

Numerical study on quantum entanglement entropy in 4d $SU(3)$ gauge theories

Etsuko Ito (KEK)

collaboration with K. Nagata (KEK), Y.Nakagawa, A. Nakamura
(Hiroshima U., RCNP) and V.I.Zakharov (Max Planck Inst.)

cf. arXiv:0911.2596 and 1104.1011:

Y.Nakagawa, A.Nakamura, S.Motoki and V.I.Zakharov



entanglement entropy

=?

a novel approach

to understand the confinement

Don't think....feel....

Outline

- Introduction (QCD in 4d)
- Definition of entanglement entropy
- Replica method
- Results for the quenched QCD

QCD

dynamics of gluons and quarks

- Yang-Mills theory
(SU(3) gauge theory)

$$\mathcal{L} = -\frac{1}{4} (F_{\mu\nu}^a)^2 + \bar{\psi} (iD_\mu \gamma^\mu - m) \psi$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + ig f^{abc} A_\mu^b A_\nu^c$$

$$D_\mu = \partial_\mu - ig A_\mu^a t^a$$

local symmetry

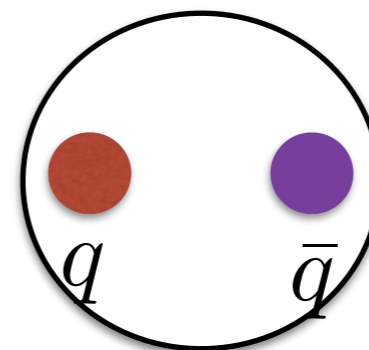
- parameter: gauge coupling, fermion mass

- adjoint reps. of SU(3)
(gluon: messenger of the force)

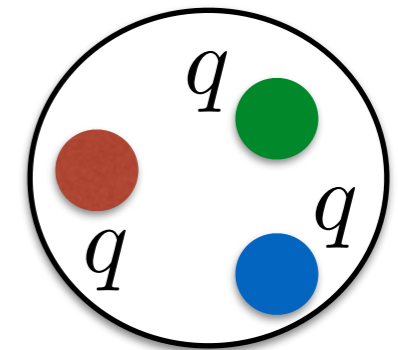
$$A_\mu^a(x) \quad \begin{cases} a = 1, \dots, 8 \\ \mu = 1, \dots, 4 \end{cases}$$

- fundamental reps. of SU(3)
(quark: fundamental element of matter)

$$\psi_\alpha^i(x), \bar{\psi}_\alpha^{\bar{i}} \quad \text{or } q, \bar{q} \quad \begin{cases} i = 1, 2, 3 \\ \alpha = 1, \dots, 4 \end{cases}$$



meson
ex) pion

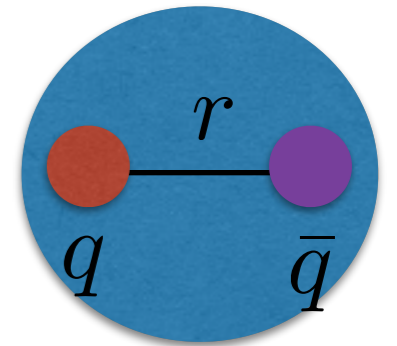


baryon
ex) proton

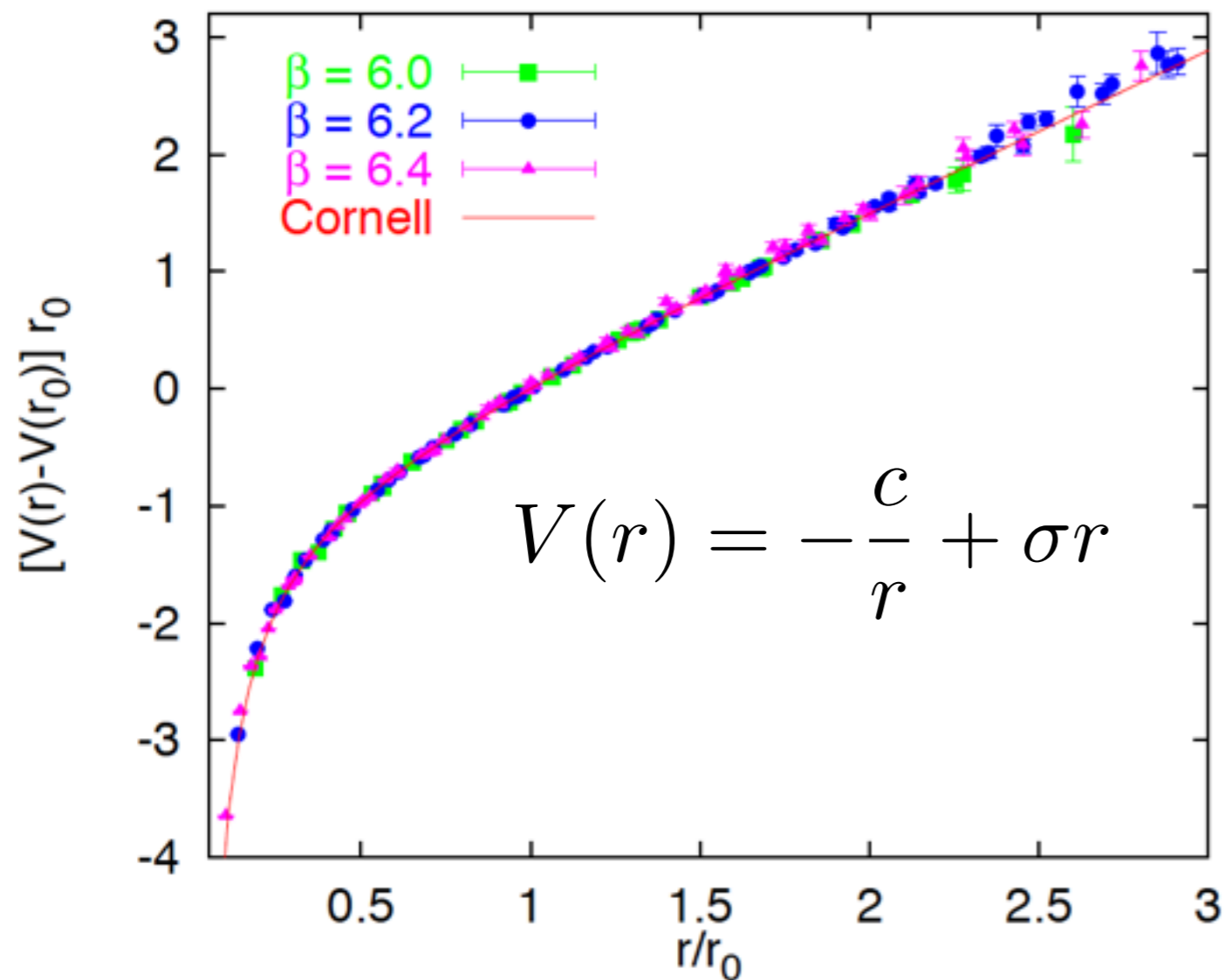
remarks: pure YM or quenched QCD => only first term
full QCD => whole lagrangian

confinement

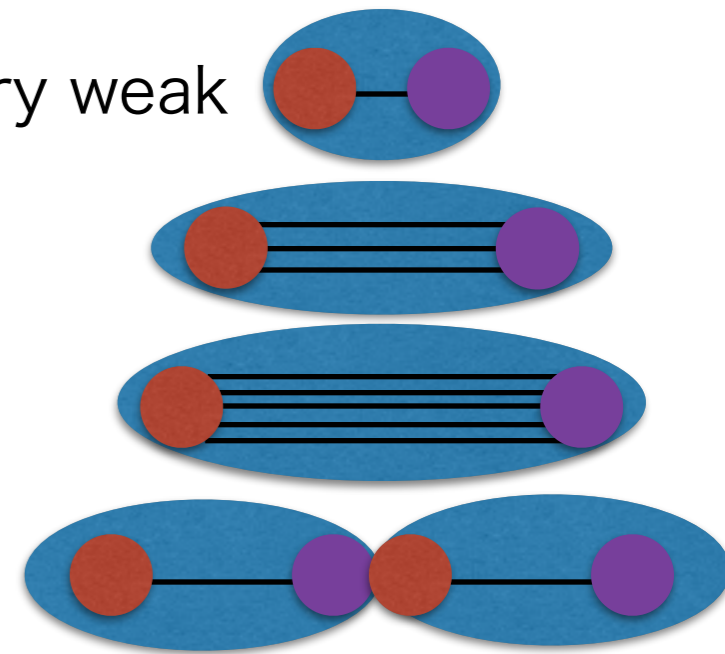
A nonperturbative property of QCD



Hadron



interaction is very weak



$$r > \Lambda_{\text{QCD}}^{-1}$$

- q-qbar potential
- center symmetry (only for pure YM)

Basic properties of E.E.

entanglement entropy for quantum system

- how much a quantum state is entangled quantum mechanically
- d.o.f of the system
- quantum properties of the ground state for the system
- in finite T system, it gives a thermal entropy

QCD theory (T=0)

A color confinement changes the d.o.f of the system

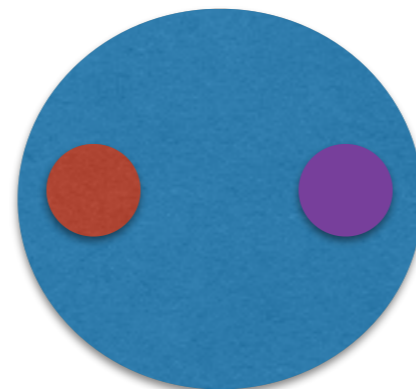
microscopically

Λ_{QCD}

macroscopically

colorful
(gluons)

$\sim O(N_c^2)$



colorless
(singlet)

$\sim O(1)$

Definition of the entanglement entropy

Entanglement entropy (E.E.)

