

# Geometric Structure of MERA networks: Relation to AdS/CFT Correspondence

Hiroaki Matsueda  
(Sendai National College of Tech.)

# Classification of Tensor Network States

Tensor product states (TPS), Tensor network states (TNS)  
→ Variational ansatz satisfying the entanglement entropy scaling

Gapped quantum systems

→ PEPS class

PEPS: Projected Entanglement Pair States

Entanglement entropy: area law scaling       $S \propto L^{d-1}$

Quantum critical systems

→ MERA class

MERA: Multi-scale Entanglement Renormalization Ansatz

Entanglement entropy: logarithmic correction ( $d=1$ )

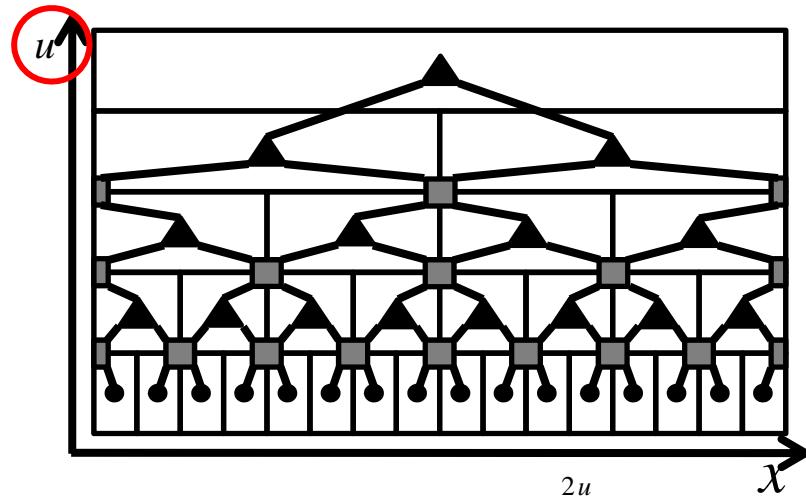
Holzhey–Larsen–Wilczek, Calabrese–Cardy

$$S = \frac{c}{3} \log L$$

# Purpose of This Talk

MERA

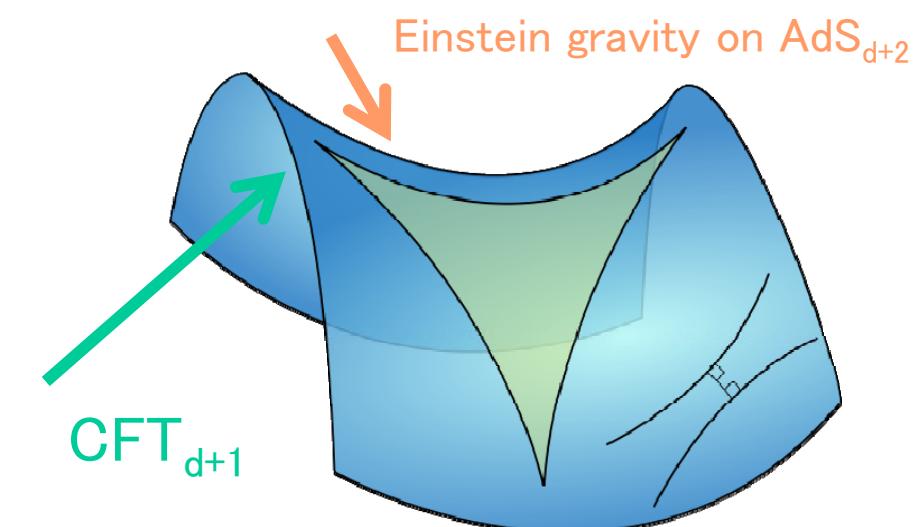
variational ansatz applicable to  
quantum critical systems



$$ds^2 = g_{uu} du^2 + \frac{e^{2u}}{\varepsilon^2} dx^2$$

AdS/CFT

quantum-classical correspondence  
in string theory



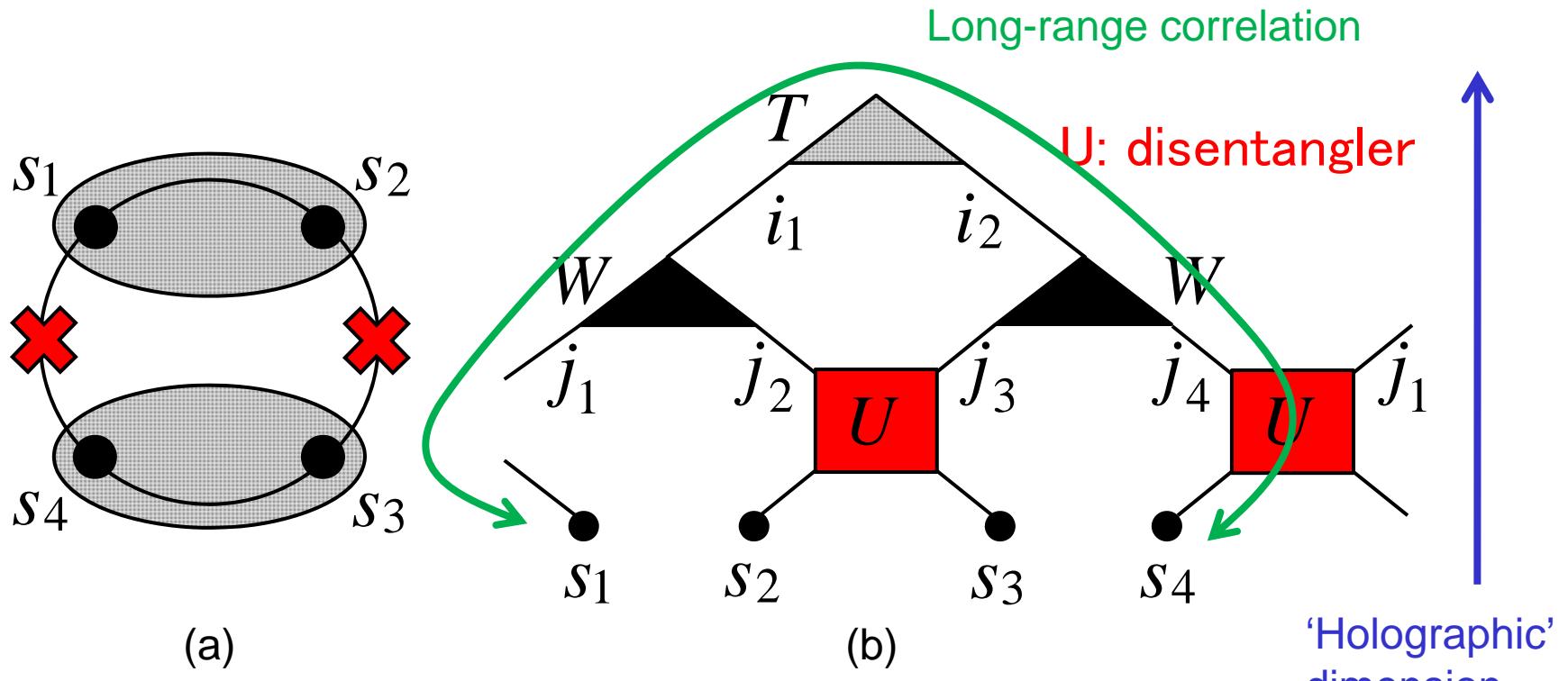
$$ds^2 = \frac{l^2}{z^2} (dz^2 + \eta_{ij} dx^i dx^j)$$

Question: how to examine their similarity ?

