nu} G^{\mu\nu}}{4m_Q} Q_v + \dots\end{aligned}$$

Heavy Quark Symmetry



$$\begin{aligned}\mathcal{L}_{\text{HQET}} = & \bar{Q}_v i v \cdot D Q_v \\ & - \bar{Q}_v \frac{D_\perp^2}{2m_Q} Q_v - \alpha(\mu) g \bar{Q}_v \frac{\sigma_{\mu\nu} G^{\mu\nu}}{4m_Q} Q_v + \dots\end{aligned}$$

Heavy Quark Symmetry



$$\mathcal{L}_{\text{HQET}} = \bar{Q}_v i v \cdot D Q_v$$

$$- \bar{Q}_v \frac{D_\perp^2}{2m_Q} Q_v - \alpha(\mu) g \bar{Q}_v \frac{\sigma_{\mu\nu} \Sigma}{4m_Q} Q_v + \dots$$

$m_D \approx m_{D^*}$

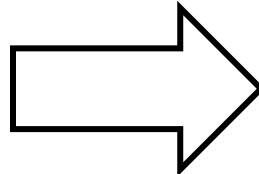
2. Symmetry and dynamics of heavy quark hadrons

3 5 150 200
u d s Λ_{QCD}

1500 $\Sigma\Sigma$ 4700 [MeV]
c b m

Chiral Symmetry

$SU(3)_L \times SU(3)_R$



Chiral Symmetry
 $SU(4)_L \times SU(4)_R$

π, K as Nambu-Goldstone bosons

Chiral multiplets (ρ - a_1 , N - N^* , ...)

Nuclei formed by tensor force

Hypernuclei as g.s. of QCD matter

...

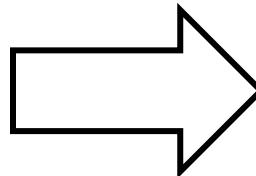
2. Symmetry and dynamics of heavy quark hadrons

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c b m

Chiral Symmetry

$SU(3)_L \times SU(3)_R$



**Heavy Quark
Symmetry**

$SU(2)_{\text{spin}} \times SU(N_h)$

π, K as Nambu-Goldstone bosons

Chiral multiplets (ρ - a_1 , N - N^* , ...)

Nuclei formed by tensor force

Hypernuclei as g.s. of QCD matter

...

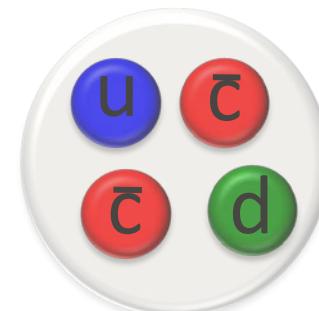
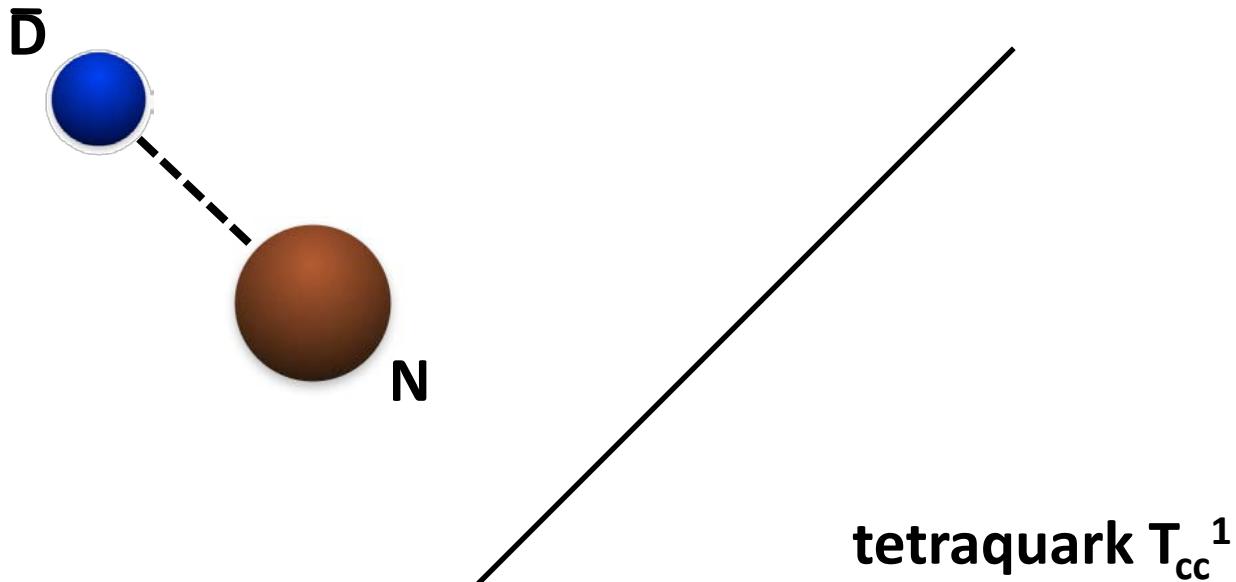
D-D* ($\Sigma c - \Sigma c^*$, $\Xi cc - \Xi cc^*$) mass degeneracy (heavy quark sym.)

More exotics? (like Tcc with no pion decay)

Charmed nuclei ?

3. New exotic hadrons and nuclei

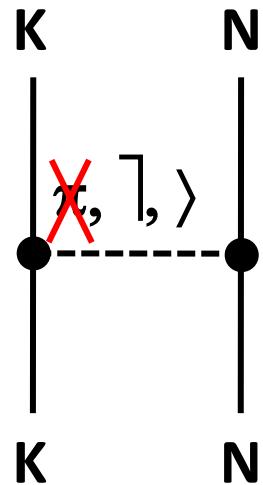
\bar{D} and nucleon



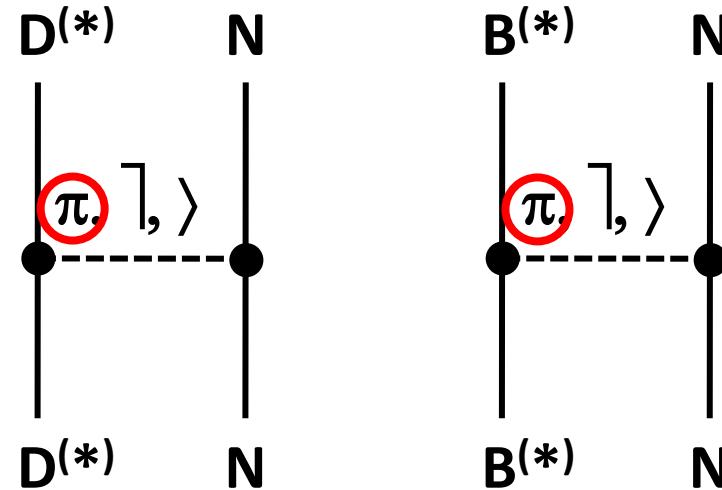
3. New exotic hadrons and nuclei

D and nucleon

KN interaction



D(B)N interaction



One pion is absent.
(short range force)

Weinberg-Tomozawa

One pion is present.
(long range force)

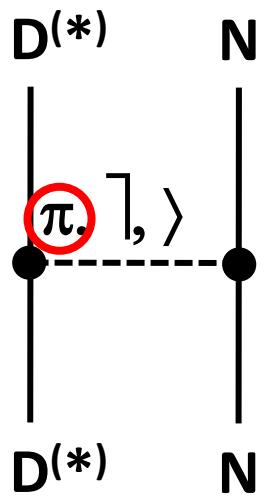
One pion exchange potential
(OPEP)

S. Y. and K. Sudoh, PRD80, 034008 (2009)

Y. Yamaguchi, S. Ohkoda, S. Y., A. Hosaka, arXiv:1105.0734

3. New exotic hadrons and nuclei

D and nucleon



Heavy Quark Symmetry

$$\mathcal{L}_{\pi HH} = g \operatorname{tr} \bar{H}_a H_b \gamma_\nu \gamma_5 A_{ba}^\nu$$

$$H_a = \frac{1 + \not{v}}{2} [P_{a\mu}^* \gamma^\mu - P_a \gamma_5]$$

vector + pseudoscalar

G. Burdman and J.F. Donoghue (1992)
 M.B. Wise (1992)
 T.-M. Yan, H.-Y. Cheng, C.-Y. Cheung,
 G.-L. Lin, Y.C. Lin and H.-L. Yu (1997)

with $\bar{H}_a = \gamma_0 H_a^\dagger \gamma_0$

$P=D, B$

π -exchange pot.

S. Y. and K. Sudoh, PRD80, 034008 (2009)

$$V_{PN \rightarrow P^* N} = -\frac{gg_{\pi NN}}{\sqrt{2}m_N f} \frac{1}{4\pi} \frac{\mu^2}{3} \left[\vec{\varepsilon}^{(\lambda)\dagger} \cdot \vec{\sigma} C(r; \mu) + S_{\varepsilon^{(\lambda)}}^\dagger T(r; \mu) \right] \vec{\tau}_P \cdot \vec{\tau}_N$$

$$V_{P^* N \rightarrow P^* N} = \frac{gg_{\pi NN}}{\sqrt{2}m_N f} \frac{1}{4\pi} \frac{m_\pi^2}{3} \left[\vec{T} \cdot \vec{\sigma} C(r; m_\pi) + S_T T(r; m_\pi) \right] \vec{\tau}_P \cdot \vec{\tau}_N$$

ω and ρ -exchange pot.