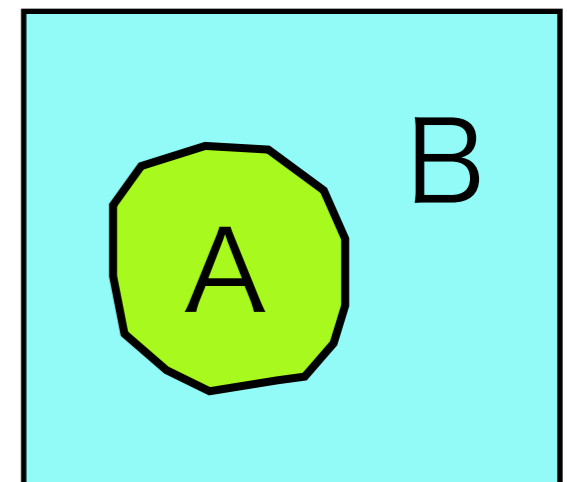
von Neumann entropy $S_{tot} = -\text{Tr} \rho_{tot} \log \rho_{tot}$

density matrix $\rho_{tot} = |\Psi\rangle\langle\Psi|$ $|\Psi\rangle$: pure ground state

decompose total Hilbert space into two subsystems $\mathcal{H}_{tot} = \mathcal{H}_A \otimes \mathcal{H}_B$

reduced density matrix $\rho_A = -\text{Tr}_{\mathcal{H}_B} \rho_{tot}$

entanglement entropy $S_A = -\text{Tr}_A \rho_A \log \rho_A$



At finite T, it is equivalent to the thermal entropy.

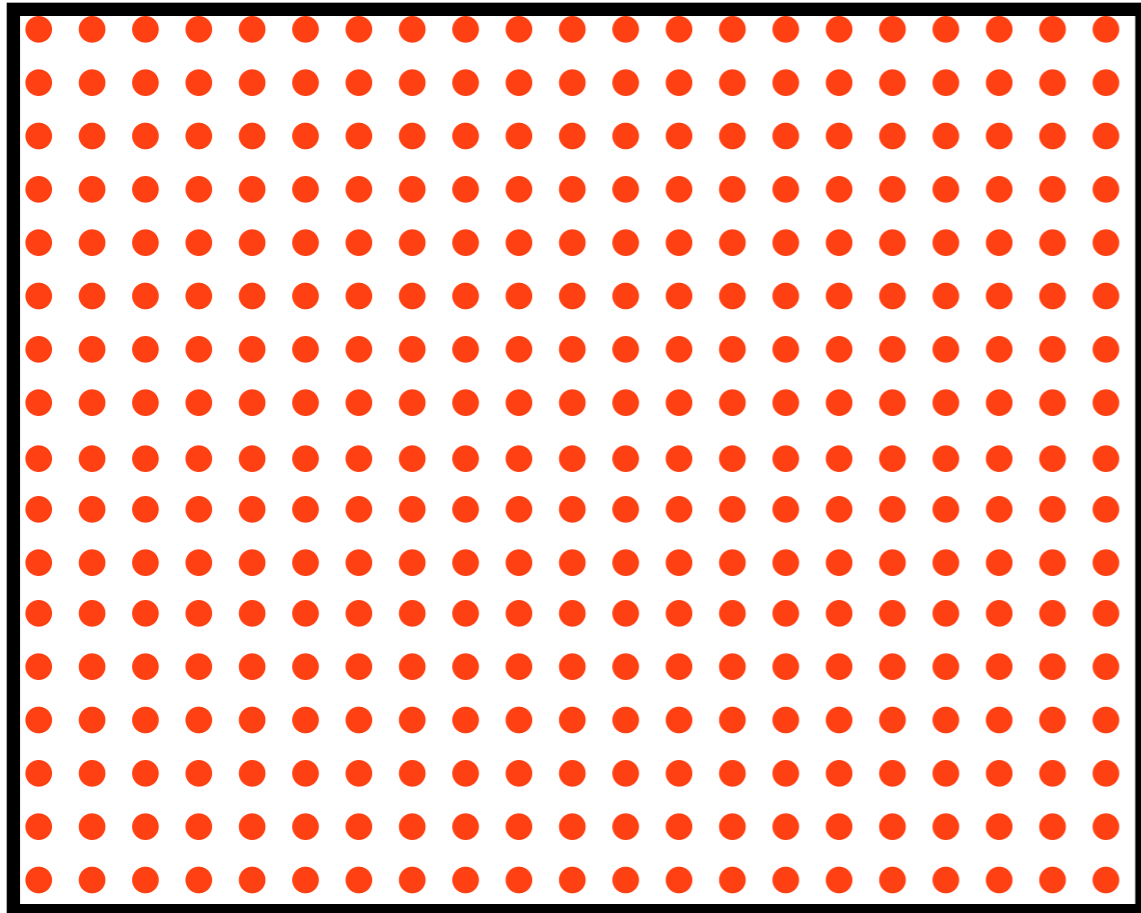
Schematic picture of E.E.

$(1+1)$ -dim. model

[Holzhey,Larsen and Wilczek: NPB424 \(1994\) 443](#)

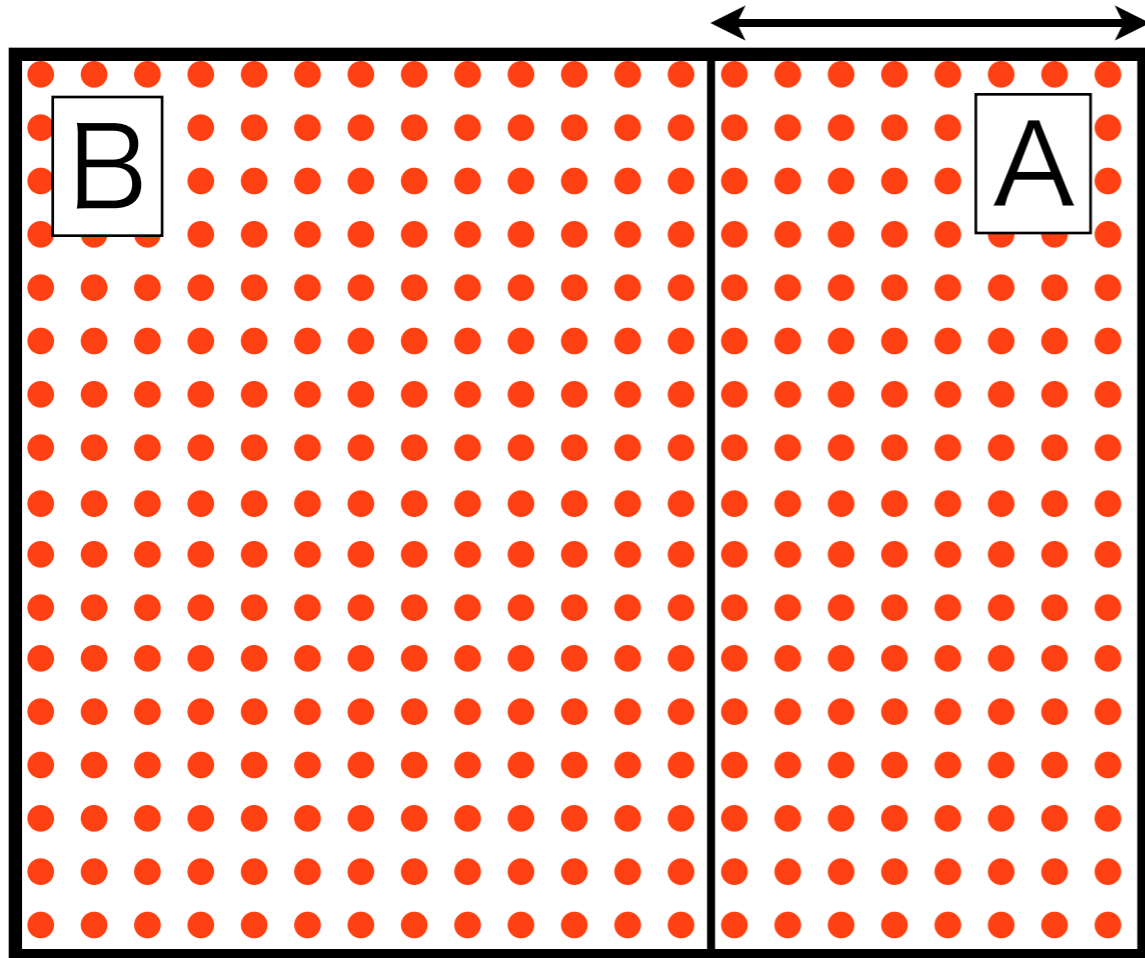
[Calabrese and Cardy: J.S.M.0406\(2004\)P06002](#)

[Calabrese and Cardy, arXiv:0905.4013](#)



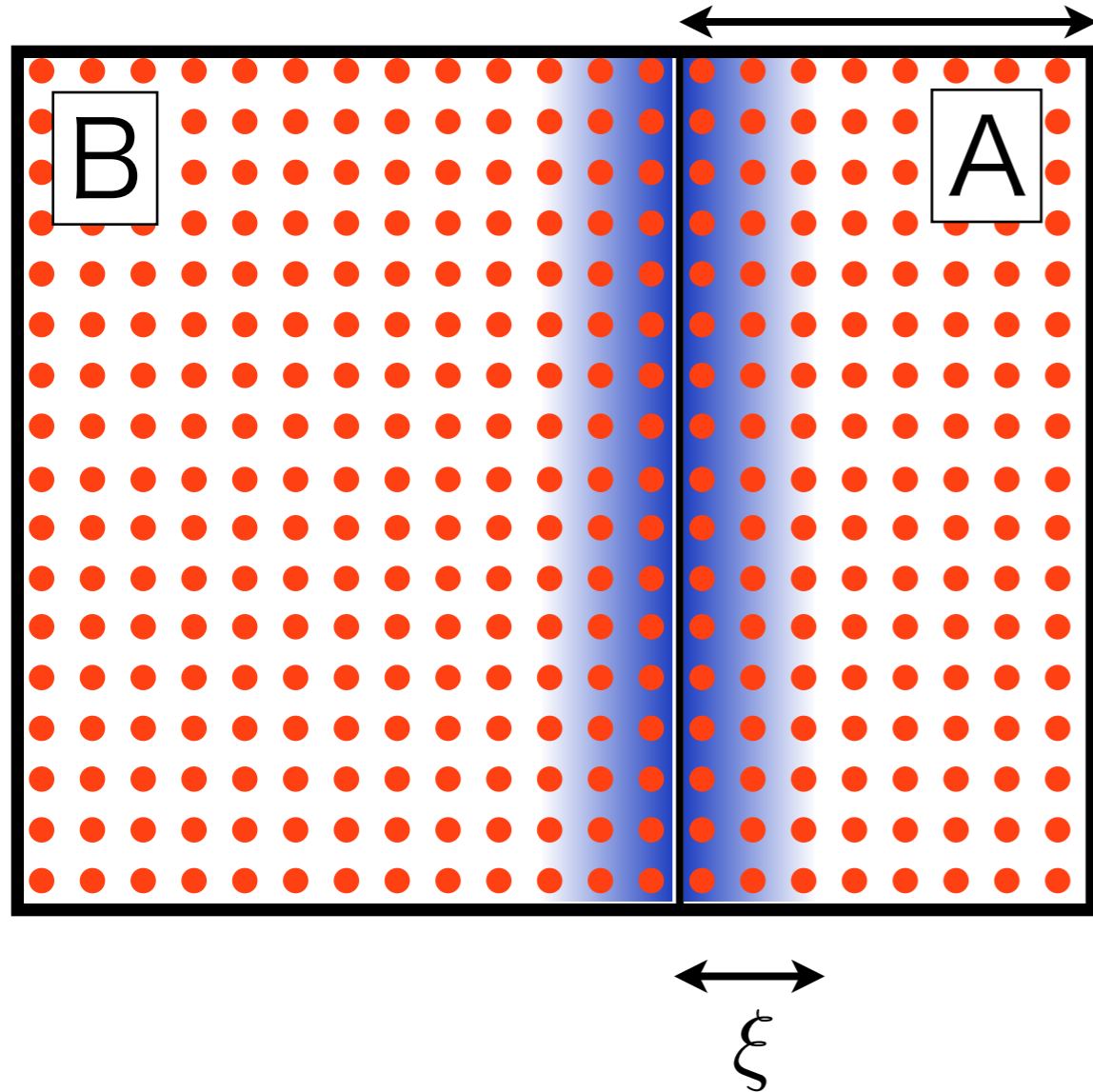
Schematic picture of E.E.