- (a) Information geometry for MERA network (Takayanagi)
- (b) Thermo-field double of MERA  $\Leftrightarrow$  BTZ black hole (HM)

# Hierarchical Tensor Network

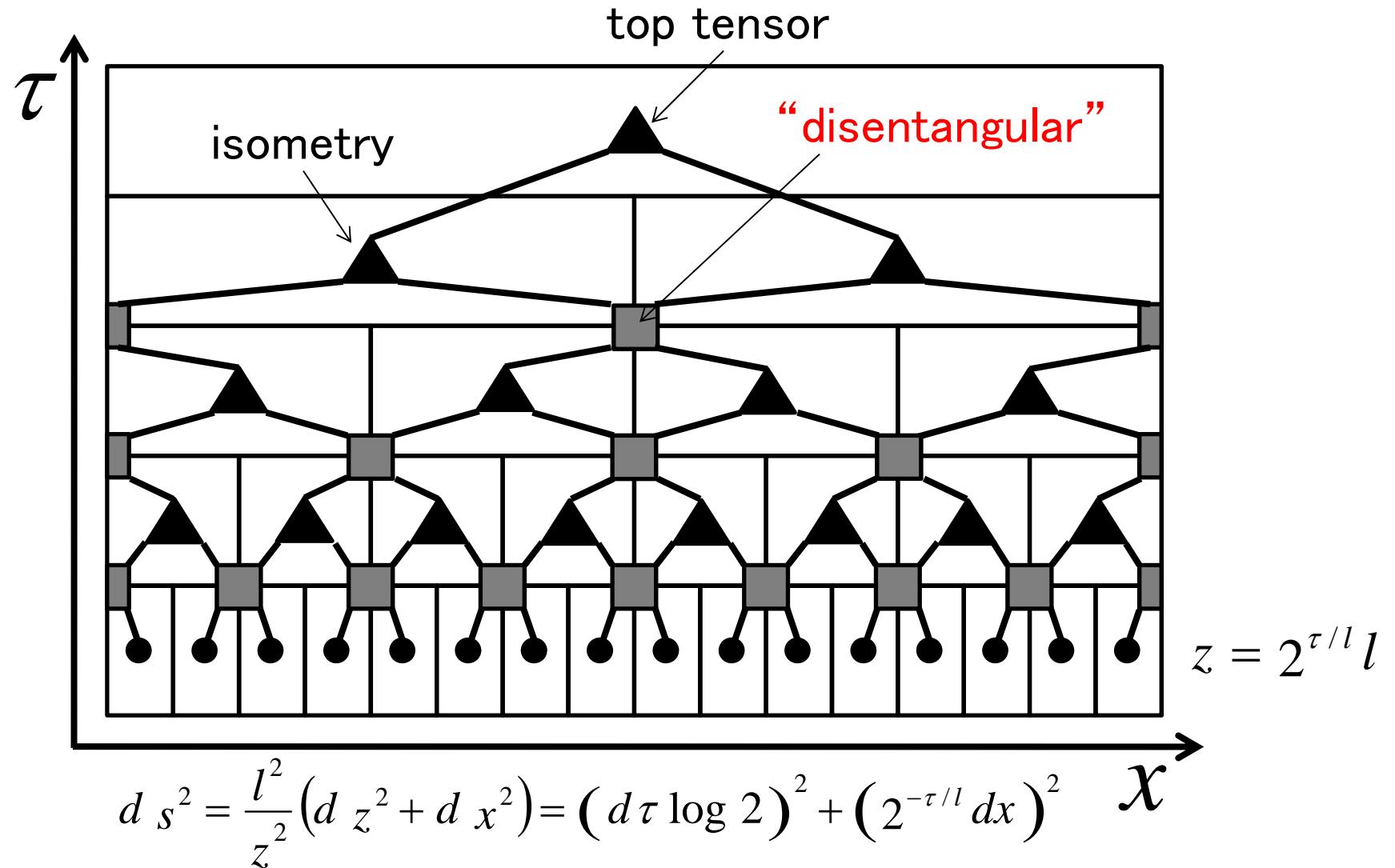
## Multiscale Entanglement Renormalization Ansatz (MERA)



$$|\psi\rangle = \sum_{s_1, s_2, s_3, s_4} T_{s_1 s_2 s_3 s_4} |s_1 s_2 s_3 s_4\rangle$$

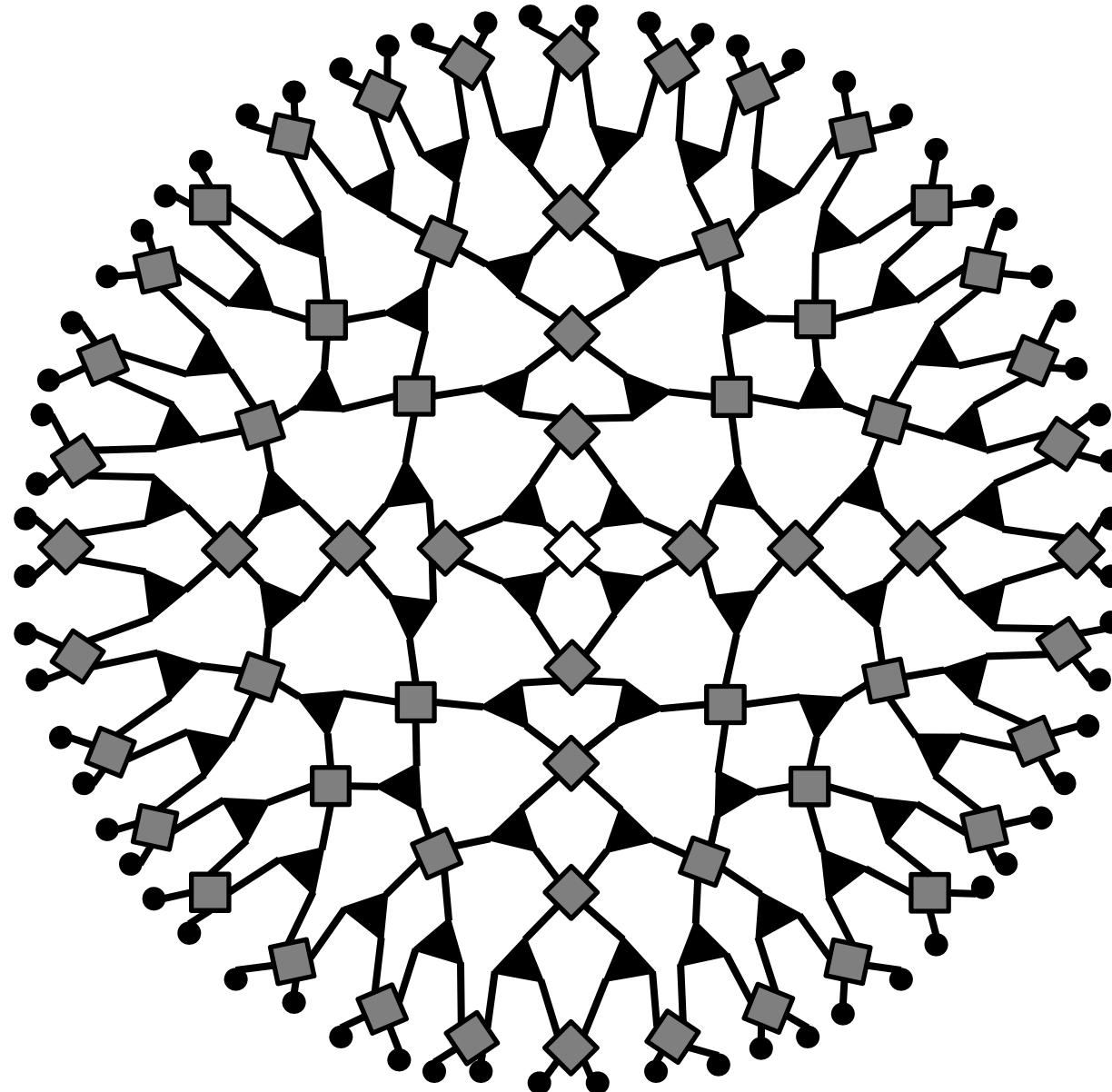
$$|\psi\rangle = \sum_{i_1, i_2} \sum_{j_1, \dots, j_4} \sum_{s_1, \dots, s_4} T_{i_1 i_2} W_{j_1 j_2}^{i_1} W_{j_3 j_4}^{i_2} U_{s_2 s_3}^{j_2 j_3} U_{s_4 s_1}^{j_4 j_1} |s_1 s_2 s_3 s_4\rangle$$