Y. Yamaguchi, S. Ohkoda, S. Y., A. Hosaka,
 arXiv:1105.0734

3. New exotic hadrons and nuclei

D and nucleon

$$I(J^P)=0(1/2^-) \quad \bar{D}N(^2S_{1/2}), \bar{D}^*N(^2S_{1/2}), \bar{D}^*N(^4D_{1/2})$$

$$V_{1/2^-}^\pi = \frac{g_\pi g_{\pi NN}}{\sqrt{2}m_N f_\pi} \frac{1}{3} \begin{pmatrix} 0 & \sqrt{3}C_{m_\pi} & -\sqrt{6}T_{m_\pi} \\ \sqrt{3}C_{m_\pi} & -2C_{m_\pi} & -\sqrt{2}T_{m_\pi} \\ -\sqrt{6}T_{m_\pi} & -\sqrt{2}T_{m_\pi} & C_{m_\pi} - 2T_{m_\pi} \end{pmatrix} \vec{\tau}_P \cdot \vec{\tau}_N$$

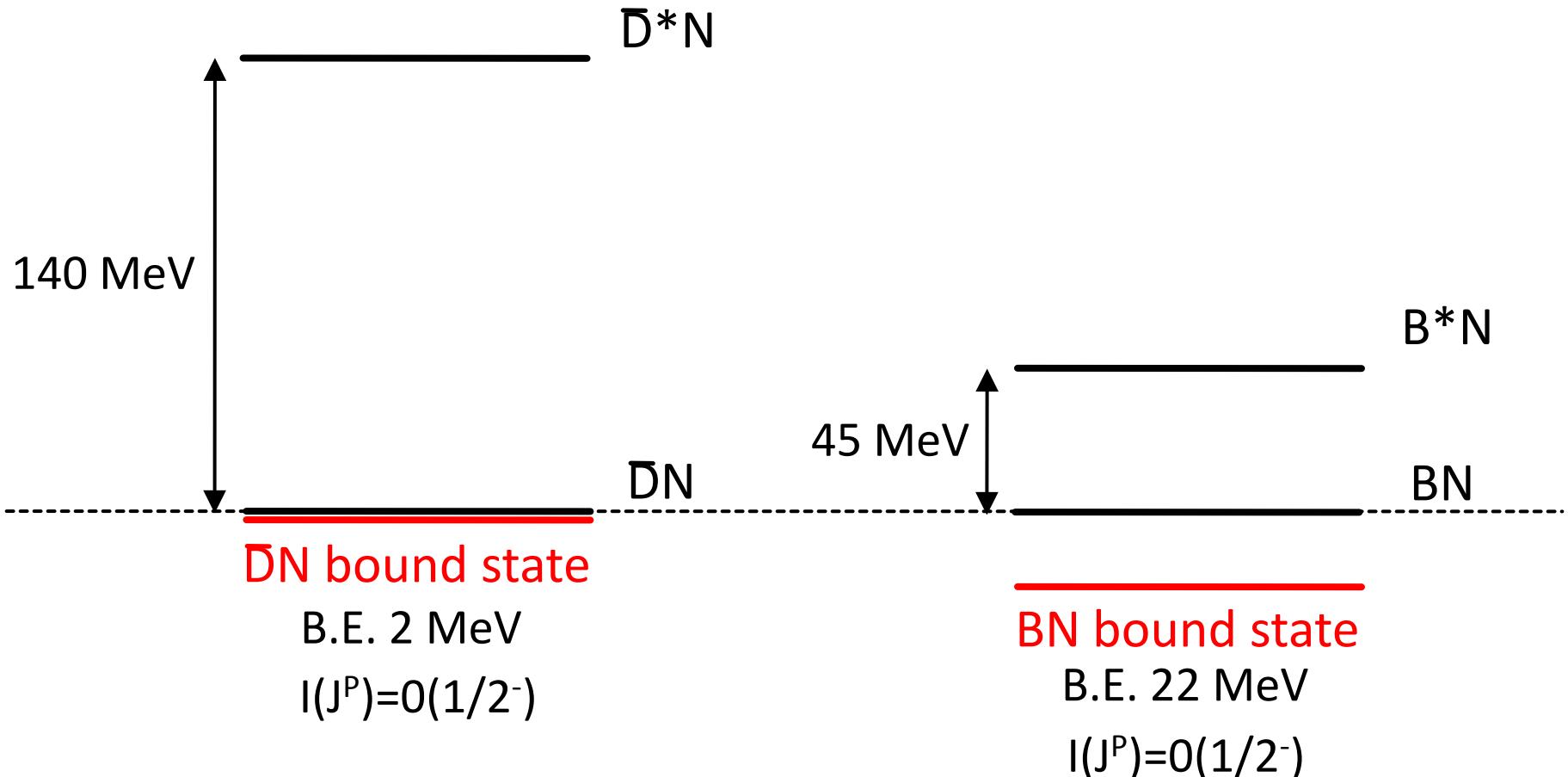
$$I(J^P)=0(3/2^-) \quad \bar{D}N(^2D_{3/2}), \bar{D}^*N(^4S_{3/2}), \bar{D}^*N(^4D_{3/2}), \bar{D}^*N(^2D_{3/2})$$

$$V_{3/2^-}^\pi = \frac{g_\pi g_{\pi NN}}{\sqrt{2}m_N f_\pi} \frac{1}{3} \begin{pmatrix} 0 & \sqrt{3}T_{m_\pi} & -\sqrt{3}T_{m_\pi} & \sqrt{3}C_{m_\pi} \\ \sqrt{3}T_{m_\pi} & C_{m_\pi} & 2T_{m_\pi} & T_{m_\pi} \\ -\sqrt{3}T_{m_\pi} & 2T_{m_\pi} & C_{m_\pi} & -T_{m_\pi} \\ \sqrt{3}C_{m_\pi} & T_{m_\pi} & -T_{m_\pi} & -2C_{m_\pi} \end{pmatrix} \vec{\tau}_P \cdot \vec{\tau}_N$$

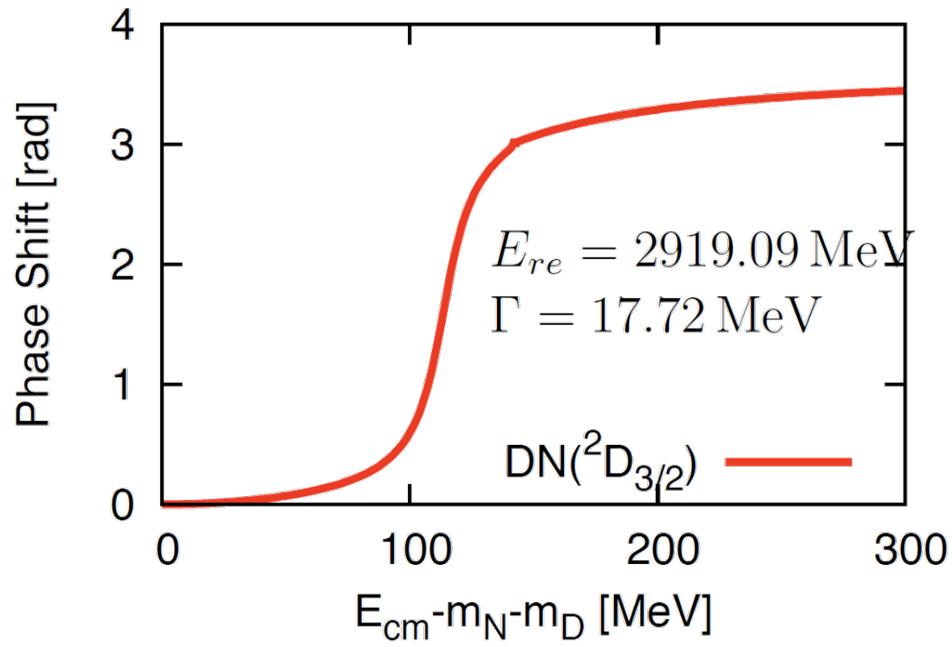
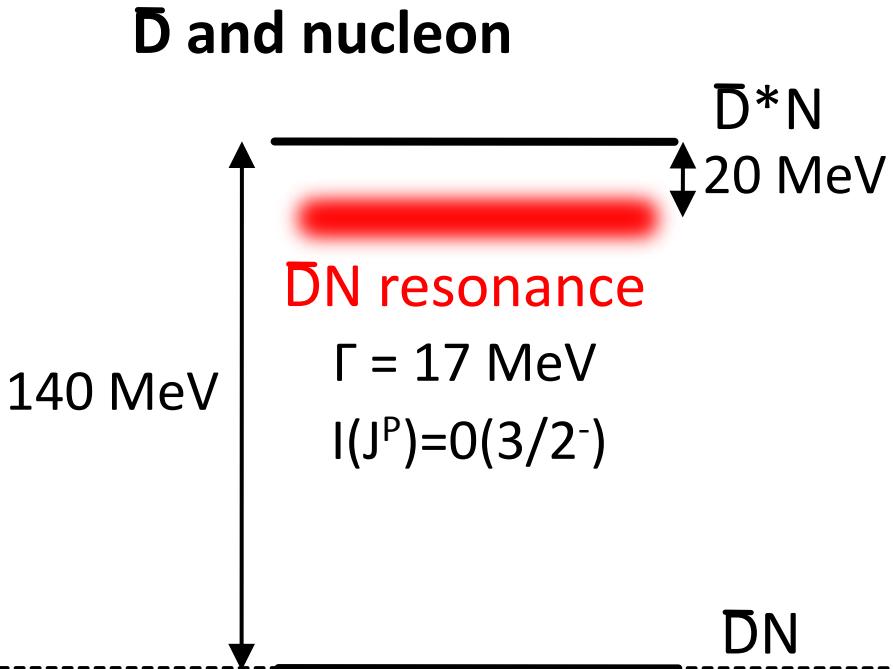
See Y. Yamaguchi, S. Ohkoda, S. Y., A. Hosaka, arXiv:1105.0734, for omega and rho-exchange potential.