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Schematic picture of E.E.

(1+1)-dim. model



At the critical point,

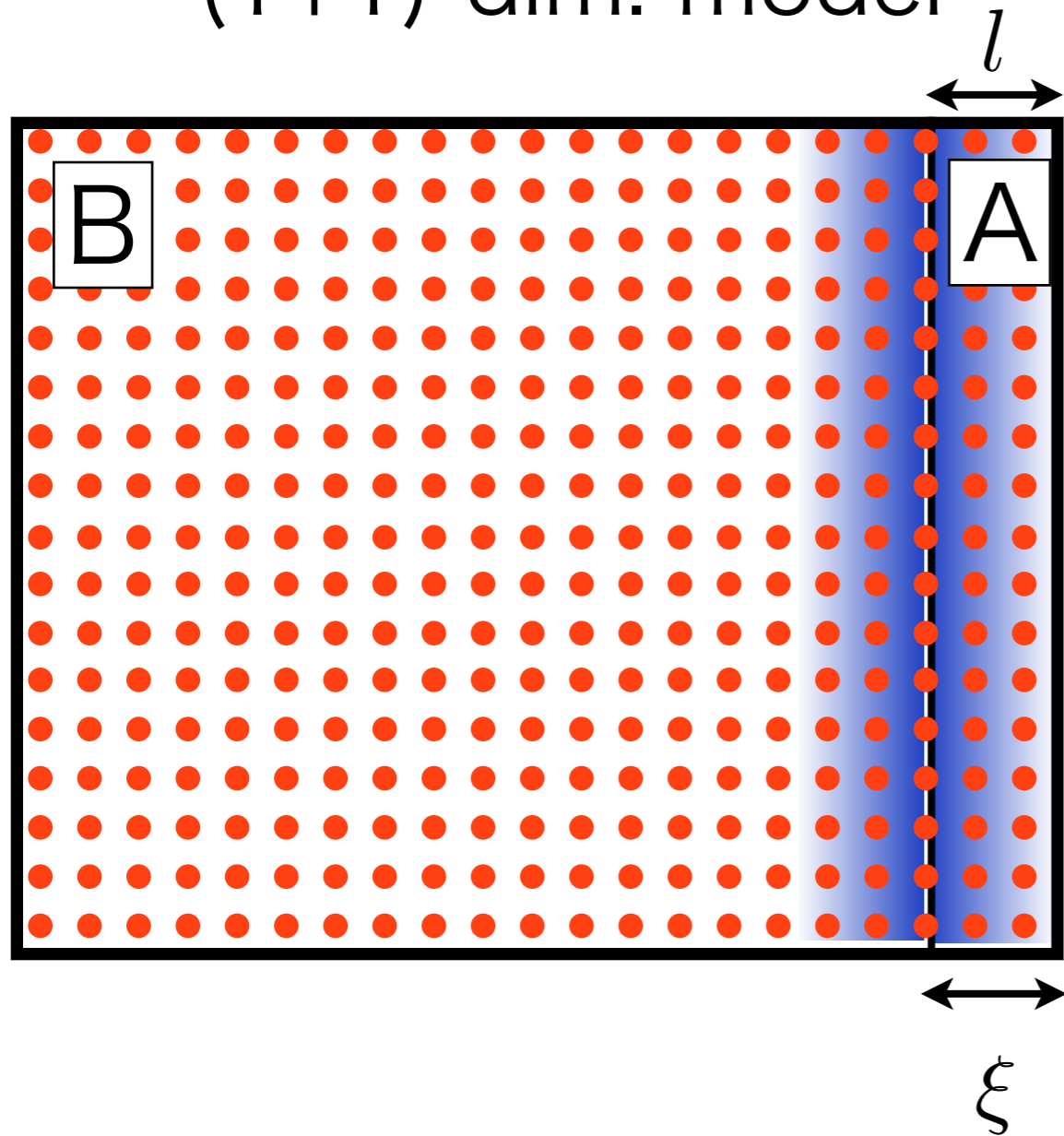
$$S_A(l) = \frac{c}{3} \log \frac{l}{a} + c_1$$

c is the central charge
in 2d CFT.

ξ : correlation length of the system

Schematic picture of E.E.

(1+1)-dim. model



In the non critical system,

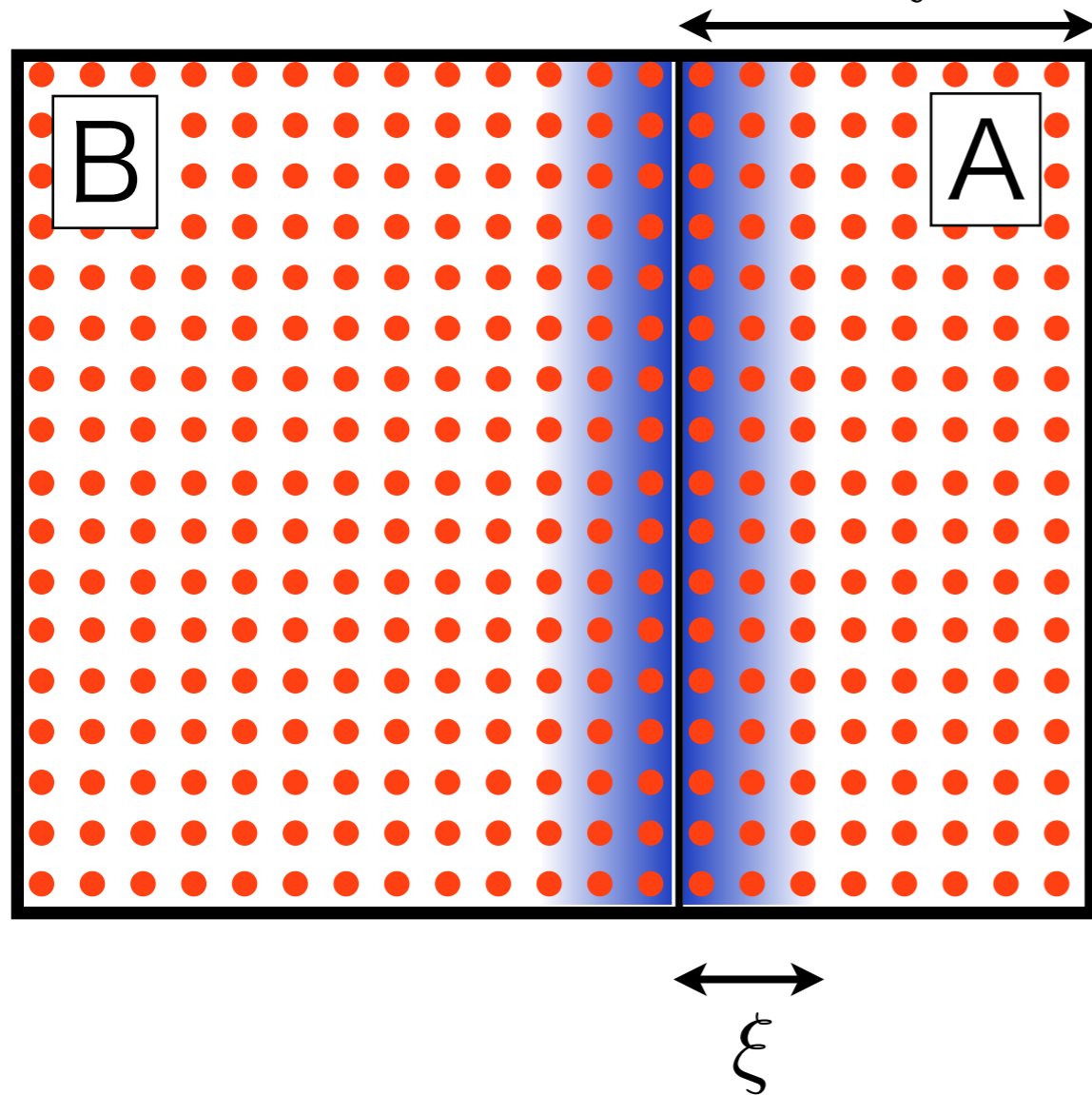
$$l \ll \xi$$

$$S_A(l) \sim \frac{c}{3} \log \frac{l}{a}$$

ξ : correlation length of the system

Schematic picture of E.E.