(Multiscale Entanglement Renormalization Ansatz, MERA)



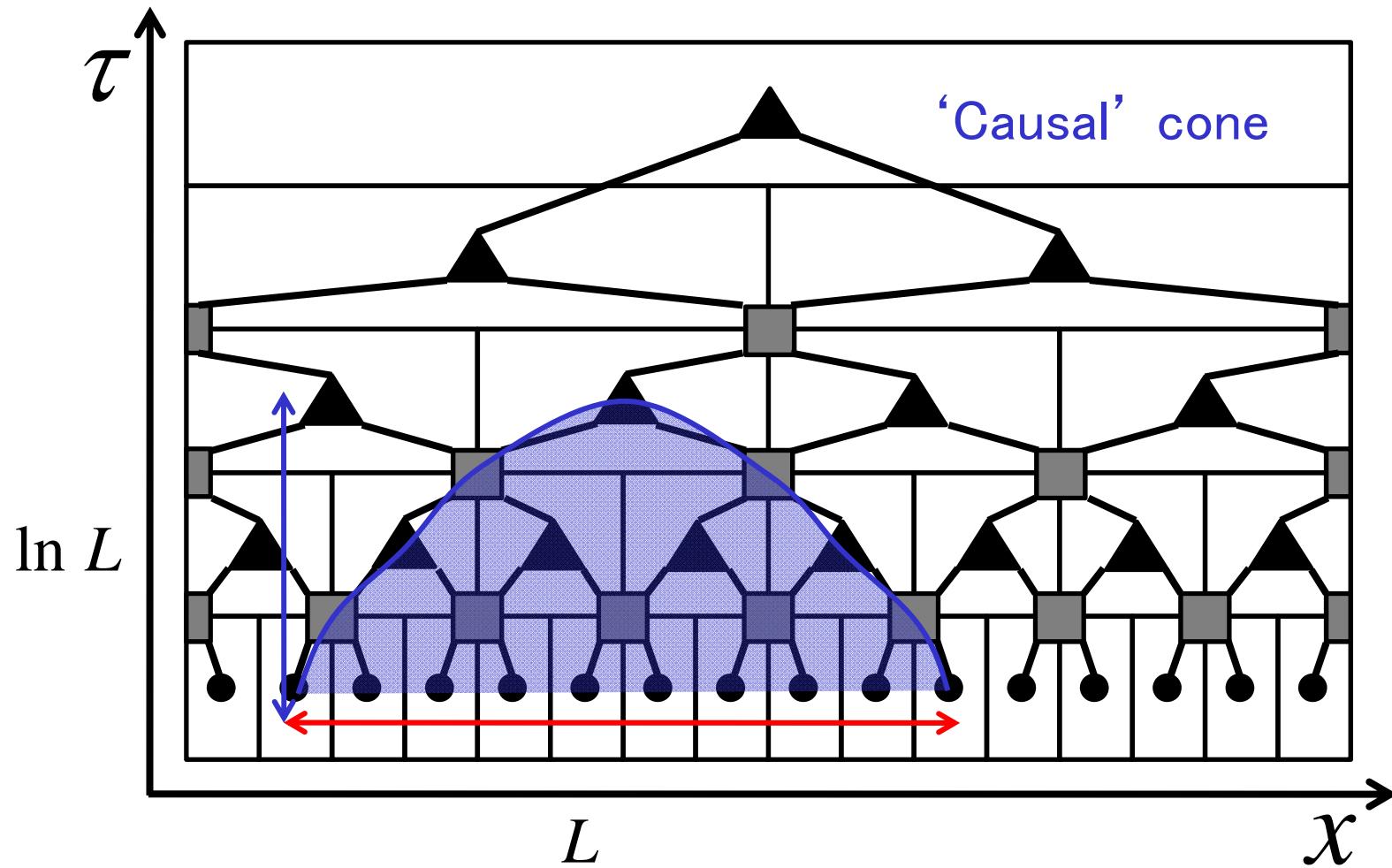
MPS → decomposed into many tensors with different function  
 Basis change (disentangler) before renormalization

# Poincare Disk Model for MERA Network



How to evaluate entanglement entropy in holographic space ?

Close connection to ‘Ryu–Takayanagi formula’  
developed in superstring theory



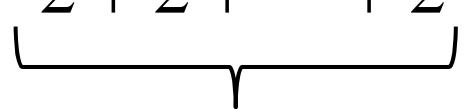
$S = \text{minimal surface area in holographic space}$

Binary decomposition



Spatially 1D cases:

$$2 + 2 + \cdots + 2 = 2 \ln L$$



No. of boundary points:  $\ln L$

Spatially 2D cases:

$$4L\left(1 + \frac{1}{2} + \frac{1}{2^2} + \cdots + \frac{1}{2^n}\right) = 4L\left(2 - \frac{1}{2^n}\right) \rightarrow 8L$$

# Anti-de Sitter (Hyperbolic) Space and CFT

Metric of AdS space

$z$ : radial axis,  $z \rightarrow 0$ : boundary

$$\eta_{ij} = \begin{pmatrix} -1 & 0 & \cdots \\ 0 & 1 & \cdots \\ \vdots & \vdots & \ddots \end{pmatrix}$$