3. New exotic hadrons and nuclei

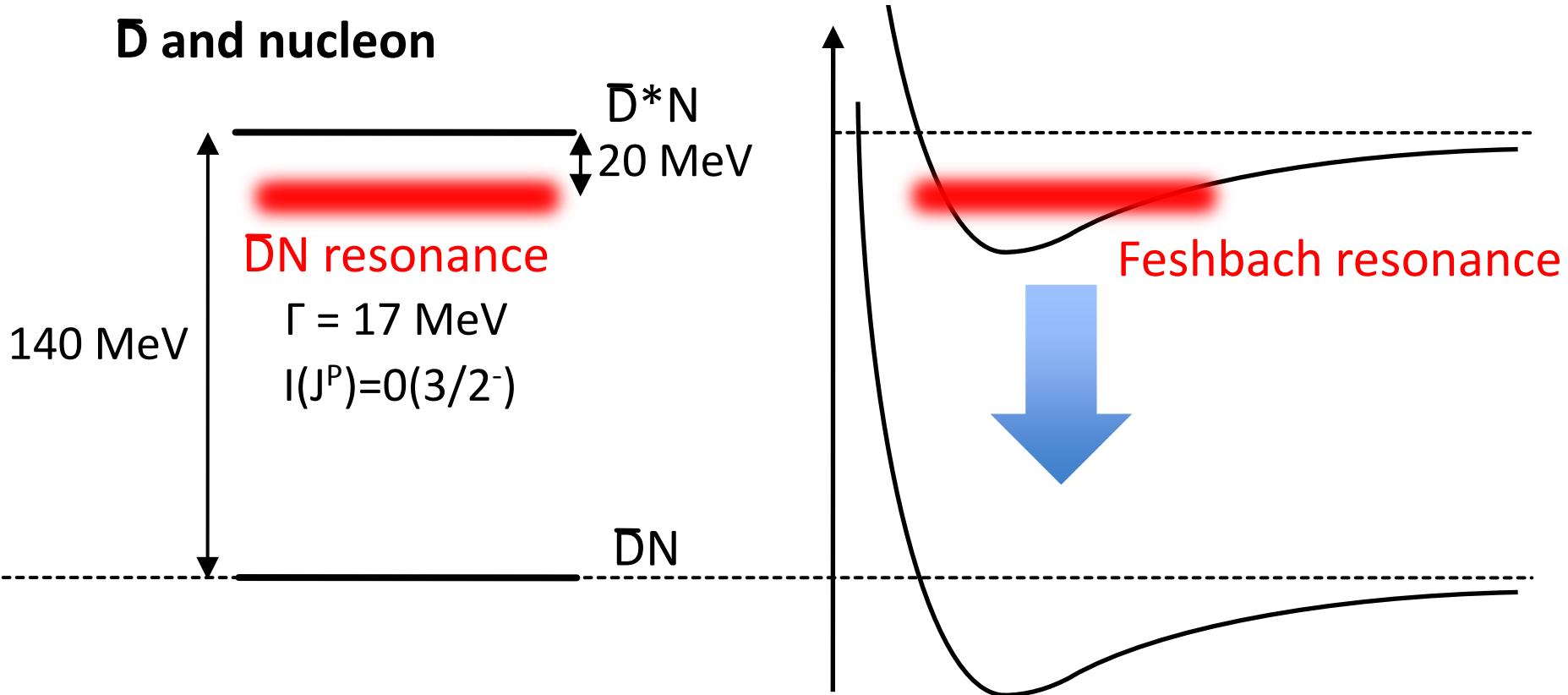
\bar{D} and nucleon



3. New exotic hadrons and nuclei

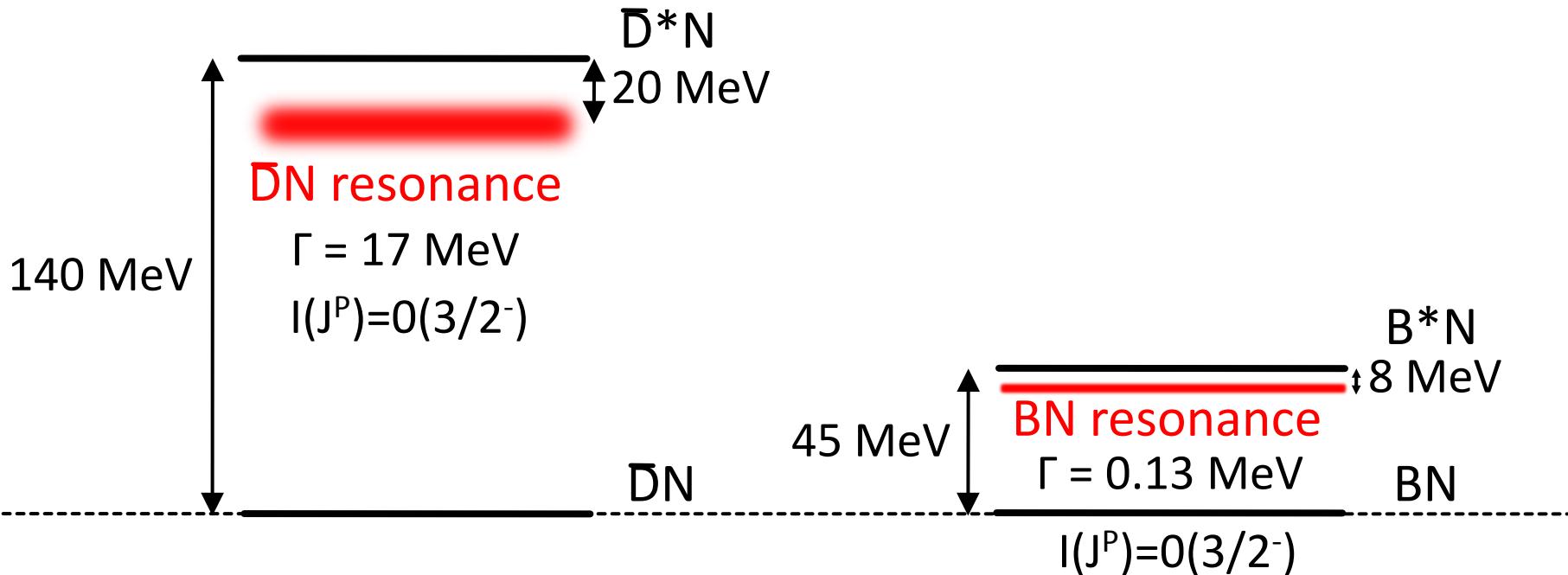


3. New exotic hadrons and nuclei



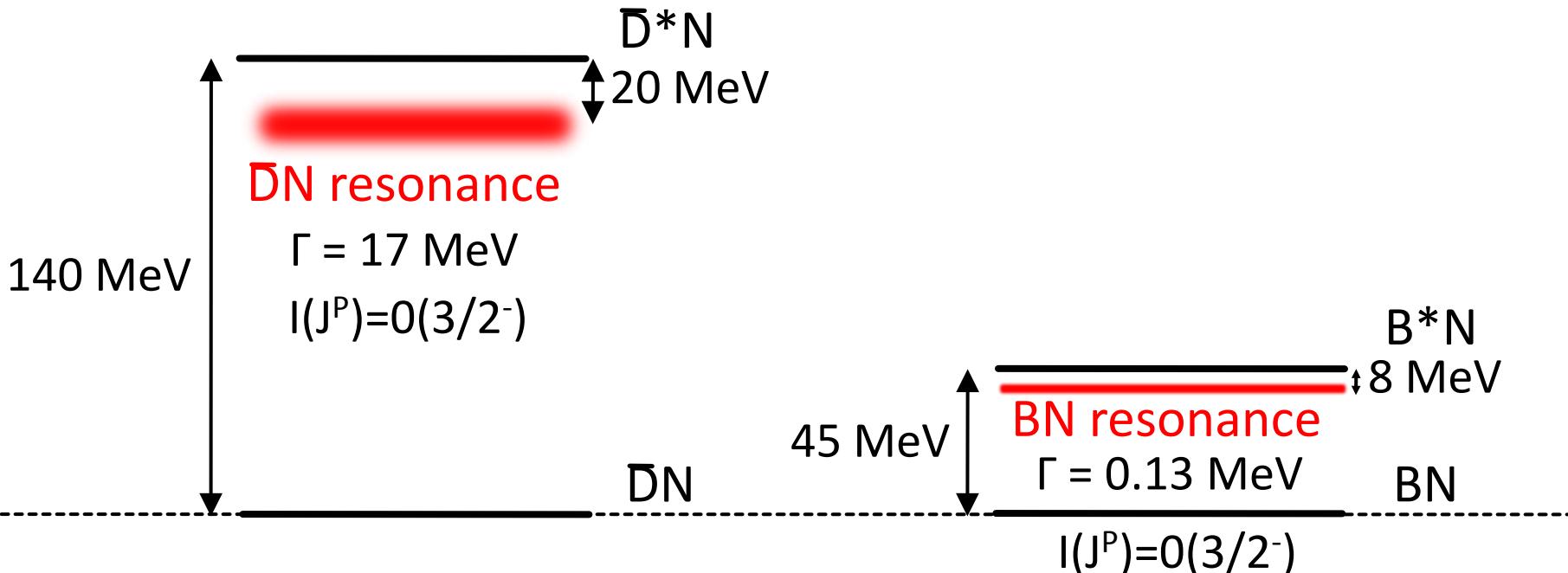
3. New exotic hadrons and nuclei

D and nucleon



3. New exotic hadrons and nuclei

D and nucleon

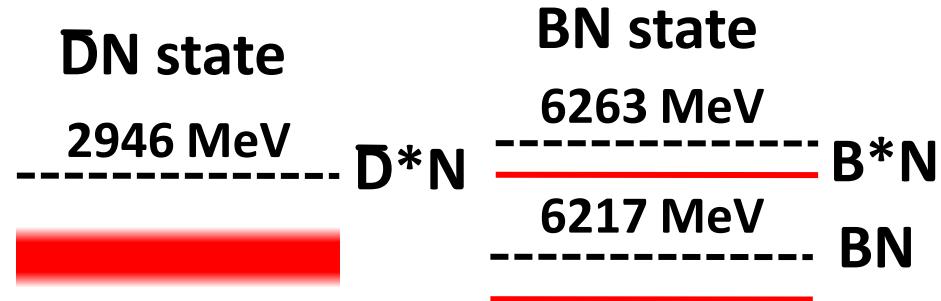


$I(J^P) = 0(1/2^-)$ bound state

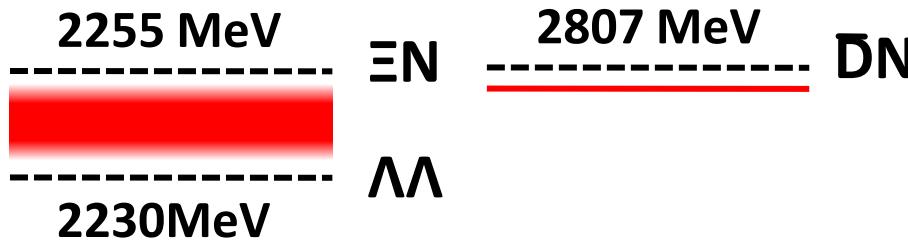
$I(J^P) = 0(3/2^-)$ narrow resonance state

heavy quark mass dependence is very interesting!