(1+1)-dim. model



In the non critical case,

$$l \ll \xi$$

$$S_A(l) \sim \frac{c}{3} \log \frac{l}{a}$$

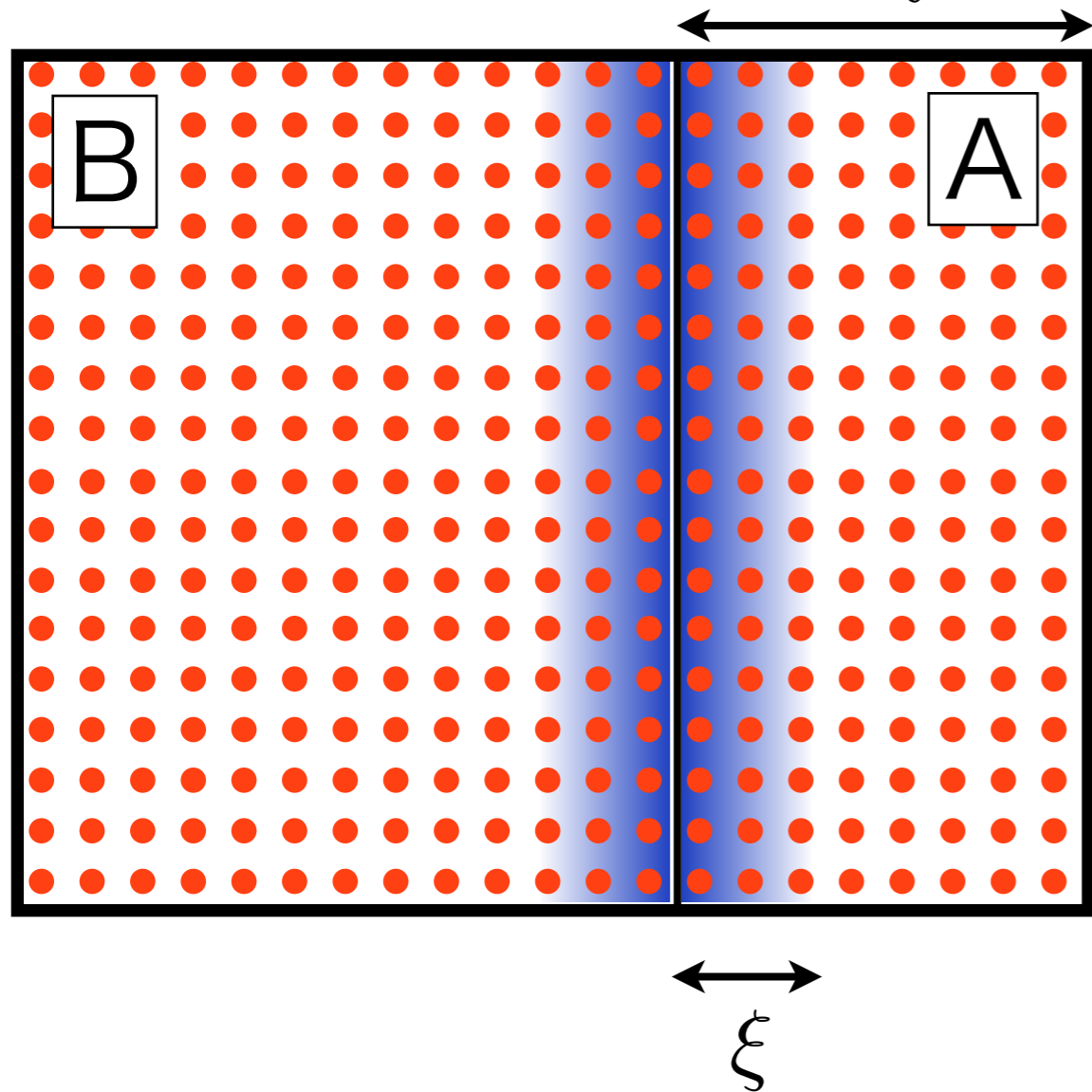
$$l \gg \xi$$

$$S_A(l) \rightarrow \frac{c}{3} \log \frac{\xi}{a}$$

ξ : correlation length of the system

Schematic picture of E.E.

(1+1)-dim. model



ξ : correlation length of the system

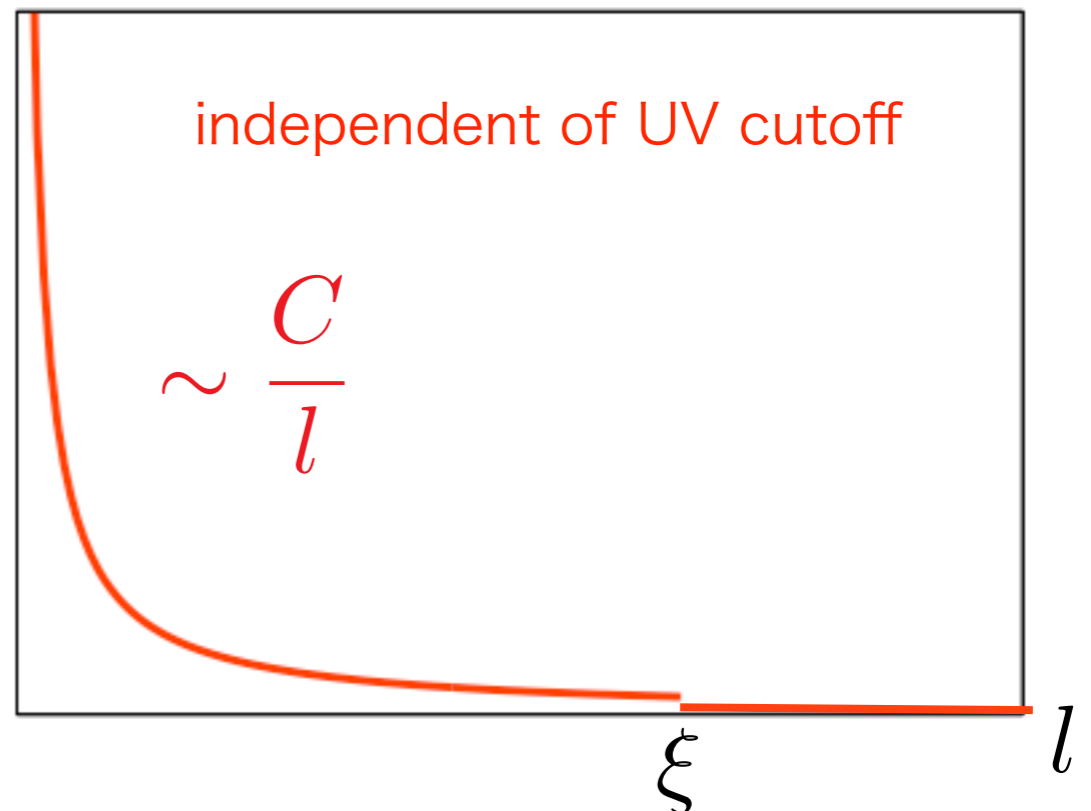
At the critical point or $l \ll \xi$ in the noncritical system

$$S_A(l) = \frac{c}{3} \log \frac{l}{a} + c_1$$

In the non-critical system,

$$S_A(l) \xrightarrow{l \gg \xi} \frac{c}{3} \log \frac{\xi}{a}$$

$$\frac{\partial S_A(l)}{\partial l}$$



Difficulties to obtain E.E. in 4d gauge theory

- UV cutoff dependence of 4d E.E.

2d

$$S_A(l) = c_0 \log(l/a)$$

4d

$$S_A(l) = c_0 \frac{\text{Area}}{a^2} - c'_0 \frac{\text{Area}}{l^2} + c_1 \log(l/a) + (\text{regular terms})$$

Entropic C-function

$$C(l) = l^3 \frac{1}{|\partial A|} \frac{dS}{dl}$$

Ryu and Takayanagi:PRL96(2006)181602
JHEP 0608(2006)045

cf) Holographic approach

C is obtained by AdS and QFT

- (local) gauge invariance

E.E. for gauge theory

- P.V.Buividovich and M.I.Polikarpov PLB670(2008)141

extended Hilbert space

- H.Casini, M.Muerta and J.A.Rosabal arXiv:1312.1183

electric b.c.(electric center), magnetic center, trivial center

- D.Radicevic arXiv:1404.1391

magnetic center

- W.Donnelly PRD85 (2012) 085004

extended lattice construction

- S.Ghosh, R.M.Soni,S.P.Trivedi arXiv:1501.02593

- S.Aoki, T.Iritani, M.Nozaki et.al. arXiv:1502.04267

maximally gauge invariant reduced density matrix

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maximally gauge invariant reduced density matrix

red definitions are inadequate
for E.E. or ρ_A

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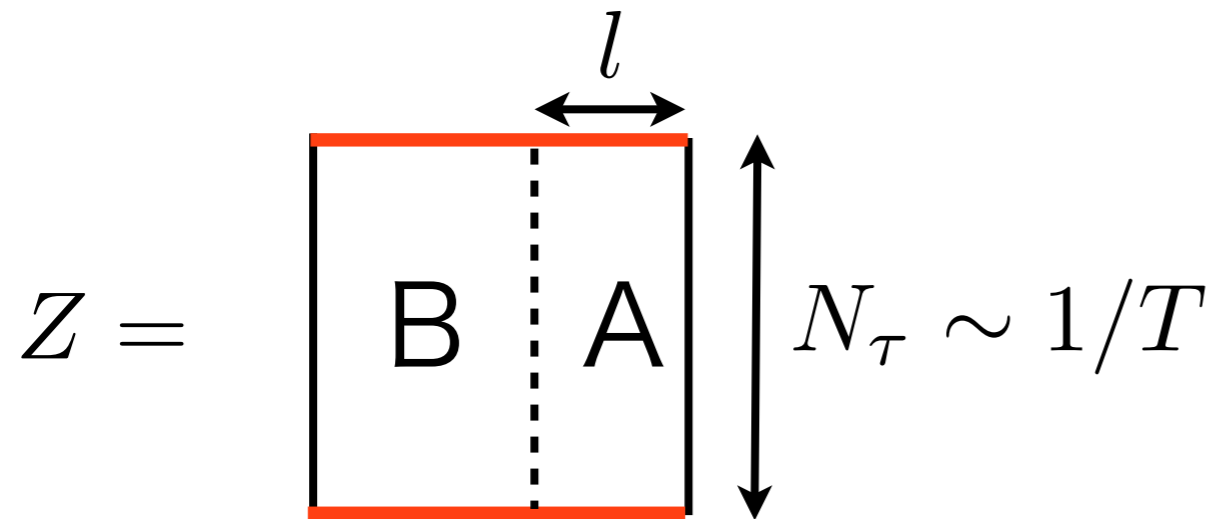
maximally gauge invariant reduced density matrix

red definitions are inadequate
for E.E. or ρ_A

Replica method

Calabrese and Cardy: J.S.M.0406(2004)P06002

Replica method



entanglement entropy:

$$S_A = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \ln \text{Tr}_A \rho_A^n$$