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu = \frac{l^2}{z^2} (dz^2 + \eta_{ij} dx^i dx^j)$$

$$ds \sim \frac{l}{z} dz$$

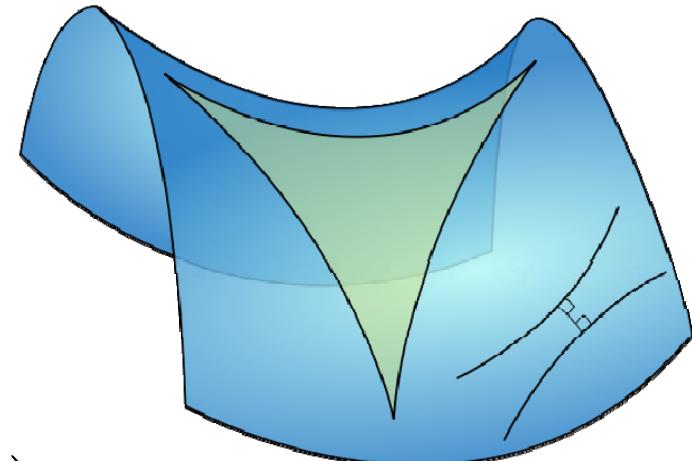
$$s \sim l \log z$$

Infinitesimal trans.

$$\begin{aligned}\bar{x}^i &= x^i + \xi^i(x) \\ \bar{z} &= z + z\zeta(x)\end{aligned}$$

$$z \rightarrow 0 \quad \downarrow$$

$$\begin{aligned}d\bar{s}^2 &= \bar{g}_{\mu\nu} d\bar{x}^\mu d\bar{x}^\nu \\ &= d s^2 + (\partial_i \xi_j + \partial_j \xi_i - 2\zeta \eta_{ij}) d x^i d x^j\end{aligned}$$



Isometry trans.  $\Rightarrow$  conformal Killing equation at  $z \rightarrow 0$

Boundary of  $\text{AdS}_{d+1} \Rightarrow \text{CFT}_d$

# AdS/CFT correspondence and Holographic Entropy

Gubser–Klevanov–Polyakov(GKP)–Witten relation

$$\langle O(x_1) \cdots O(x_n) \rangle_{CFT} = \frac{\delta}{\delta \phi(x_1)} \cdots \frac{\delta}{\delta \phi(x_n)} \exp \left( -\frac{1}{2\kappa} I(\phi(x)) \right) \Big|_{\phi=0}$$

Ryu–Takayanagi Formula (extension of Beckenstein–Hawking)

2006

$$S = \frac{\gamma}{4G}$$

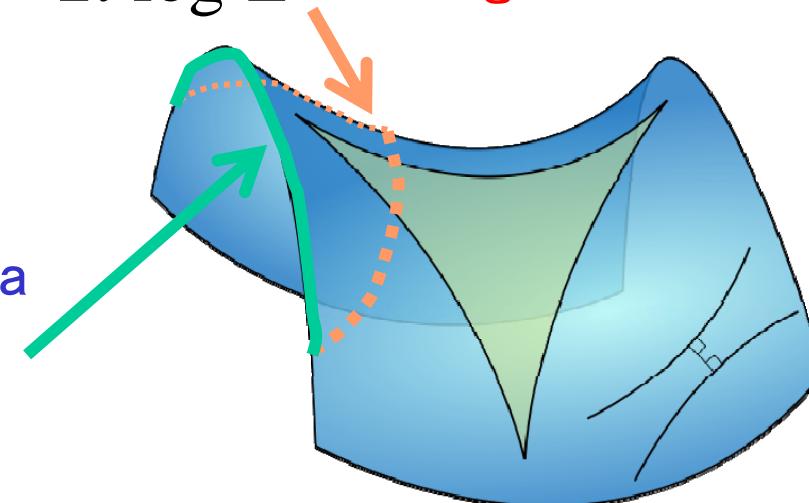
$$\gamma = 2l \log L$$

2D: geodesic distance

$\gamma$ : area of minimal surface

Calabrese–Cardy formula

$$S = \frac{1}{3} c \log L$$



$$c = \frac{3l}{2G}$$

Brown–Henneaux central charge

## Geometry of quantum states

Difference between two similar states

→ distance on information space (Bures metric)

$$D(\theta) = 1 - \left| \langle \psi(\theta) | \psi(\theta + d\theta) \rangle \right|^2$$

$\Theta$ : vector valued, internal parameters of  $\Psi$

Expansion of D by  $d\theta$  up to the second order

$$D(\theta) = \chi_{\mu\nu}(\theta) d\theta^\mu d\theta^\nu$$

$$\chi_{\mu\nu}(\theta) = \langle \partial_\mu \psi(\theta) | \partial_\nu \psi(\theta) \rangle - \langle \partial_\mu \psi(\theta) | \psi(\theta) \rangle \langle \psi(\theta) | \partial_\nu \psi(\theta) \rangle$$

$$\chi_{\mu\nu}(\theta) = \chi_{\nu\mu}^*(\theta)$$

Fisher metric g and Berry curvature F

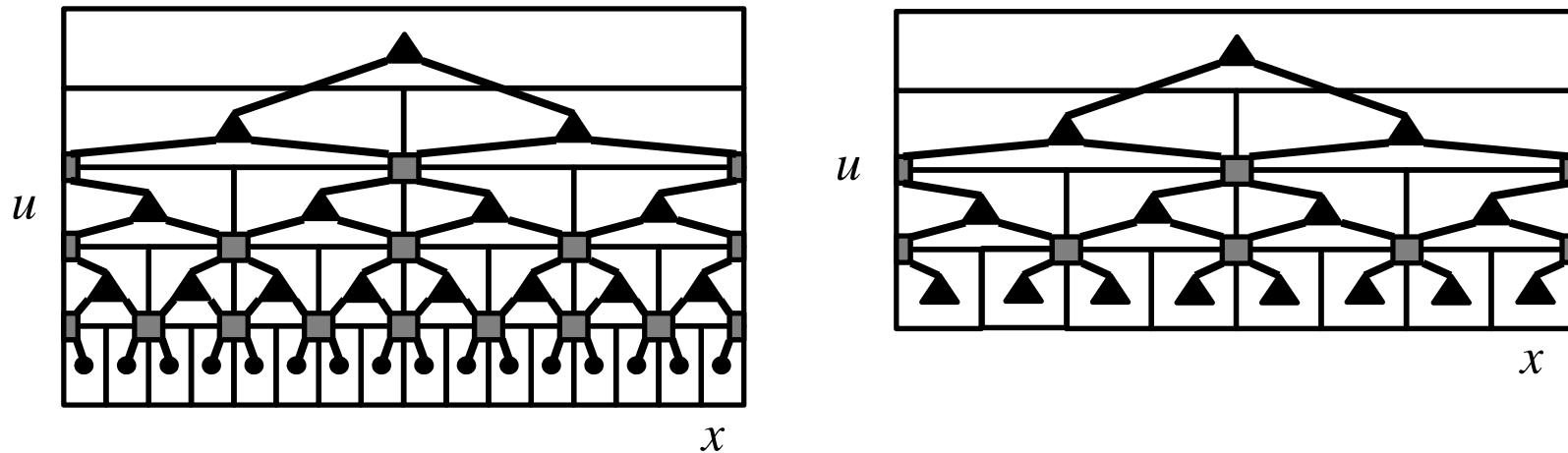
$$g_{\mu\nu}(\theta) = \frac{\chi_{\mu\nu}(\theta) + \chi_{\nu\mu}(\theta)}{2} \quad F_{\mu\nu}(\theta) = i(\chi_{\mu\nu}(\theta) - \chi_{\nu\mu}(\theta))$$