3. New exotic hadrons and nuclei



H dibaryon



$\Lambda(1405)$

1433 MeV

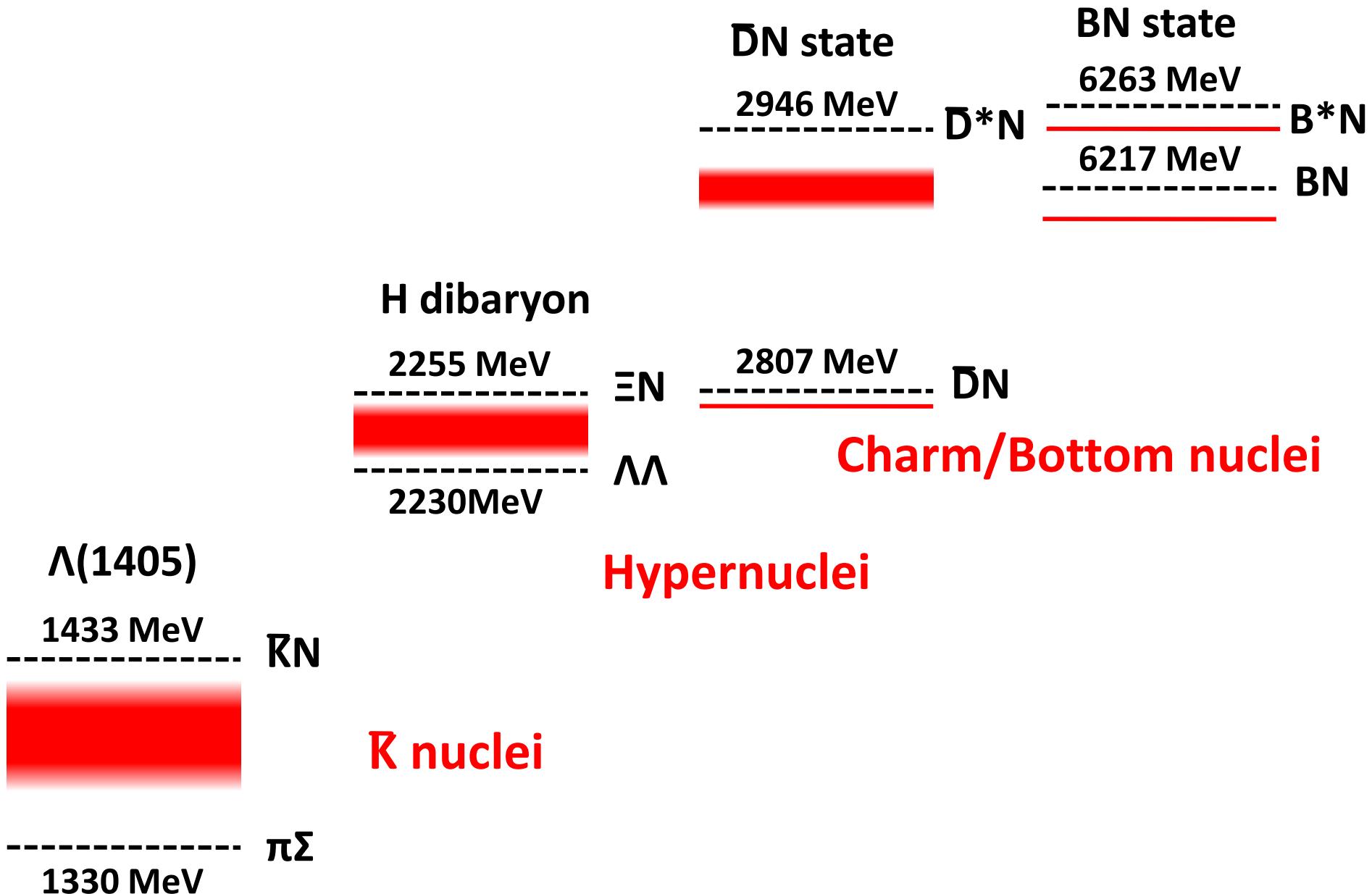
RN



$\pi \Sigma$

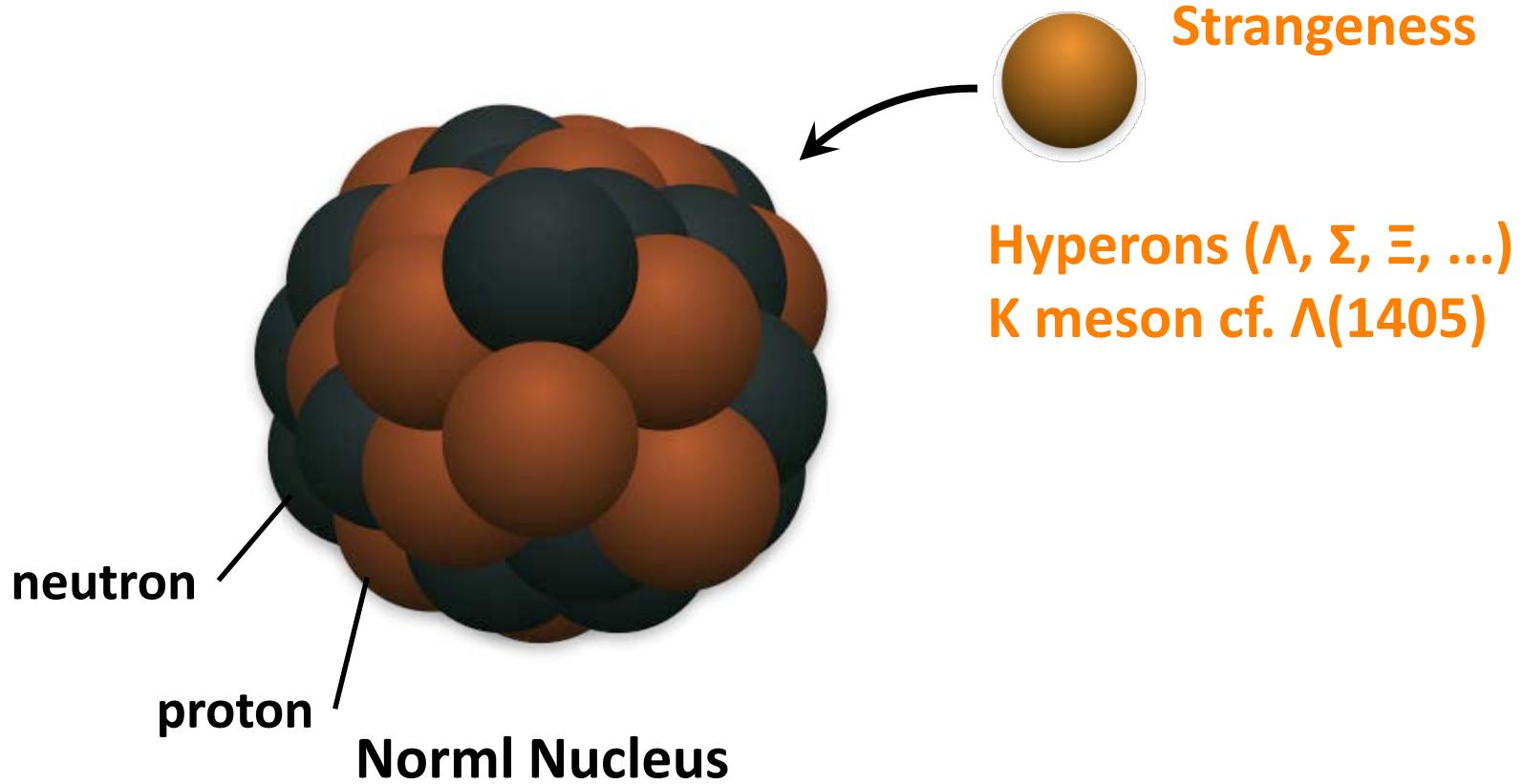
1330 MeV

3. New exotic hadrons and nuclei



3. New exotic hadrons and nuclei

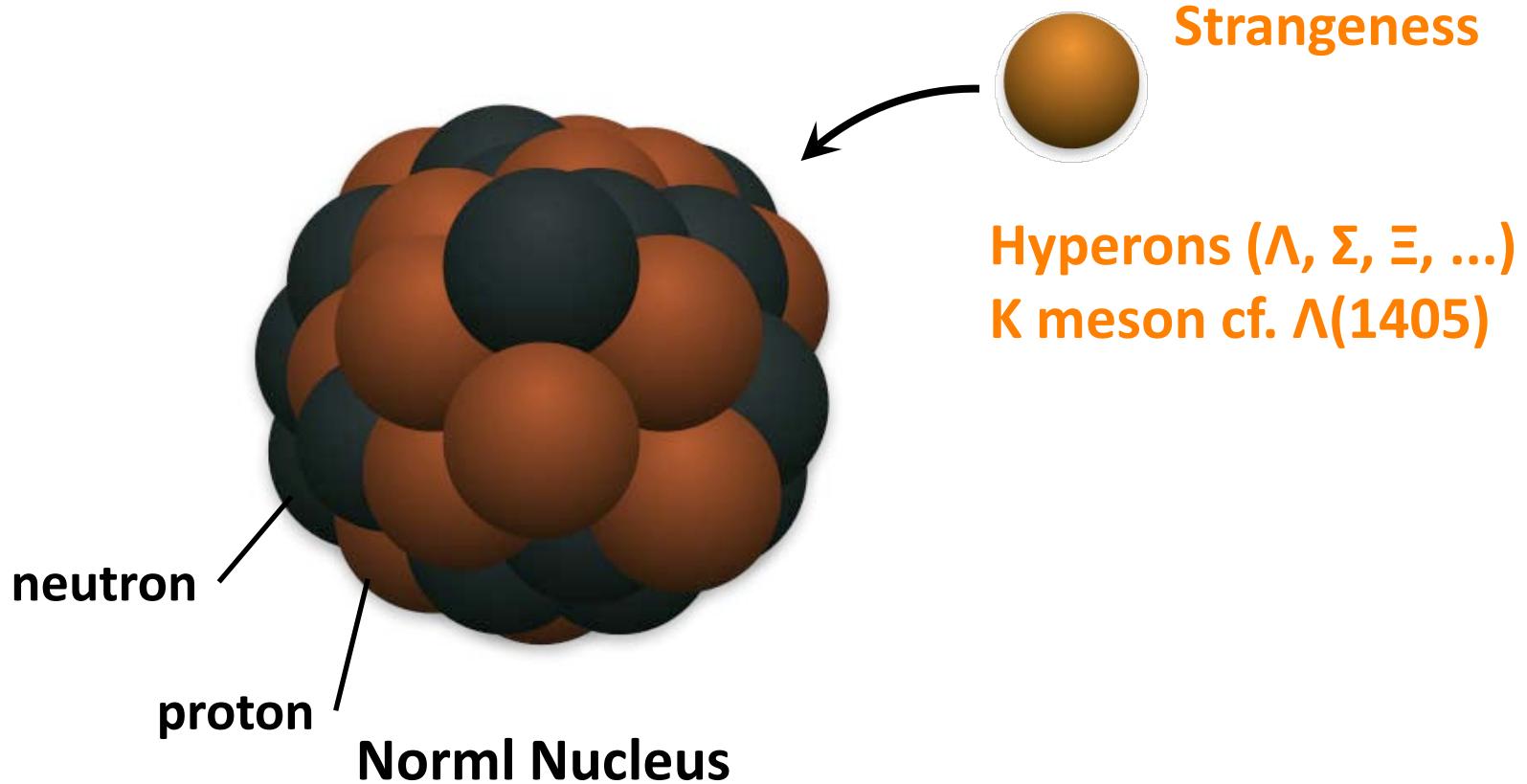
$| \text{Strange Nuclei} \rangle = U | \text{Normal Nuclei} \rangle$



$\exists U : \text{unitary transformation } SU(3)_f \text{ for some rep.}$

3. New exotic hadrons and nuclei

$| \text{Strange Nuclei} \rangle = U | \text{Normal Nuclei} \rangle$



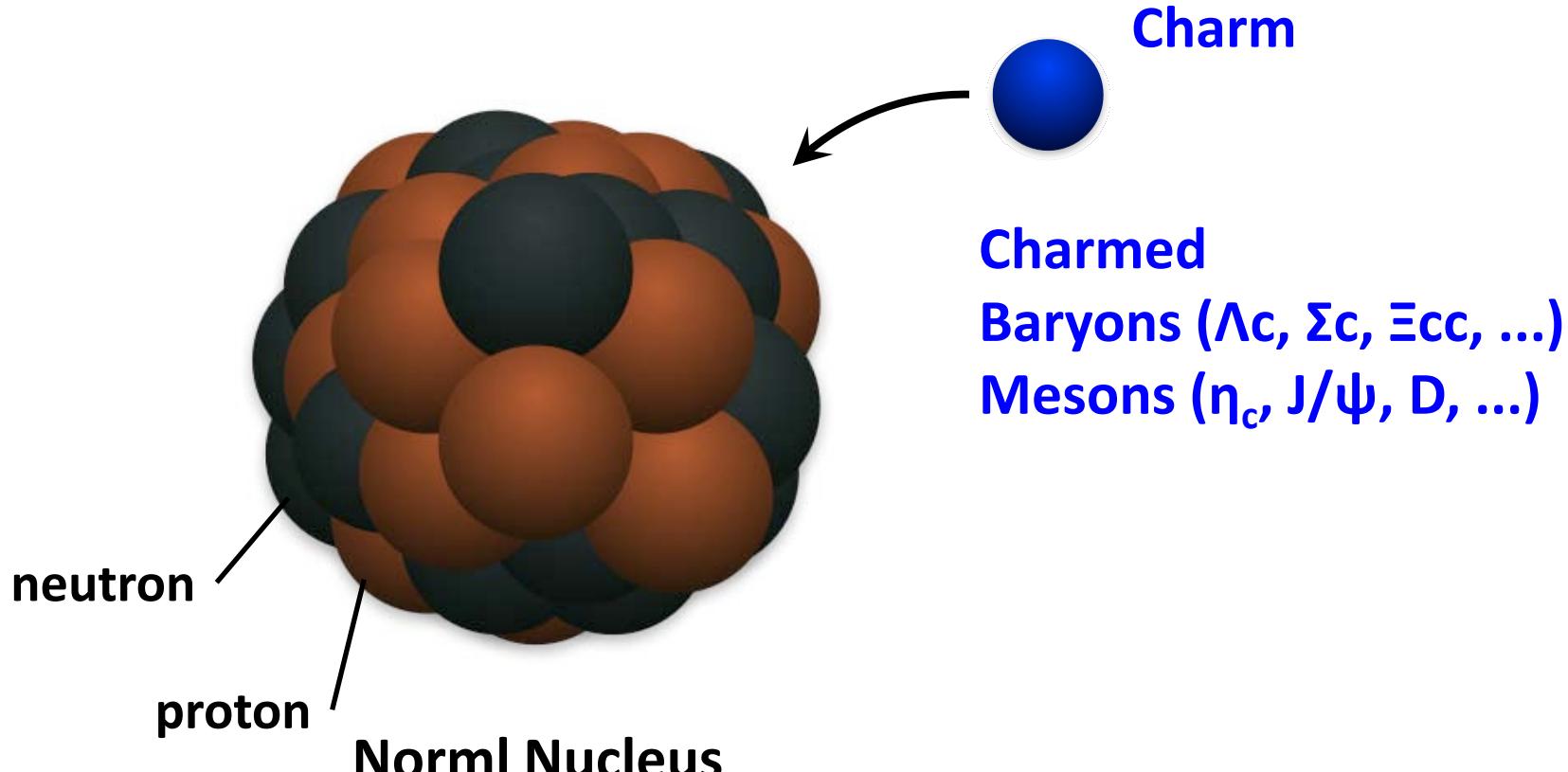
$\exists U : \text{unitary transformation } SU(3)_f \text{ for some representation}$