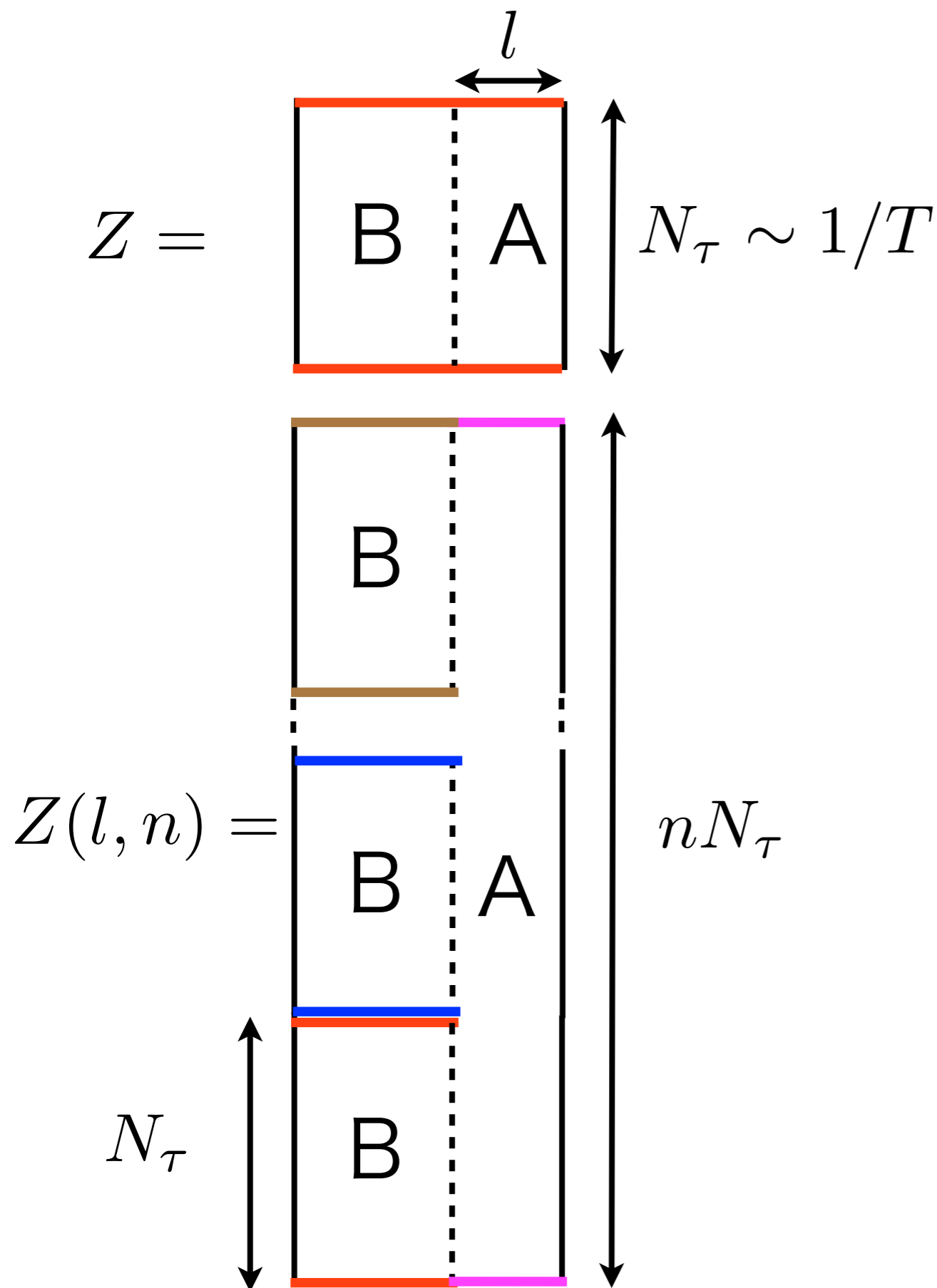
$$S_A(l) = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \ln \left(\frac{Z(l, n)}{Z^n} \right)$$

Replica method

entanglement entropy:

$$S_A = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \ln \text{Tr}_A \rho_A^n$$

$$S_A(l) = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \ln \left(\frac{Z(l, n)}{Z^n} \right)$$



Replica method

entanglement entropy:

$$S_A = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \ln \text{Tr}_A \rho_A^n$$

$$S_A(l) = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \ln \left(\frac{Z(l, n)}{Z^n} \right)$$

observable:

$$\frac{\partial S_A(l)}{\partial l} = \lim_{n \rightarrow 1} \frac{\partial}{\partial l} \frac{\partial}{\partial n} F[l, n] \quad \text{:free energy}$$

$$\rightarrow \frac{F[l + a, n = 2] - F[l, n = 2]}{a}$$

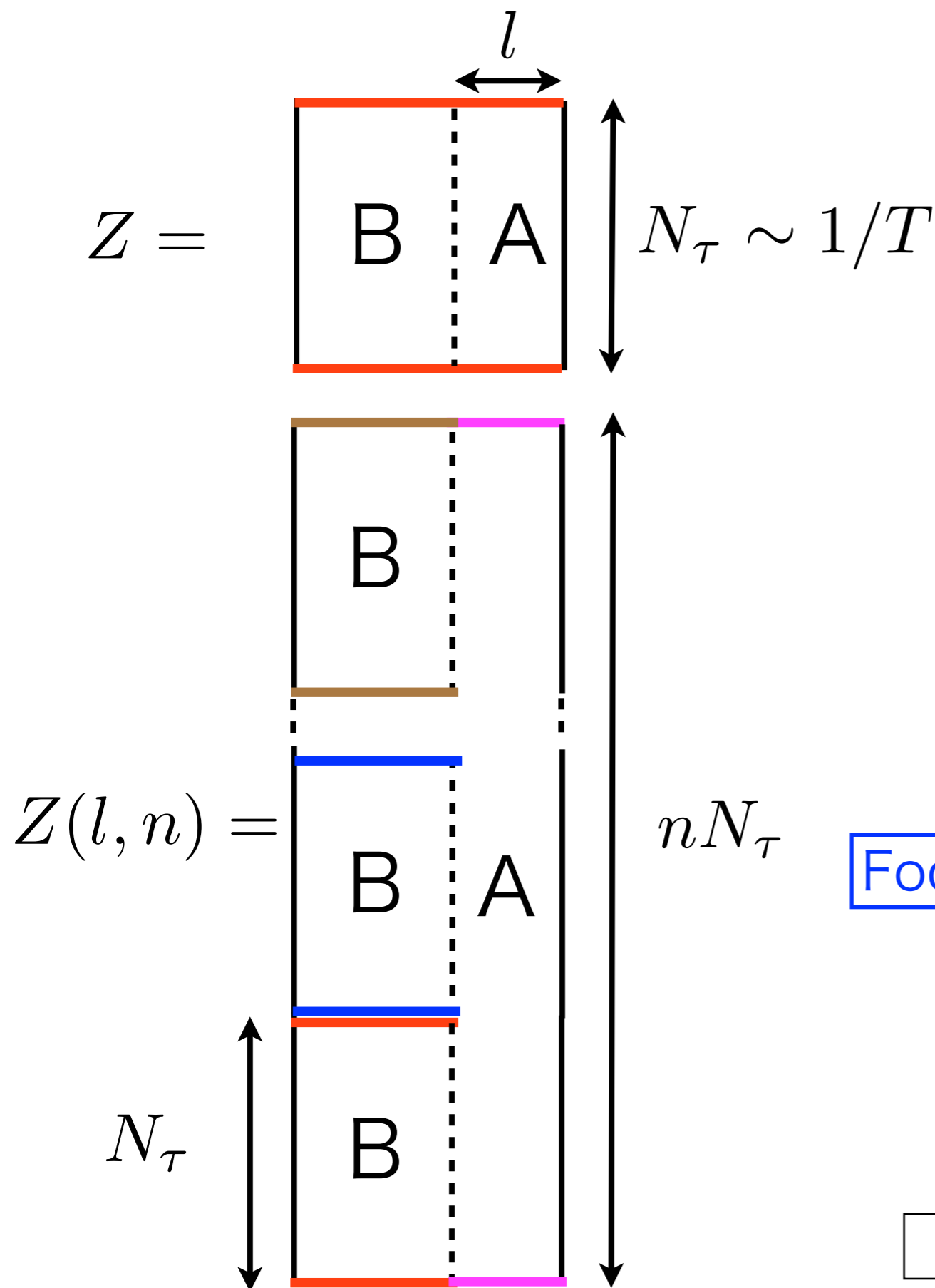
Fodor (2009)