# Emergent AdS spacetime from continuous MERA

Question:

Does the holographic direction of MERA network really correspond to the radial axis of AdS spacetime ?

$$u \subset \theta$$



→ Calculation of  $g_{uu}(u)$

$$D(u) = 1 - \left| \langle MERA(u) | MERA(u + du) \rangle \right|^2$$

RG changes the number of effective lattice sites and it is hard to take the inner product directly → field-theoretical treatment

# Physical Information and Geometric Distance

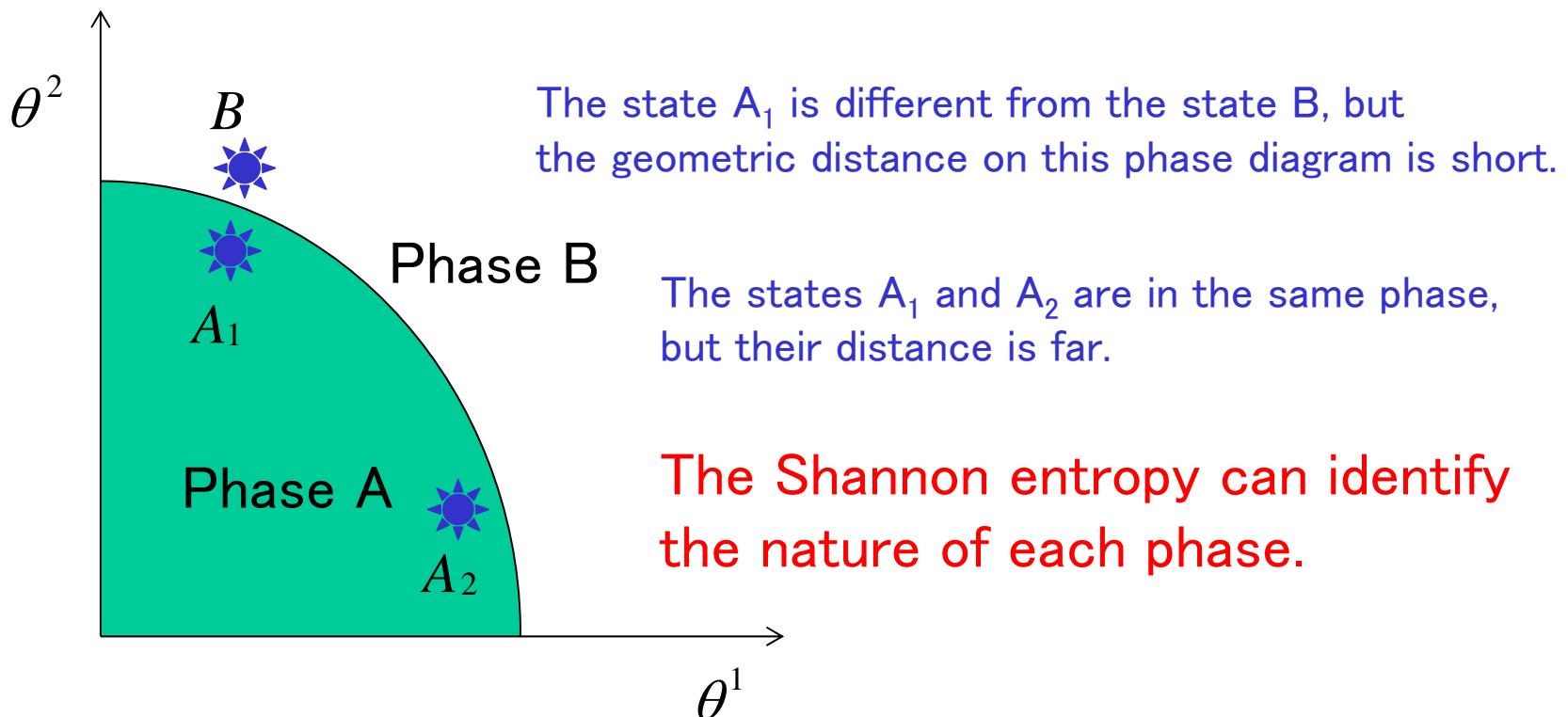
Phase diagram for states of matter

→ ‘highly compressed’ quantum information

Only the information entropy (amount of information, not quality) has some physical meaning.

Equation of states for ideal gas  $PV = nRT$

(Avogadro number → only three parameters P,V,T)



## Relative information entropy → Fisher Metric

Probability distribution  $\sum_n p_n(\theta) = 1$

$\theta$  : internal parameter (vector valued)

→ This parameter set determines a particular physical state.

‘Relative’ information entropy

$$D(\theta) = -\sum_n p_n(\theta) \log p_n(\theta) + \sum_n p_n(\theta) \log p_n(\theta + d\theta)$$

$$\approx \frac{1}{2} \sum_n p_n(\theta) \frac{\partial \log p_n(\theta)}{\partial \theta^\mu} \frac{\partial \log p_n(\theta)}{\partial \theta^\nu} d\theta^\mu d\theta^\nu$$

Fisher metric

$$\gamma_n(\theta) = -\log p_n(\theta)$$

$$g_{\mu\nu}(\theta) = \sum_n p_n(\theta) \frac{\partial \log p_n(\theta)}{\partial \theta^\mu} \frac{\partial \log p_n(\theta)}{\partial \theta^\nu} = \langle \partial_\mu \gamma \partial_\nu \gamma \rangle$$

# cMERA

Continuous version of MERA (cMERA)

Evolution operator along holographic direction

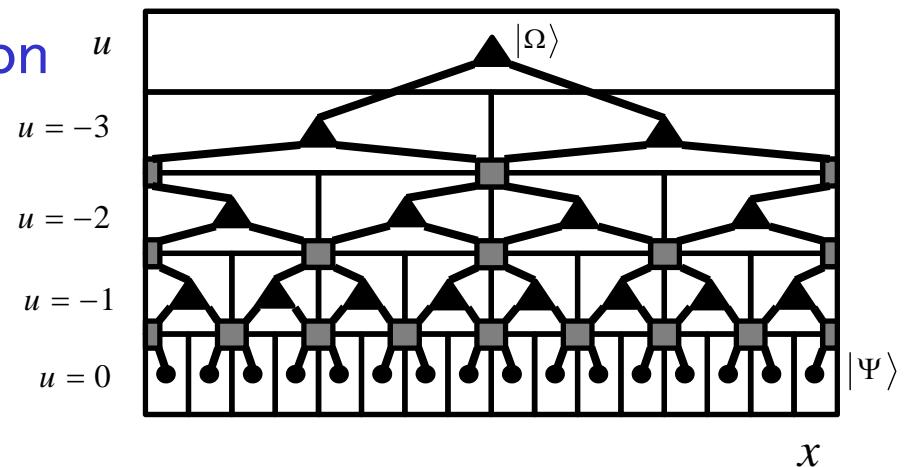
$$U(u_1, u_2) = P \exp \left[ -i \int_{u_1}^{u_2} (K(u) + L) du \right]$$