BB system $8 \times 8 = 1 + 8_s + 8_a + 10 + 10^* + 27$

Inoue et al.
PTP124, 591 (2010)

3. New exotic hadrons and nuclei

$| \text{Charmed Nuclei} \rangle \neq | \text{Normal Nuclei} \rangle$



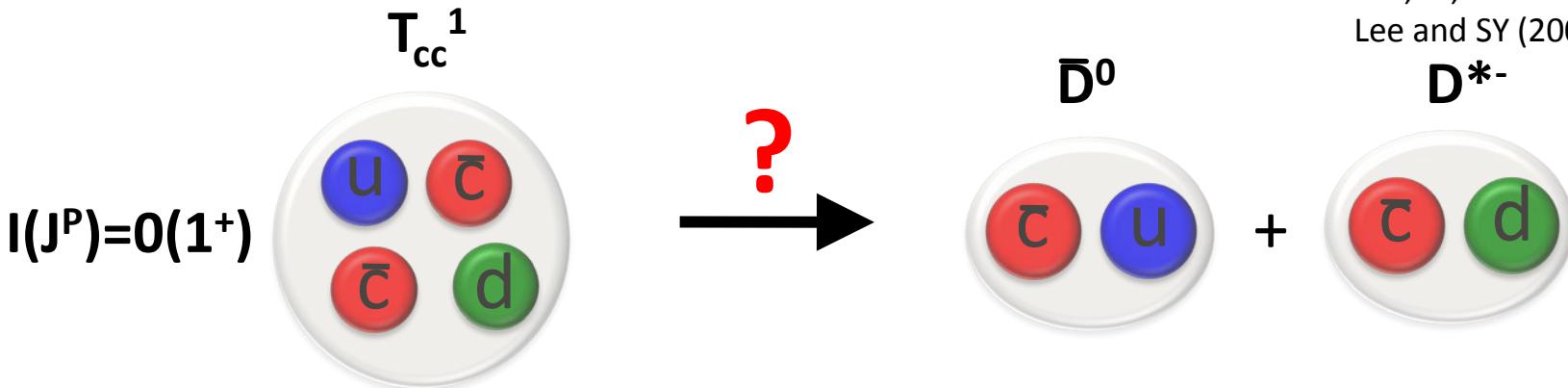
No unitary transformation

3. New exotic hadrons and nuclei

tetraquark T_{cc}^1 „explicitly“ exotic

Carlson, Heller, Tjon (1988)
Silvestre-Brac and Semay (1993)
Manohar and Wise (1993)
Lee, SY, Liu and Ko (2008)
Lee and SY (2009)

D^{*-}



Spin-color int. $\vec{s} \cdot \vec{B}^a$ ($a=1,..,8$) induces

$$H_{int} = \sum_{i>j} \frac{C_H}{m_i m_j} \vec{s}_i \cdot \vec{s}_j \quad C_H = v_0 \vec{\lambda}_i \cdot \vec{\lambda}_j \langle \delta(r_{ij}) \rangle$$