$$= \int_0^1 d\alpha \langle S_{l+a}[U] - S_l[U] \rangle_\alpha$$

using the interpolation action

$$S_{int} = (1 - \alpha) S_l[U] + \alpha S_{l+a}[U]$$

we measure the diff. of the action density

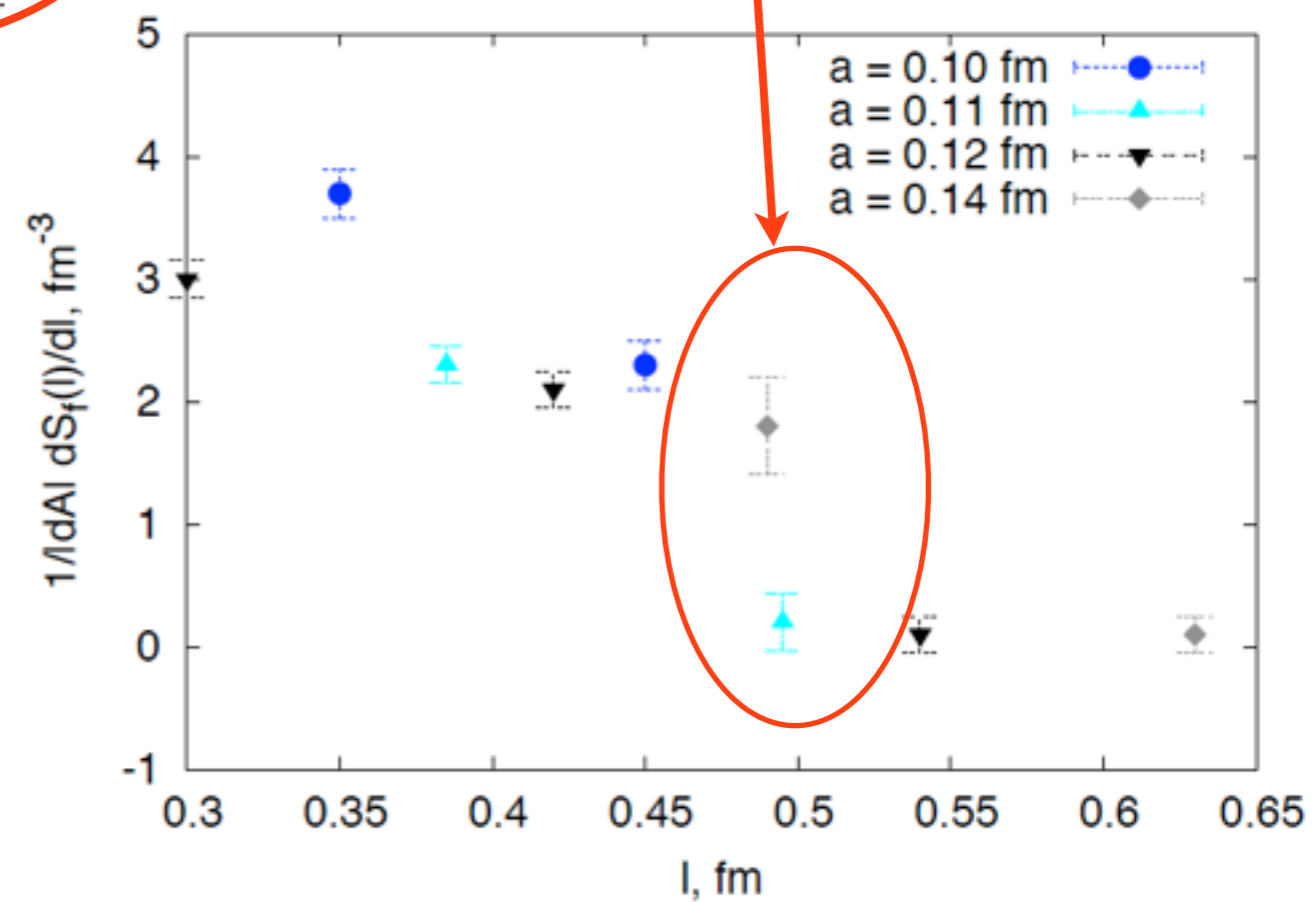
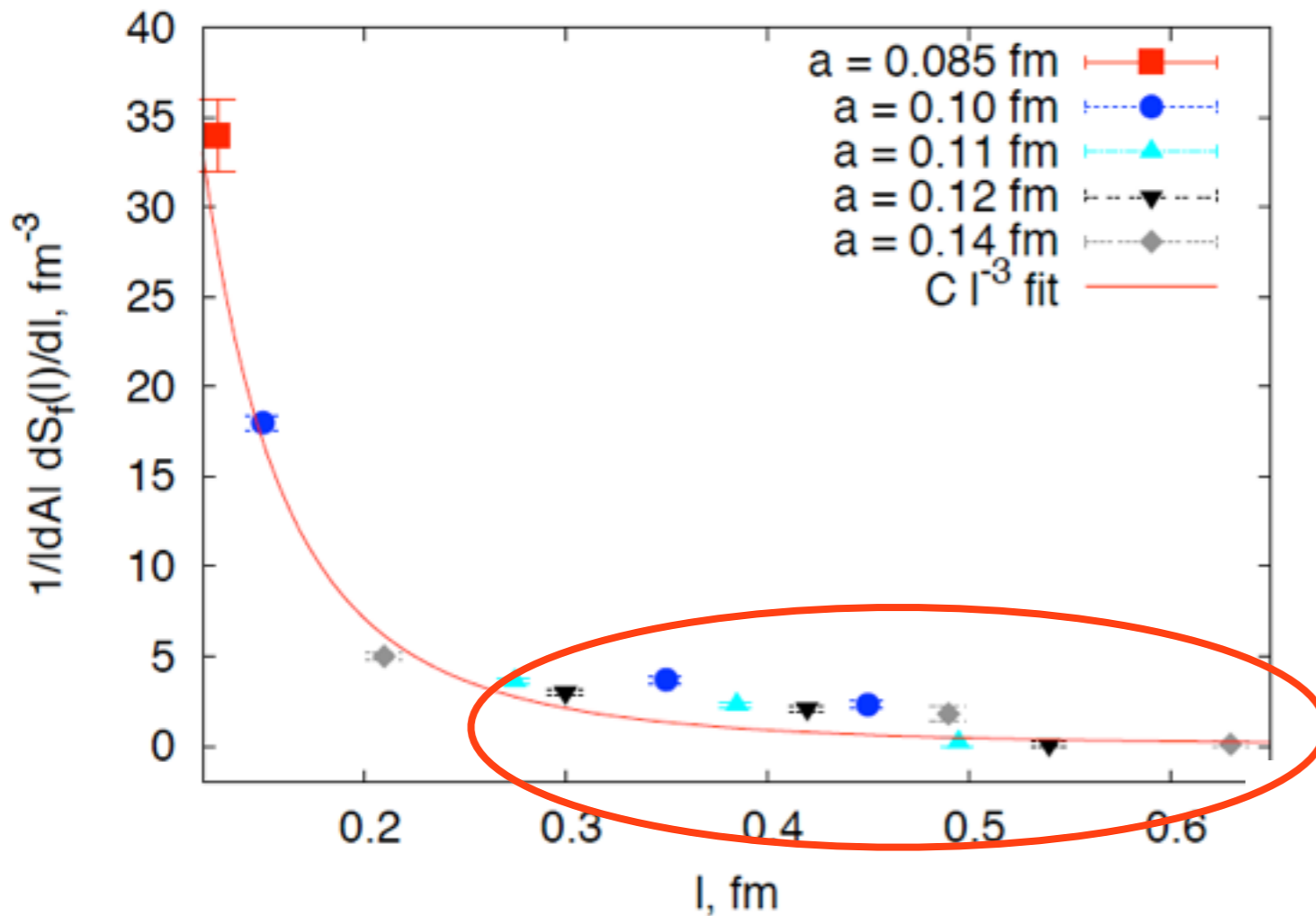


Simulation results

Lattice results for quenched SU(2)

Buividovich and Polikarpov:
NPB802(2008)458

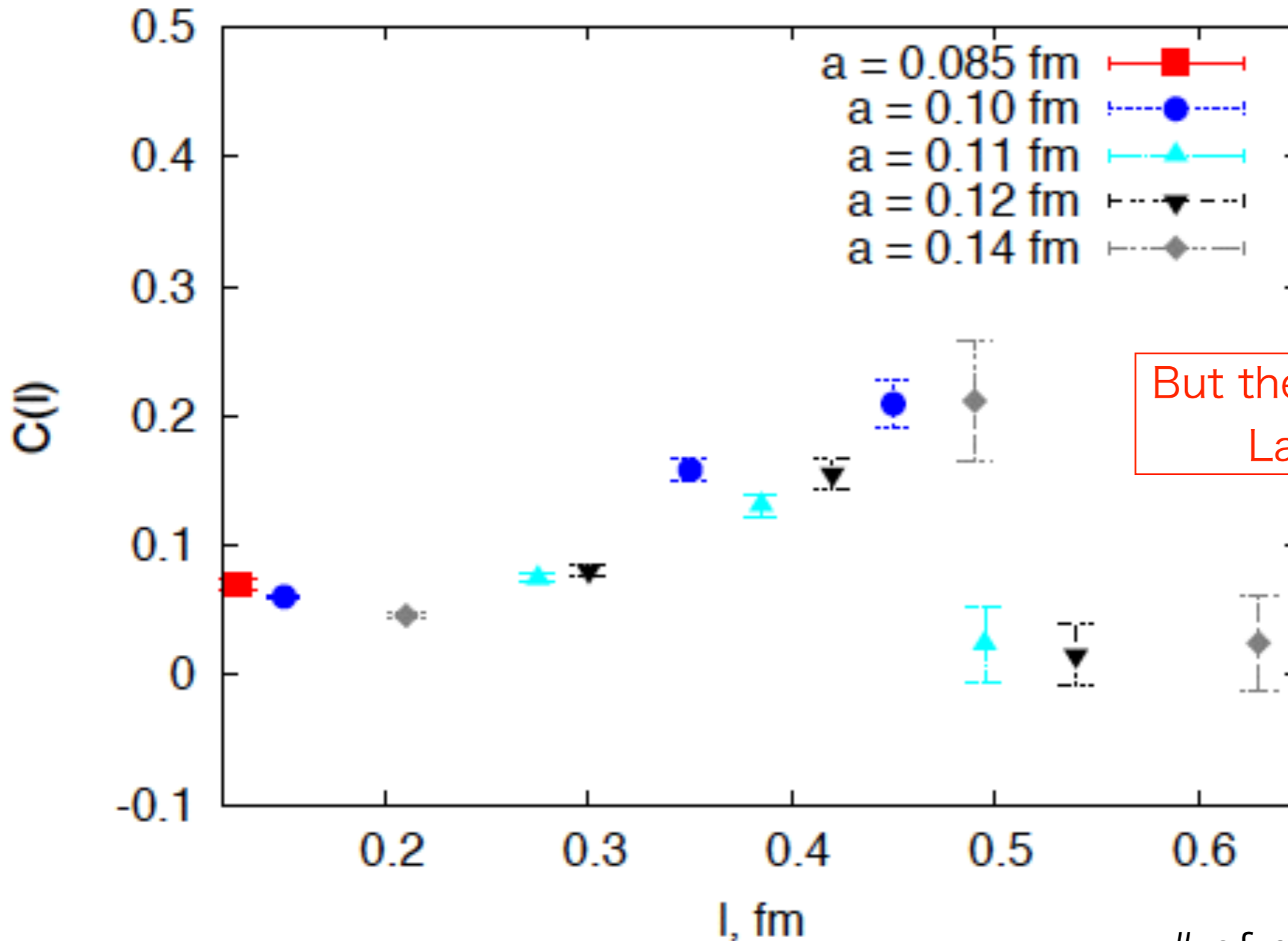
- in short range, $1/l^3$ scaling
- discontinuity is clear!?



Entropic C-function

$$C(l) = l^3 \frac{1}{|\partial A|} \frac{dS}{dl}$$

should be constant in short l region

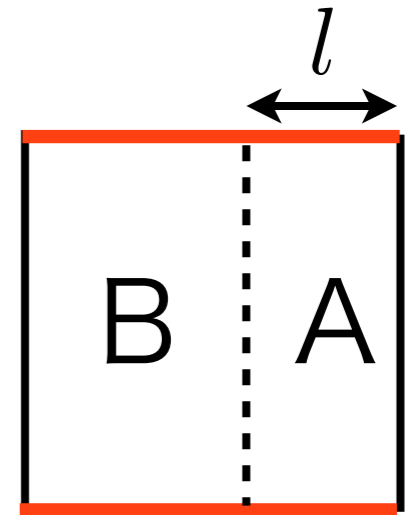


But the data is enhanced
Lattice artifact?