K: entangler, L: scale transformation

$$k \leq \Lambda e^u$$

IR & UV limit

$$|\Omega\rangle = |\Psi(u_{IR})\rangle, |\Psi\rangle = |\Psi(u_{UV} = 0)\rangle$$



Any state with a particular scale  $u$  is constructed by the evolution from a reference state at IR or UV.

$$|\Psi\rangle = U(0, u)|\Psi(u)\rangle \quad |\Psi(u)\rangle = U(u, u_{IR})|\Omega\rangle \quad L|\Omega\rangle = 0$$

## Interaction representation

$$K_I(u) = e^{iuL} K(u) e^{-iuL} \quad U(u_1, u_2) = e^{-iu_1 L} P \exp \left[ -i \int_{u_1}^{u_2} K_I(s) ds \right] e^{iu_2 L}$$

$$|\Psi(u)\rangle = e^{-iuL} P \exp \left[ -i \int_{u_{IR}}^u K_I(s) ds \right] |\Omega\rangle = e^{-iuL} |\Phi(u)\rangle$$

### Explicit form of entangler (assumption) for 1D free scalar case

$$H = \frac{1}{2} \int dk \left[ \pi(k) \pi(-k) + \varepsilon_k^2 \phi(k) \phi(-k) \right] \quad \varepsilon_k = \sqrt{k^2 + m^2}$$

$$K(u) = \frac{1}{2} \int dk g(k, u) [\phi(k) \pi(-k) + \pi(k) \phi(-k)]$$

Energy minimization

Scale invariance of  $\Omega$

Scaling properties of  $\Phi, \pi$



$$g(k, u) \sim \chi(u) \theta(\Lambda - |k|)$$

$$k \leq \Lambda e^u$$

$$\chi(s) = \frac{1}{2} \frac{e^{2u}}{e^{2u} + m^2 / \Lambda^2}$$

## Derivation of Bures metric

$$g_{uu}(u) = \langle \Phi(u) | K_I(u)^2 | \Phi(u) \rangle - \langle \Phi(u) | K_I(u) | \Phi(u) \rangle^2$$

$$g_{uu}(u) = \chi(u)^2 = \frac{e^{4u}}{4(e^{2u} + m^2/\Lambda^2)^2}$$



$$e^{2u} = \frac{1}{\Lambda^2 z^2} - \frac{m^2}{\Lambda^2}$$

$$ds^2 = \frac{dz^2}{4z^2} + \left( \frac{1}{\Lambda^2 z^2} - \frac{m^2}{\Lambda^2} \right) dx^2 + g_{tt} dt^2$$

Massless case  $\rightarrow$  pure AdS

Finite mass  $\rightarrow$  truncation of the network  $z < \frac{1}{m}$

## Thermofield dynamics (TFD) for finite-T wavefunction

Purpose: finite-T MERA and its relation with AdS/CFT

Finite-T  $\rightarrow$  thermal average

TFD form  $\rightarrow$  ‘thermal vacuum’

Identity state (maximally entangled)  $|I\rangle = \sum_n |n\rangle \otimes |\tilde{n}\rangle$

General representation theorem

$$|I\rangle = \sum_n |n\rangle \otimes |\tilde{n}\rangle = \sum_{\alpha} |\alpha\rangle \otimes |\tilde{\alpha}\rangle$$

$$|O(\beta)\rangle = \rho^{1/2} |I\rangle$$

$$\begin{aligned} \langle O(\beta) | A | O(\beta) \rangle &= \sum_{m,n} \langle m \tilde{m} | \rho^{1/2} A \rho^{1/2} | n \tilde{n} \rangle \\ &= \sum_{m,n} \langle m | \rho^{1/2} A \rho^{1/2} | n \rangle \delta_{\tilde{m}\tilde{n}} = \text{tr}(\rho A) \end{aligned}$$

## Thermal state in TFD

$$|\psi(\beta)\rangle = \sum_{\{m_j\}} \sum_{\{\tilde{n}_j\}} c^{\{m_j\}\{\tilde{n}_j\}} |\{m_j\}\{\tilde{n}_j\}\rangle$$