$\bar{c}c$ pair $1/m_Q^2$ ← further suppressed

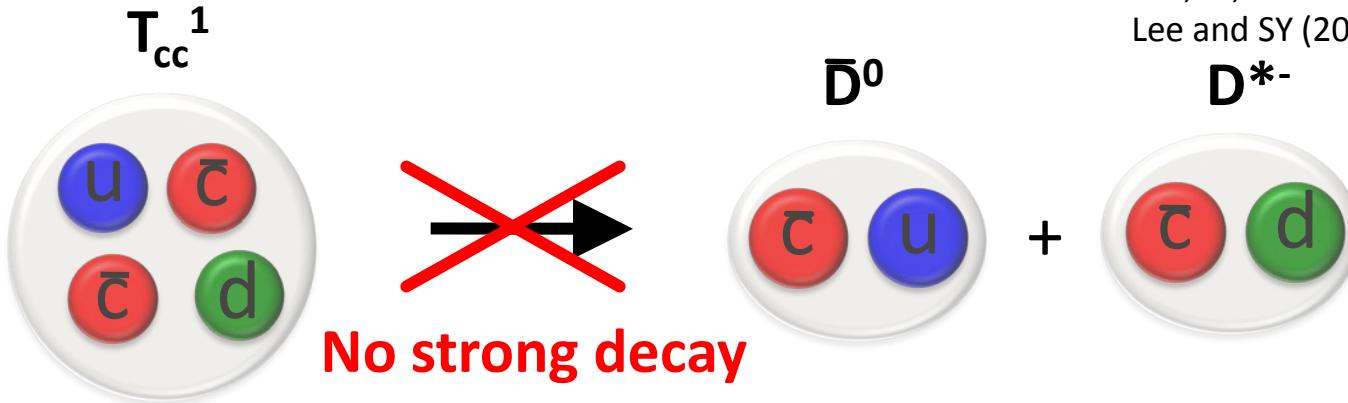
$\bar{c}u$ pair $1/m_Q^1$ ← suppressed

ud pair $1/m_Q^0$ ← dominant attraction ($\bar{3}_c, \bar{3}_f, s=0$)

3. New exotic hadrons and nuclei

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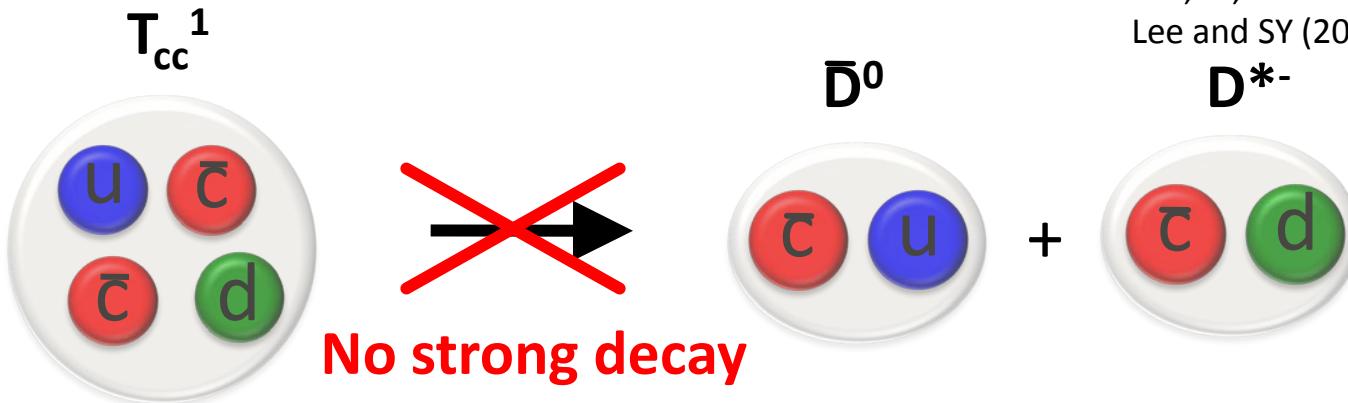
$$H_{int} = \sum_{i>j} \frac{C_H}{m_i m_j} \vec{s}_i \cdot \vec{s}_j \quad C_H = v_0 \vec{\lambda}_i \cdot \vec{\lambda}_j \langle \delta(r_{ij}) \rangle$$

$\bar{c}c$ pair	$1/m_Q^2$	\leftarrow	further suppressed
$\bar{c}u$ pair	$1/m_Q^1$	\leftarrow	suppressed
ud pair	$1/m_Q^0$	\leftarrow	dominant attraction ($\bar{3}_c, \bar{3}_f, s=0$)

3. New exotic hadrons and nuclei

tetraquark T_{cc}^1 „explicitly“ exotic

Carlson, Heller, Tjon (1988)
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Lee and SY (2009)



Binding energy of $T_{cc(bb)}^1$ [MeV]

	$ud\bar{c}\bar{c}$	$us\bar{c}\bar{c}$	$ds\bar{c}\bar{c}$
T_{cc}^1	-74.9	-4.3	-4.3
	$\bar{D}^0 + D^{*-}, \bar{D}^{*0} + D^-$	$\bar{D}^0 + D_s^{*-}$	$D^- + D_s^{*-}$
	$ud\bar{b}\bar{b}$	$us\bar{b}\bar{b}$	$ds\bar{b}\bar{b}$
T_{bb}^1	-123.8	-61.4	-61.4
	$B^+ + B^{*0}, B^{*+} + B^0$	$B^+ + B_s^{*0}$	$B^0 + B_s^{*0}$

→ $T_{cc(bb)}^1$ are stable as $\bar{3}_f$ multiplet of $SU(3)_f$.