Our result

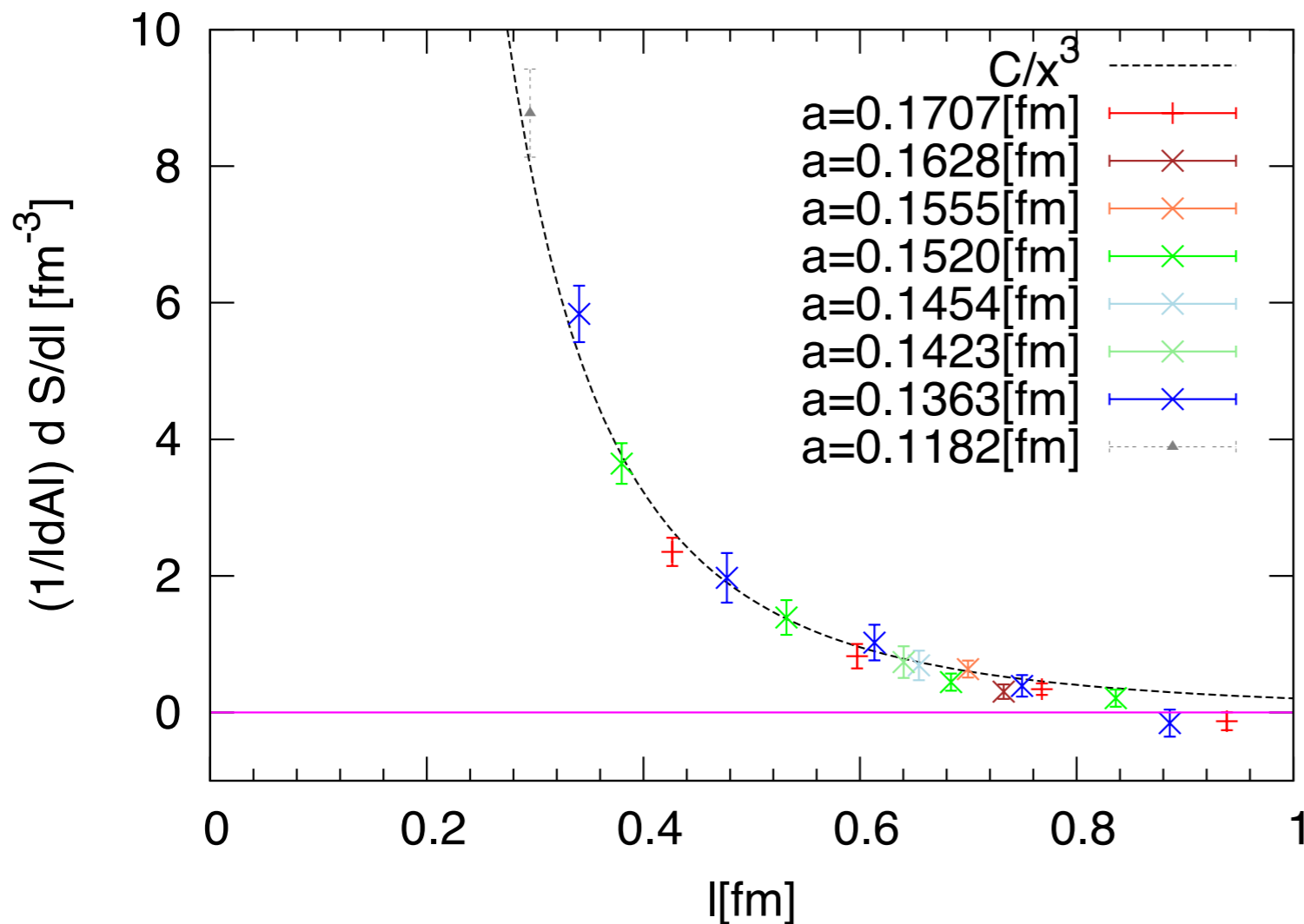
Simulation setup

- Wilson plaquette gauge action
- $N_s=N_t=16, 32$
- $l/a=2,3,4,5,(6)$
- $\beta=5.70 - 5.87$
- # of configuration 12,000~84,000
- scale setting $r_0 = 0.5$ fm and ALPHA coll.



Lattice results for quenched SU(3)

T=0, quenched QCD



We measure ~84,000 configs.

- in short range, $1/l^3$ scaling
- the coefficient is $C=0.2064(73)$

cf.) $C \sim 0.09$ in SU(2)

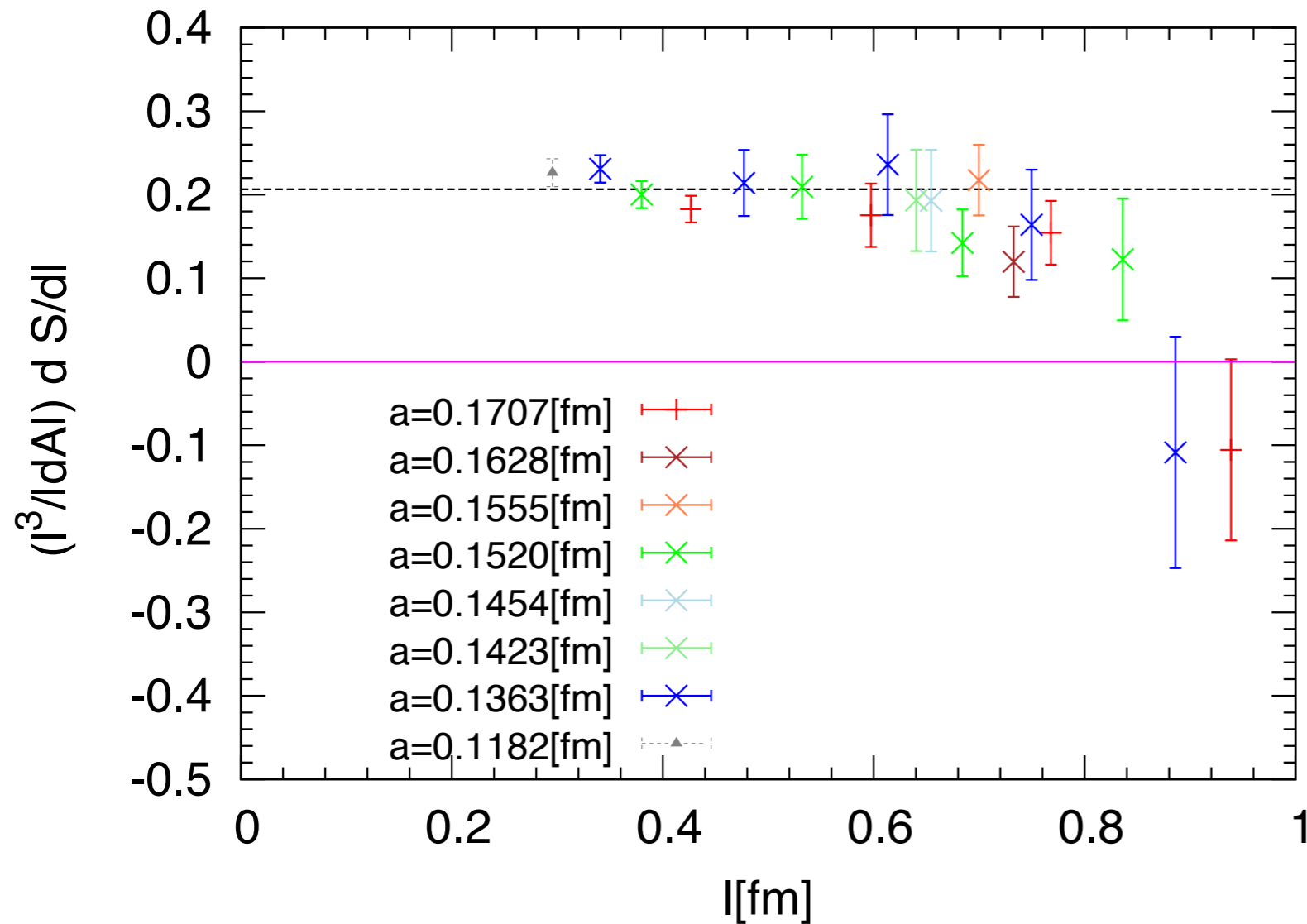
The N_c dependence is

$$\frac{0.2064}{0.09} = 2.29 \sim \frac{3^2}{2^2}$$

Entropic C-function

independent of UV cutoff

$$C(l) = l^3 \frac{1}{|\partial A|} \frac{dS}{dl}$$



In short l region, C is constant.

The discontinuity is not clear.

cf.) $\frac{1}{\Lambda_{QCD}} \sim 0.7$ [fm]

Comparison with Ryu-Takayanagi results

Ryu and Takayanagi:PRL96(2006)181602
JHEP 0608(2006)045

Holographic (or field theoretical) approach

$$(3+1)\text{-dim. CFT} \quad \frac{1}{|\partial A|} S_A(l) = c \frac{N_c^2}{a^2} - c' \frac{N_c^2}{l^2}$$

c' is obtained by AdS and QFT

$$c' \sim 0.0049 \quad \text{for free real scalar theory}$$

Estimation for non-abelian gauge theory $A_{\mu}^a \quad \left\{ \begin{array}{l} a = 1, \dots, 8 \\ \mu = 1, \dots, 4 \end{array} \right.$

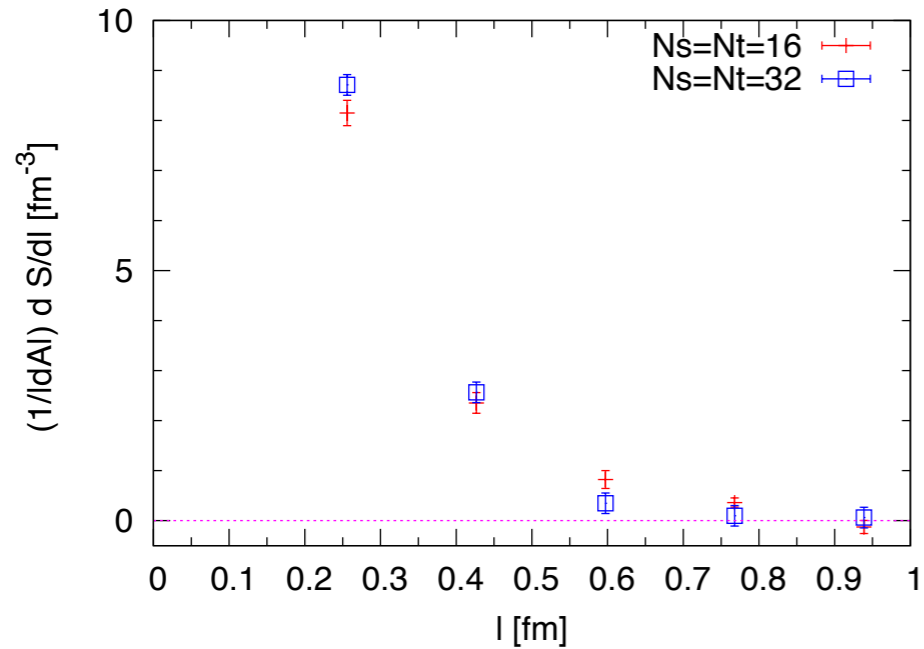
$$C_{\text{gauge}} \sim 2c' \cdot 2 \cdot 8 \sim 0.1568$$