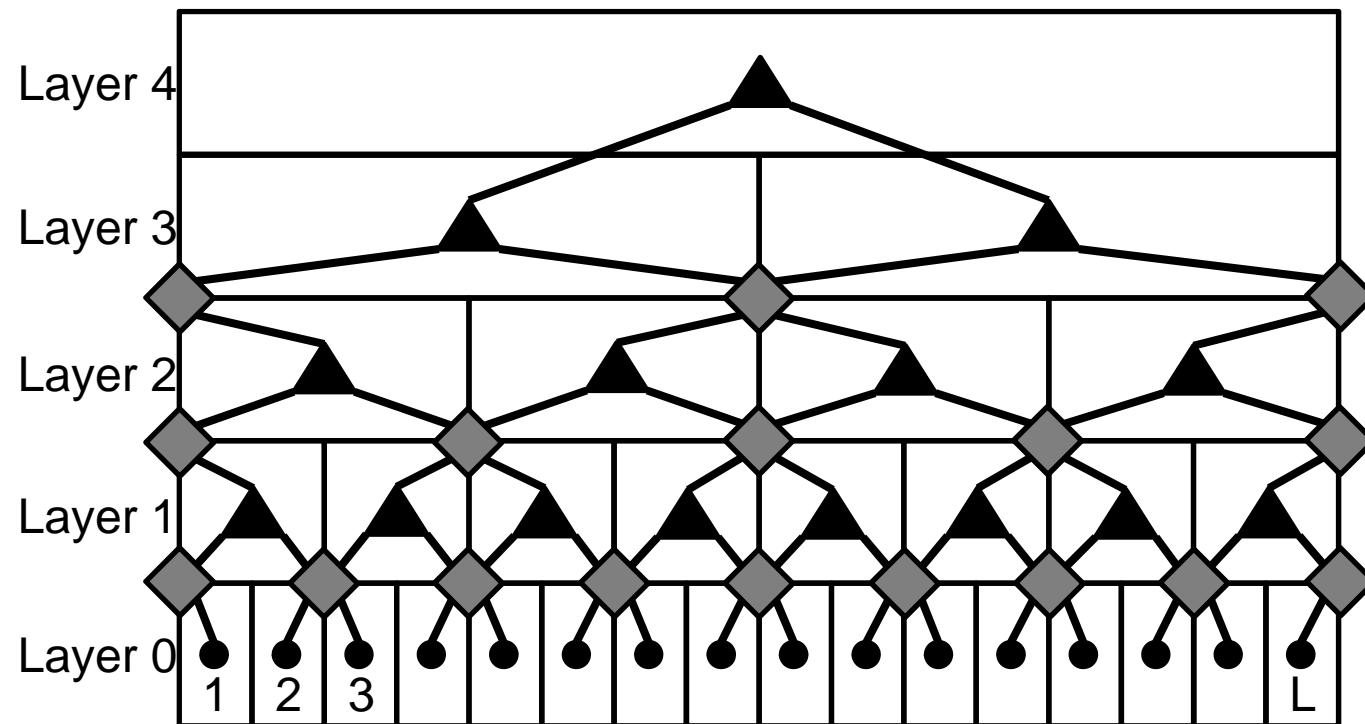
Singular value decomposition

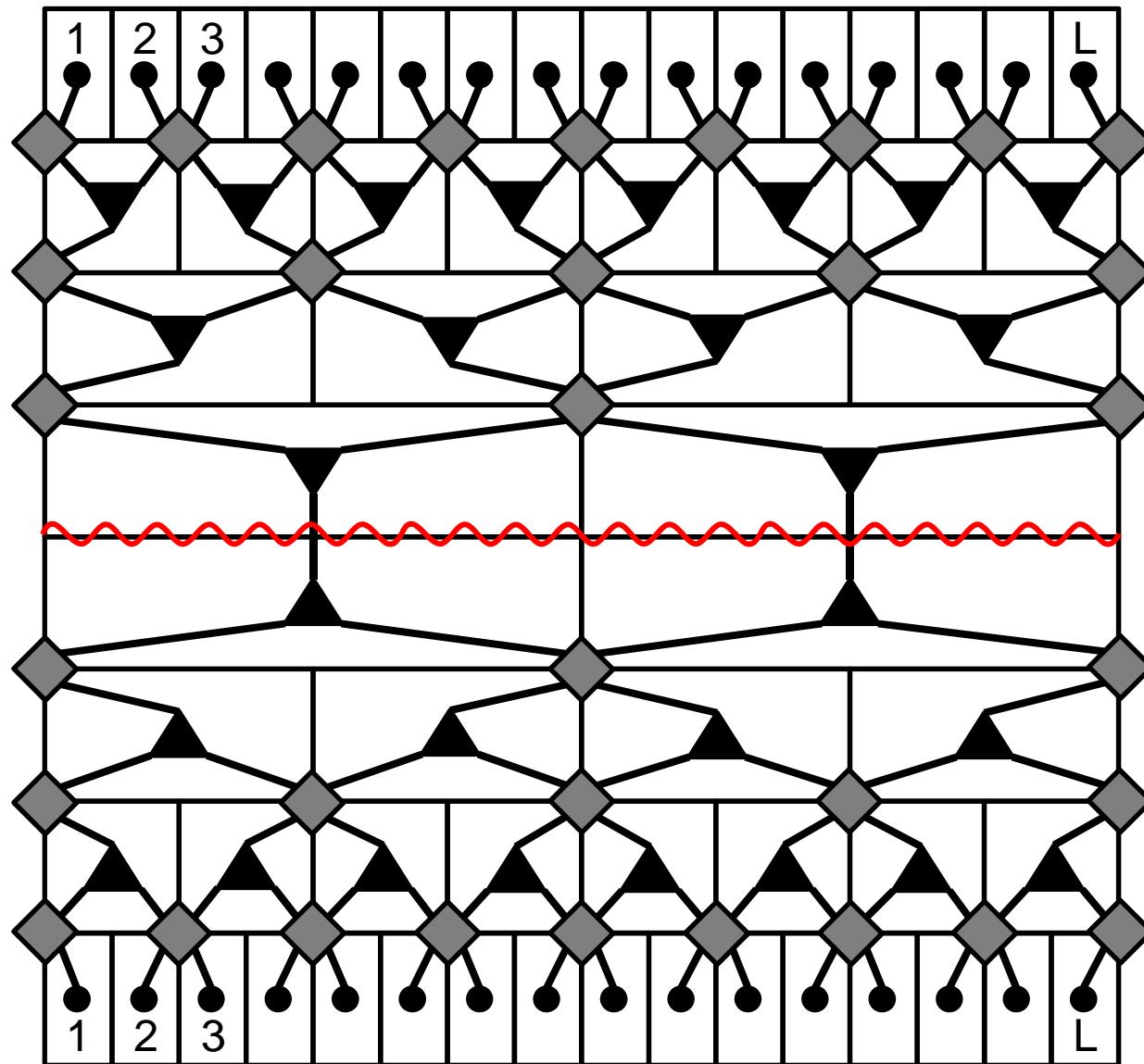
$$c^{\{m_j\}\{\tilde{n}_j\}} = \sum_{\alpha=1}^{\chi} A_\alpha^{\{m_j\}} A_\alpha^{\{\tilde{n}_j\}}$$

$\alpha$  : event horizon  $\rightarrow$  black hole entropy  
= maximally entanglement entropy

Imagine Penrose diagrams ...

$T=0$





## Banados–Teitelboim–Zanelli (BTZ) metric: black hole solution

Exact solution of Einstein eq. in (2+1)D  
with negative cosmological term

\* Schwarzschild solution  
in (3+1)D flat spacetime

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R + g_{\mu\nu} \Lambda = 0 \quad f(z) = 1 - \frac{a}{z}$$

$$ds^2 = \frac{1}{z^2} \left( -f(z) dt^2 + \frac{1}{f(z)} dz^2 + dx^2 \right)$$

**Event horizon:**  $z = z_H$

$$f(z) = 1 - \left( \frac{z}{z_H} \right)^2$$

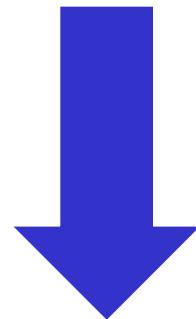
$f(z)=1 \rightarrow$  anti-de Sitter (AdS) spacetime

$$ds^2 = \frac{1}{z^2} \left( -dt^2 + dz^2 + dx^2 \right)$$

## Maximally-extended of BTZ spacetime

Coordinate transformation

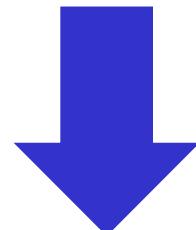
$$ds^2 = \frac{1}{z^2} \left( -f(z) dt^2 + \frac{1}{f(z)} dz^2 + dx^2 \right)$$



$$j = \ln \left( \frac{2z/\varepsilon}{1 + \sqrt{f(z)}} \right) \quad j_H = \ln \left( \frac{2z_H}{\varepsilon} \right)$$