3. New exotic hadrons and nuclei

\bar{D} and nucleon ✓

\bar{D}



e⁺e⁻ collisions
pp (pp^{bar}) collisions
heavy ion collisions

raquark T_{cc}¹

✓

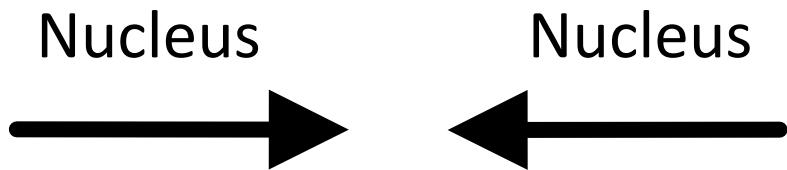


4. Experimental researches

Exotic hadrons from heavy ion collisions at RHIC and LHC

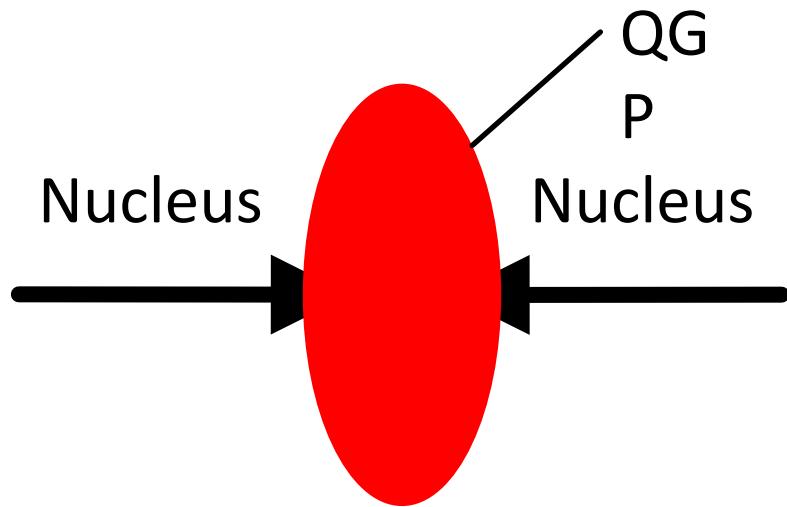
4. Experimental researches

Exotic hadrons from heavy ion collisions at RHIC and LHC



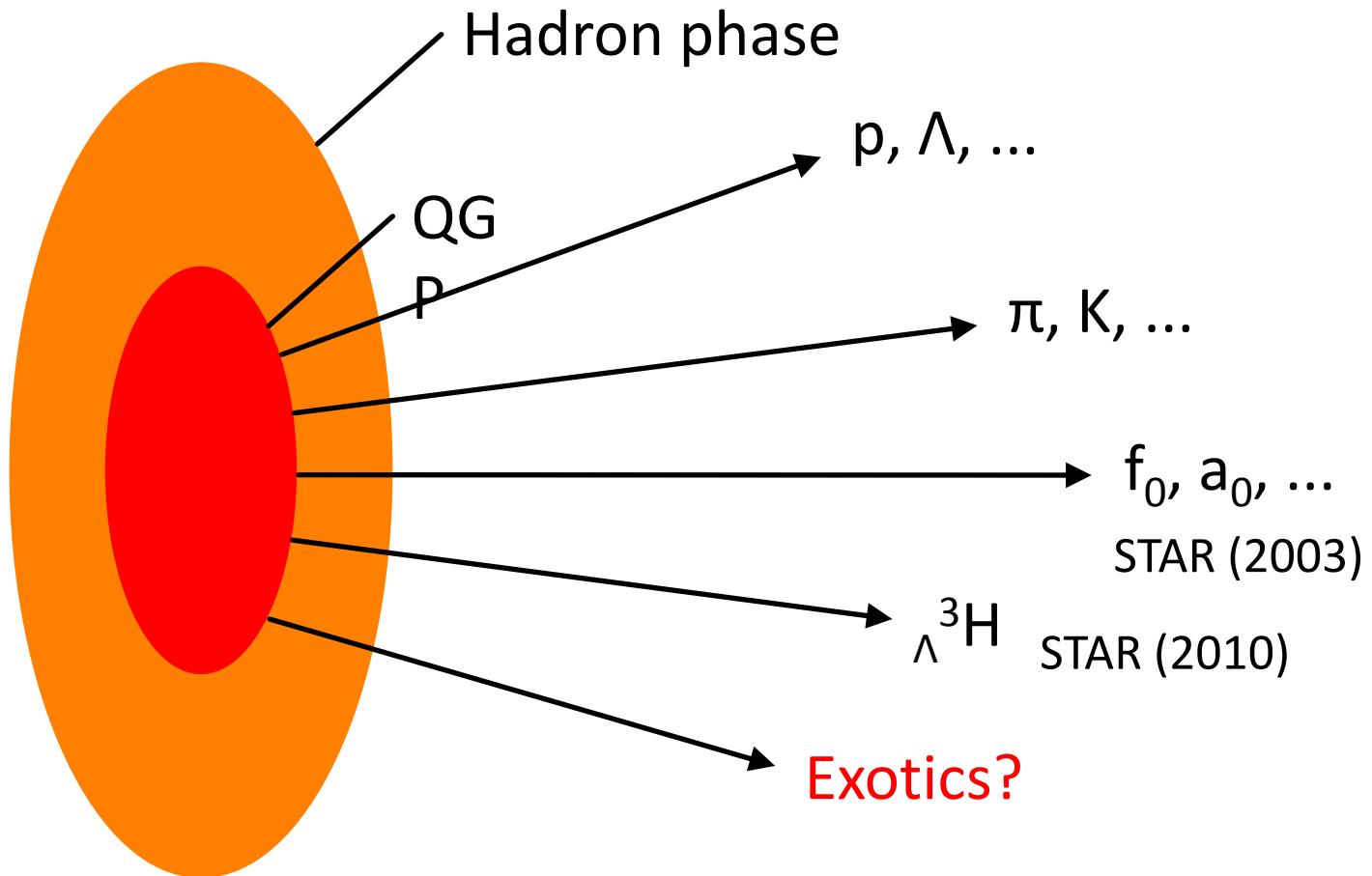
4. Experimental researches

Exotic hadrons from heavy ion collisions at RHIC and LHC



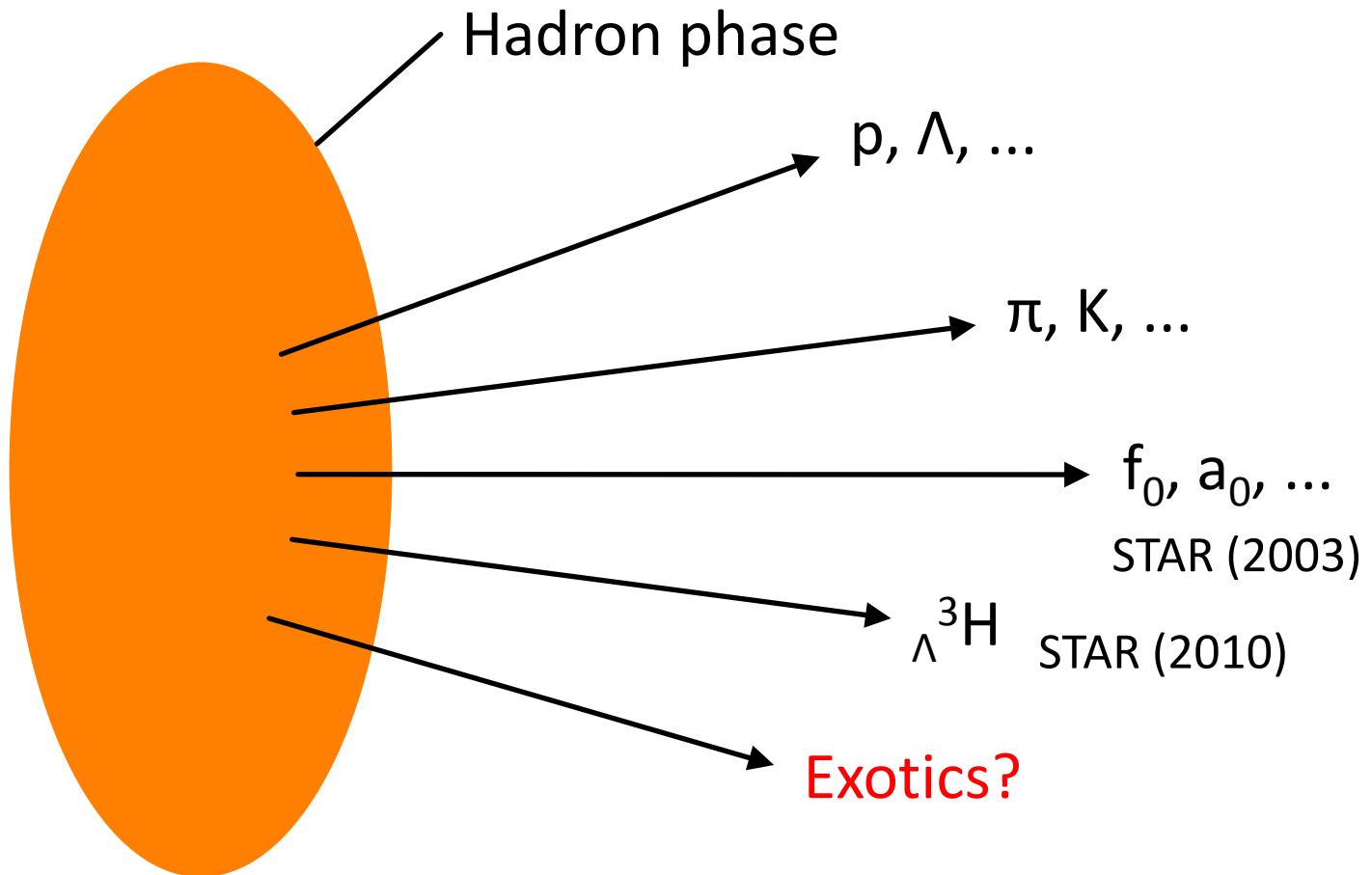
4. Experimental researches

Exotic hadrons from heavy ion collisions at RHIC and LHC



4. Experimental researches

Exotic hadrons from heavy ion collisions at RHIC and LHC



4. Experimental researches

Yields of exotic hadrons at RHIC and LHC

S. Cho *et al.* (the ExHIC collaboration), PRL106, 212001 (2011)

	RHIC				LHC			
	2q/3q/6q	4q/5q/8q	Mol.	Stat.	2q/3q/6q	4q/5q/8q	Mol.	Stat.
$f_0(980)$	3.8, 0.73($s\bar{s}$)	0.10	13	5.6	10, 2.0 ($s\bar{s}$)	0.28	36	15
$a_0(980)$	11	0.31	40	17	31	0.83	1.1×10^2	46
$D_s(2317)$	1.3×10^{-2}	2.1×10^{-3}	1.6×10^{-2}	5.6×10^{-2}	8.7×10^{-2}	1.4×10^{-2}	0.10	0.35
$X(3872)$	—	4.0×10^{-5}	7.8×10^{-4}	2.9×10^{-4}	—	6.6×10^{-4}	1.3×10^{-2}	4.7×10^{-3}
$\Lambda(1405)$	0.81	0.11	1.8–8.3	1.7	2.2	0.29	4.7–21	4.2
$\bar{K}KN$	—	0.019	1.7	0.28	—	5.2×10^{-2}	4.2	0.67
$\bar{D}N$	—	2.9×10^{-3}	4.6×10^{-2}	1.0×10^{-2}	—	2.0×10^{-2}	0.28	6.1×10^{-2}
$\bar{K}NN$	5.0×10^{-3}	5.1×10^{-4}	0.011–0.24	1.6×10^{-2}	1.3×10^{-2}	1.4×10^{-3}	0.026–0.54	3.7×10^{-2}
$\bar{D}NN$	—	2.9×10^{-5}	1.8×10^{-3}	7.9×10^{-5}	—	2.0×10^{-4}	9.8×10^{-3}	4.2×10^{-4}

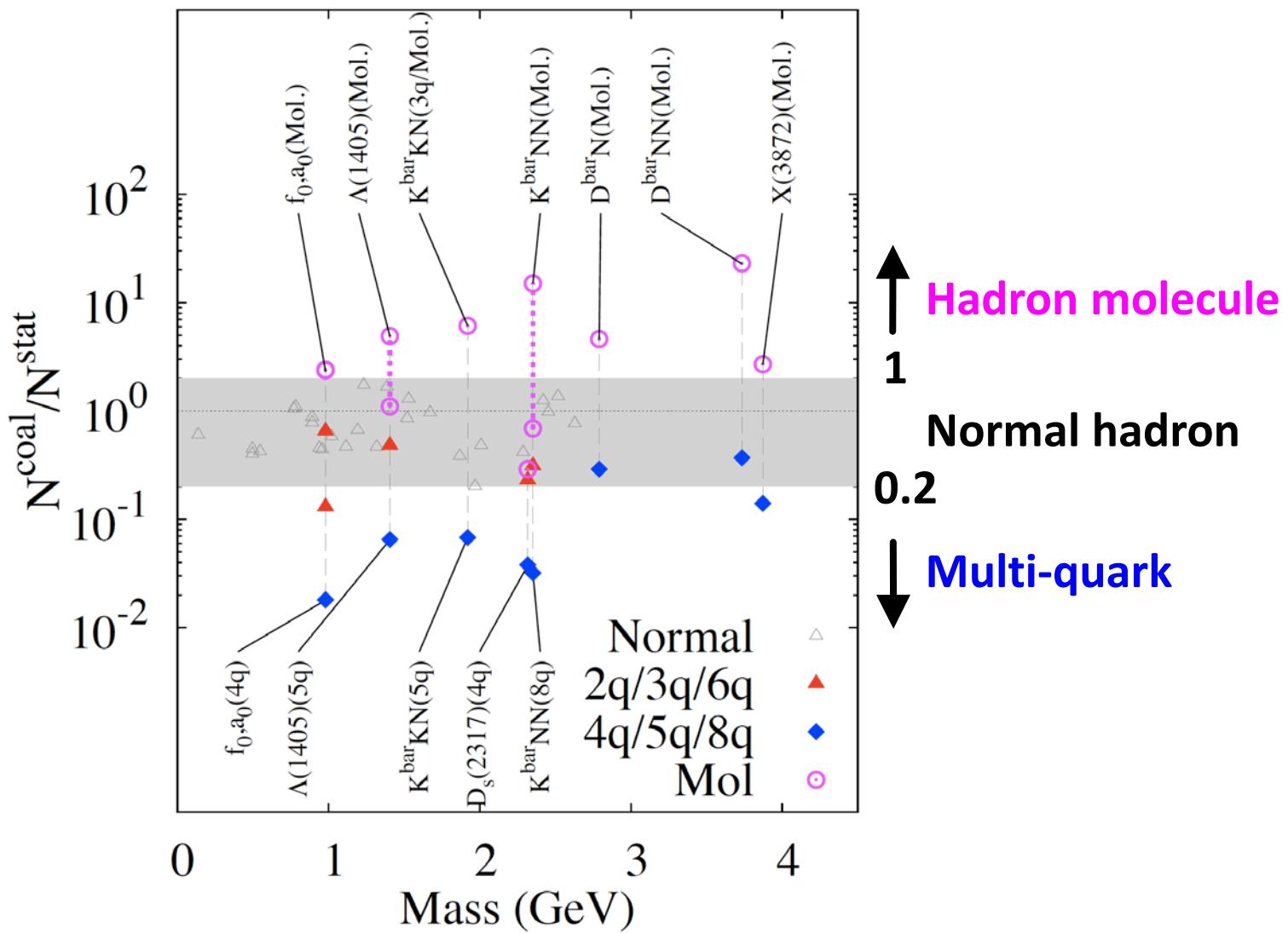
ExHIC (Exotics from Heavy Ion Collisions) collaboration