Our numerical result

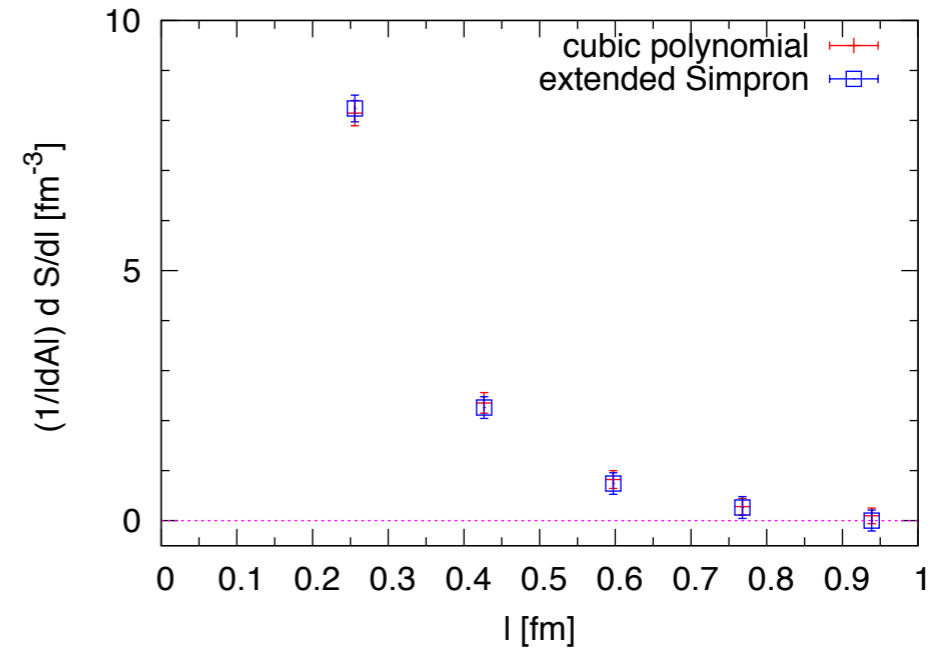
$$C_{\text{gauge}} \sim 0.2064$$

Detail analyses

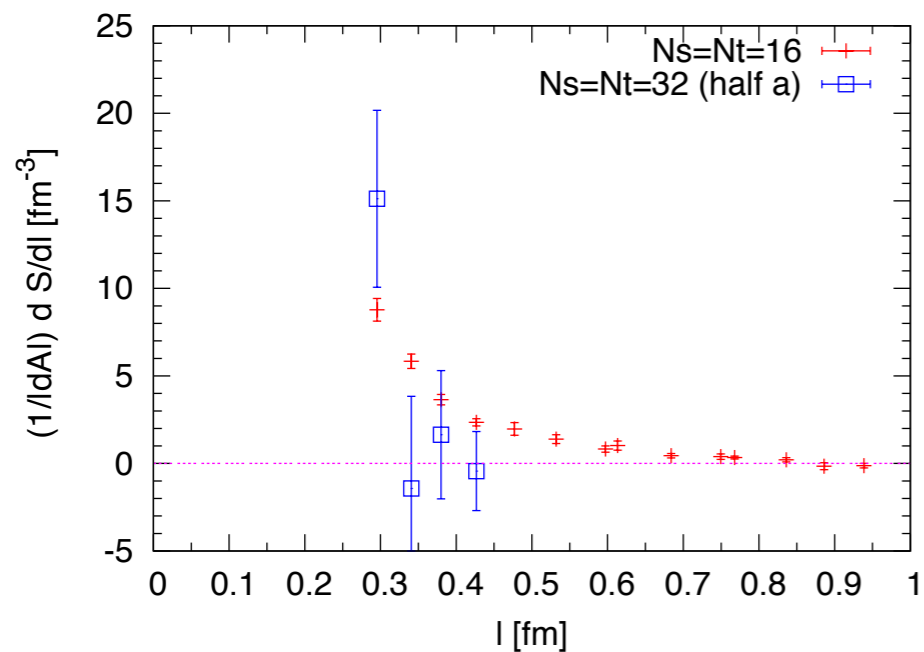
finite vol. effect



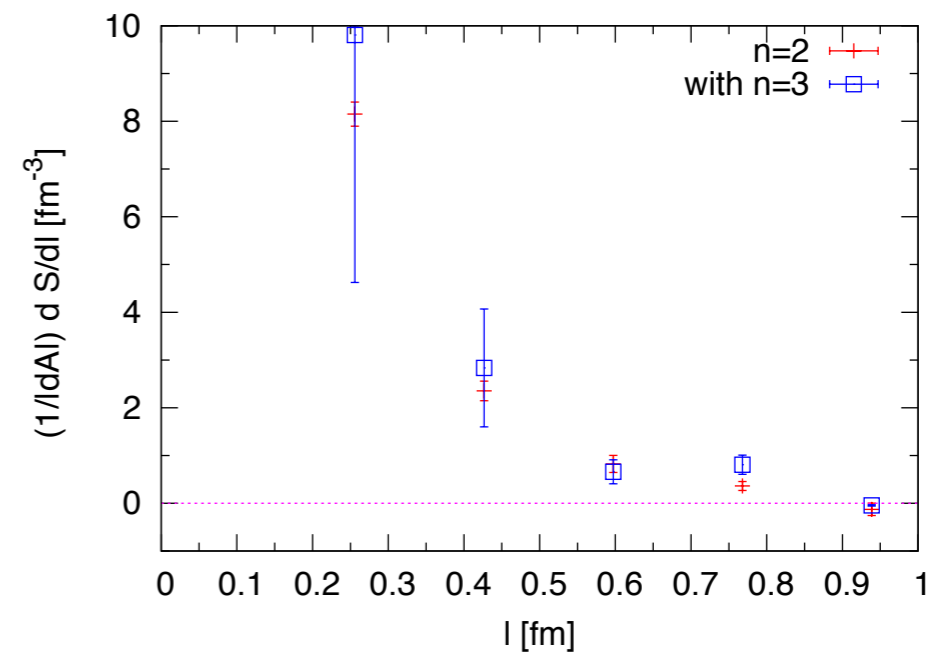
algorithm dependence of the numerical integration



UV cutoff dependence



replica number dependence

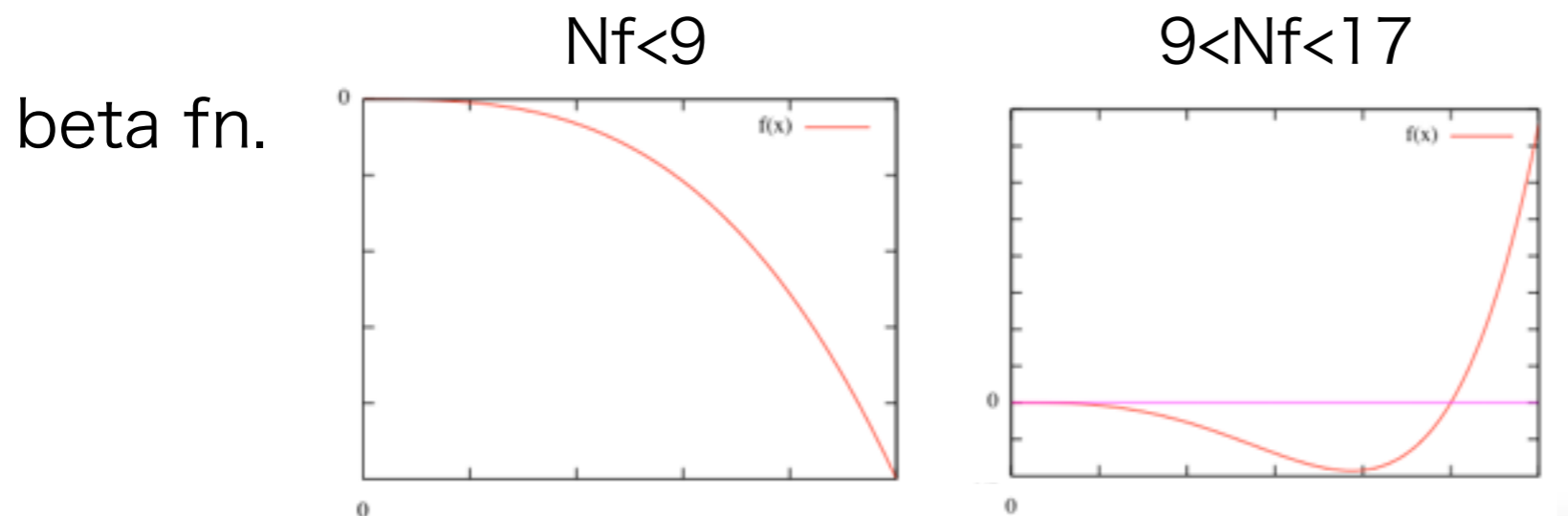


Summary

- This is the first precise determination of E.E. for quenched QCD
- N_c dependence in the short l region is N_c^2 as expected by AdS/CFT and field theoretical insights
- No discontinuity exists as contrast with SU(2) results
- Entropic C-function shows UV cutoff independence
- Value of C-function agrees with Ryu-Takayanagi work
- replica number ($n \rightarrow 1$) dependence

Future directions for E.E. using the lattice

- QCD at zero T
 - give a novel observation for confinement
 - even in full QCD case
- QCD at finite T
 - gives the thermal entropy and the correlation length in QGP phase
- conformal window in 4dim N_f flavor QCD
 - would give the a -function and central charge



nontrivial IR fixed point is found by lattice simulation
cf) E.I. PTEP(2013)083B01