S. Cho (Yonsei U.), T. Furumoto (YITP and RIKEN), T. Hyodo (Tokyo Inst. Technology), D. Jido (YITP), C.-M. Ko (Texas A&M U.), S. H. Lee (Yonsei U. and YITP), M. Nielsen (U. de Paulo), A. Ohnishi (YITP), T. Sekihara (YITP and Kyoto U.), K. Yazaki (YITP and RIKEN), S. Y. (KEK)

4. Experimental researches

Hadron molecule vs. multi-quark

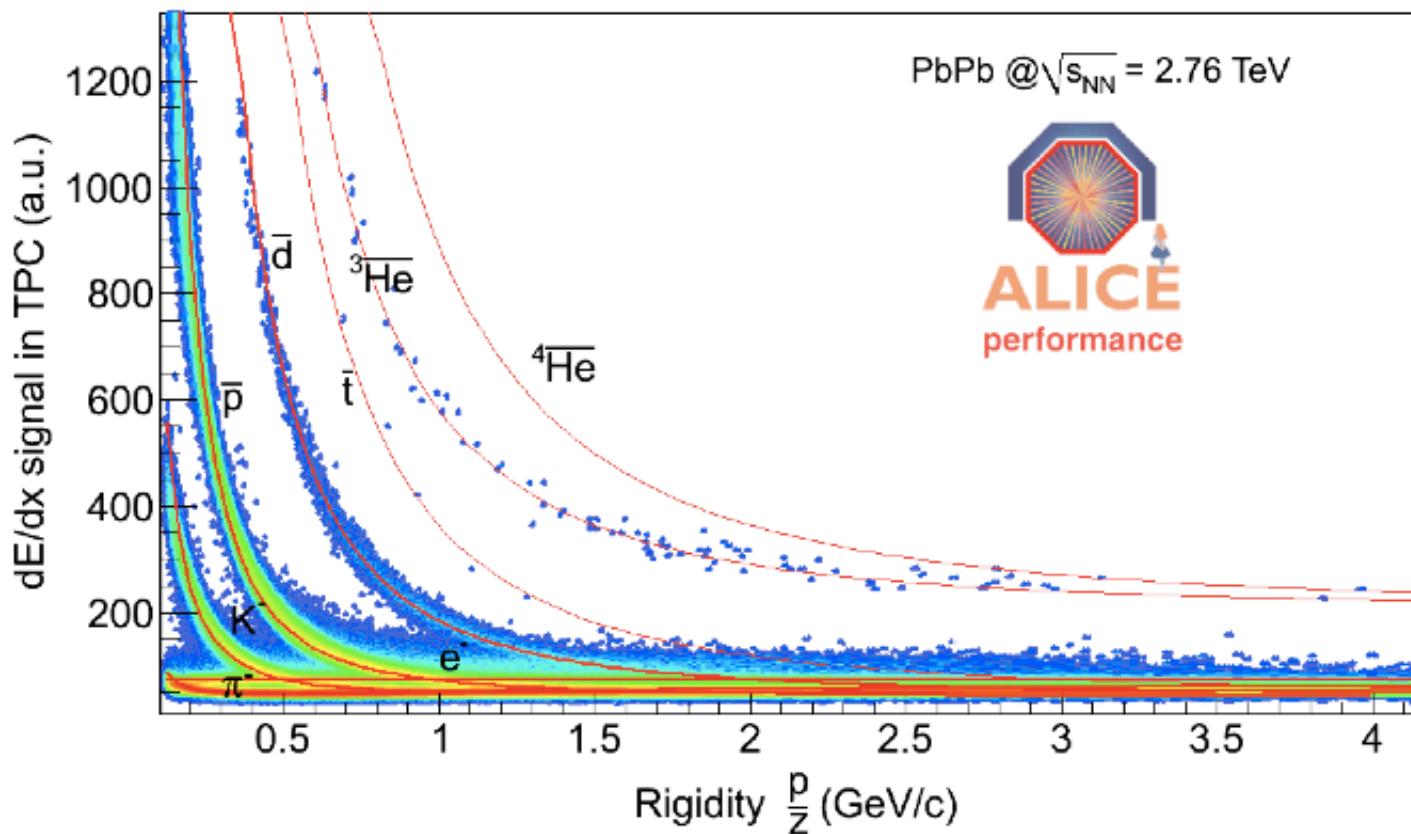
Coal. / Stat. ratio at RHIC



4. Experimental researches



Anti-Nuclei

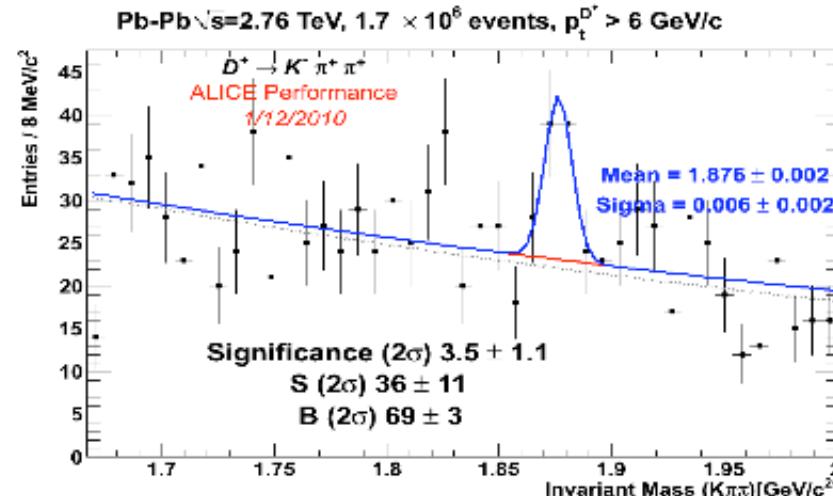
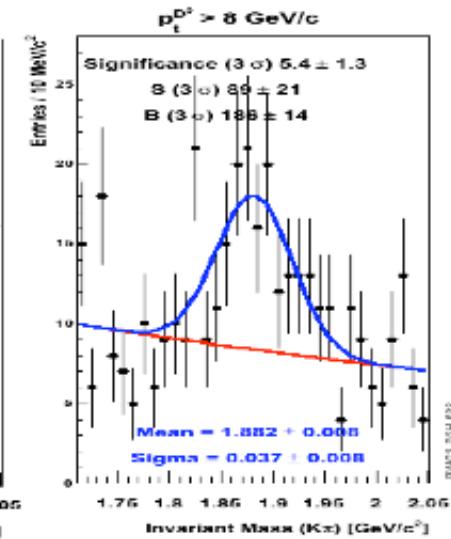
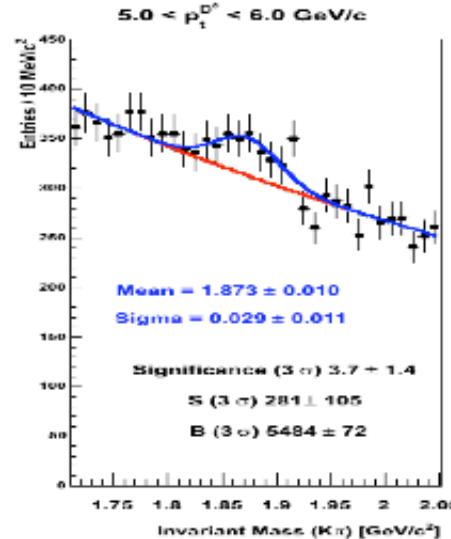
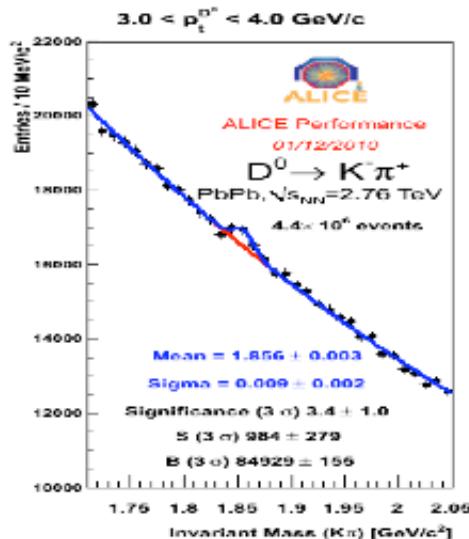


~ 2 M Pb-Pb Min Bias events

4. Experimental researches



Charm in Pb-Pb



'Jet quenching' with heavy quarks:
Energy loss depends on
- color charge (quark/gluon)
- mass (light/heavy quarks)

5. Summary

Heavy quarks have new symmetry and dynamics.

Heavy quark symmetry

D-D* mixing

...

Many exotic heavy hadrons and nuclei.

T_{cc}^1

Z_b ($I=1$)

$D^{\bar{b}}$ N bound and resonance states

...

Searches in e^+e^- , pp ($pp^{\bar{b}}$) and heavy ion collisions.