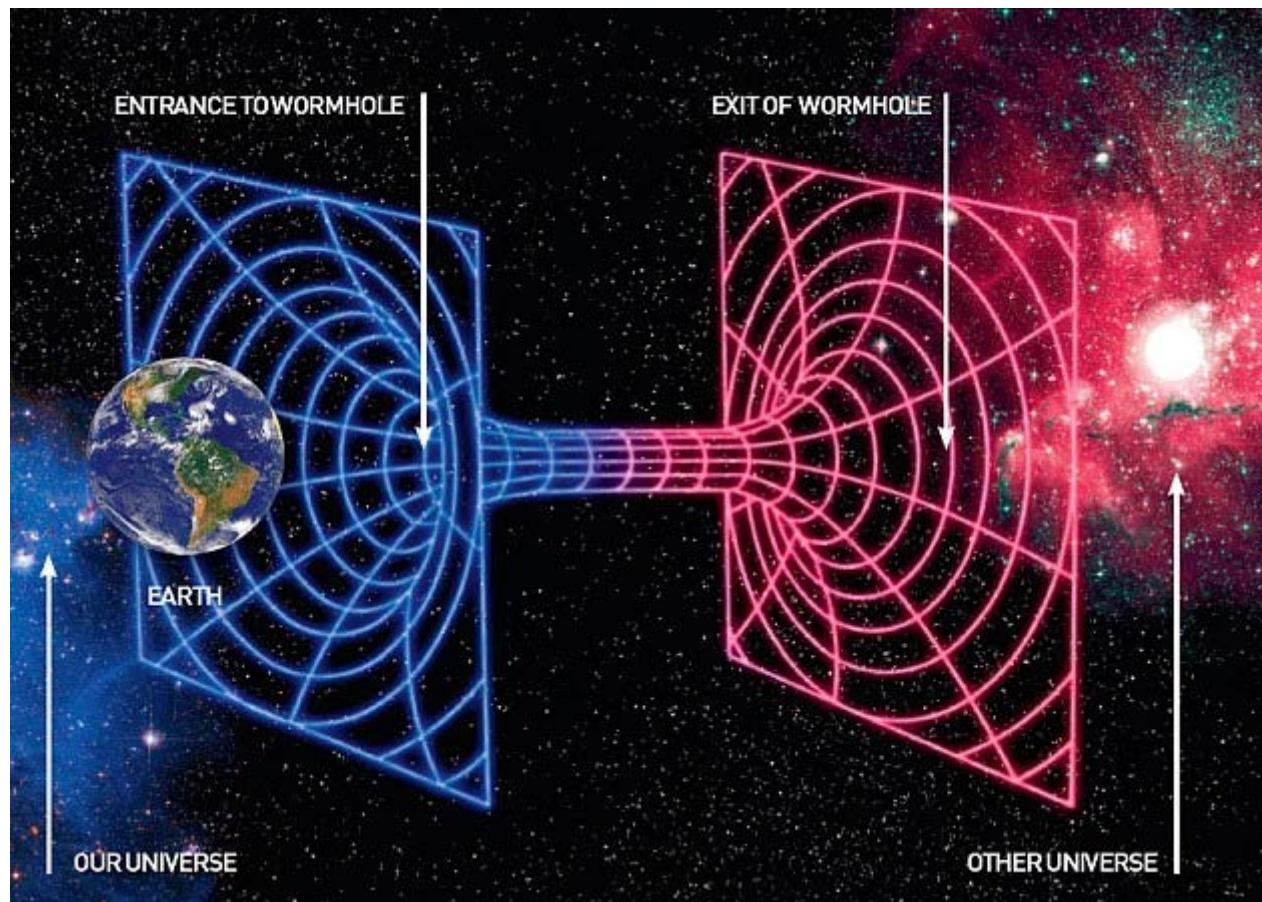
$\varepsilon$  : UV cut-off

$$ds^2 = -h(j) dt^2 + d(j \ln \eta)^2 + g(j) dx^2$$

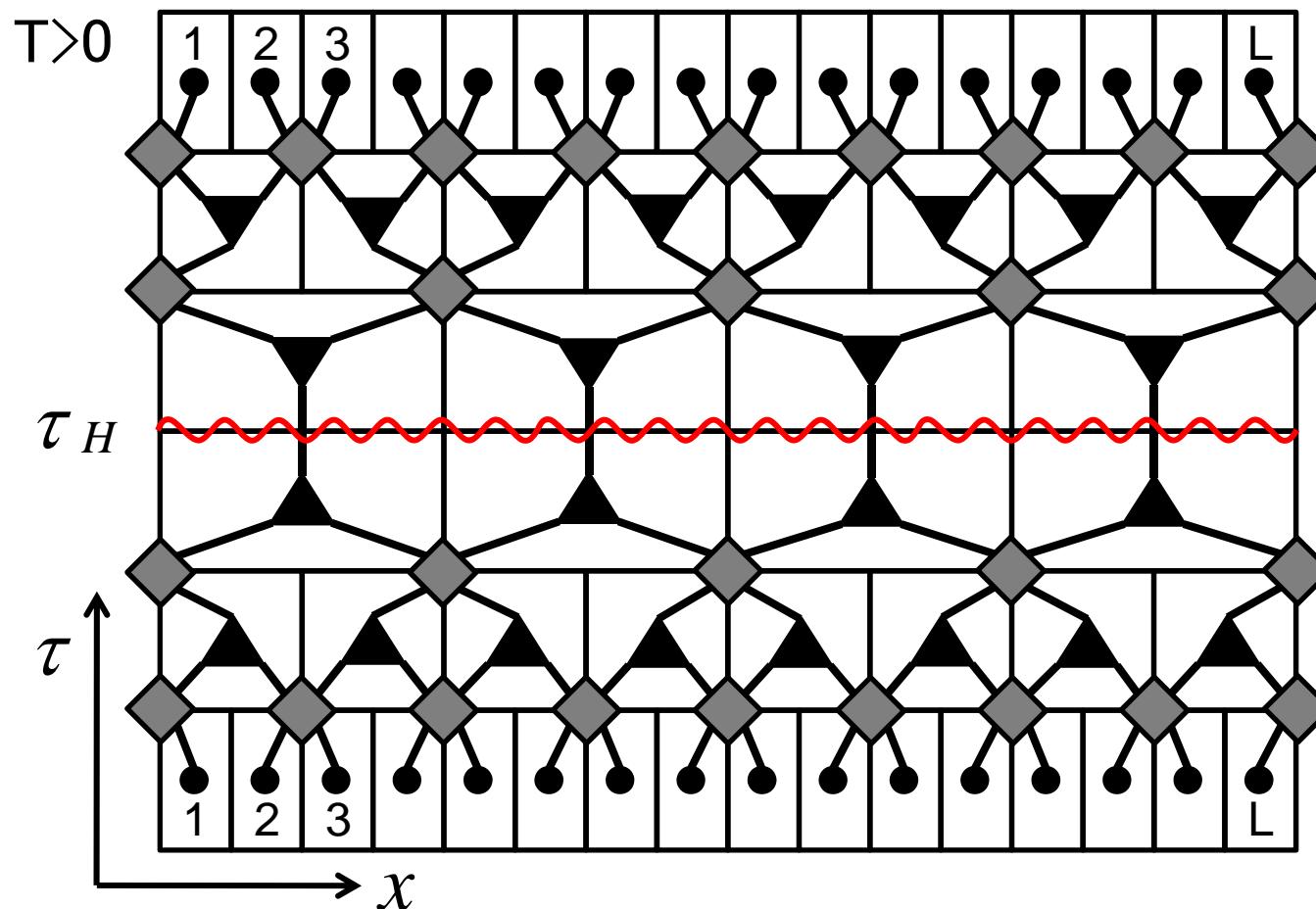


$$\alpha = j - j_H \quad \beta = 2e^{-j_H} \frac{x}{\varepsilon}$$

$$ds^2 = d\alpha^2 + (\cosh \alpha)^2 d\beta^2$$



# Finite-T MERA Network and AdS Black Hole

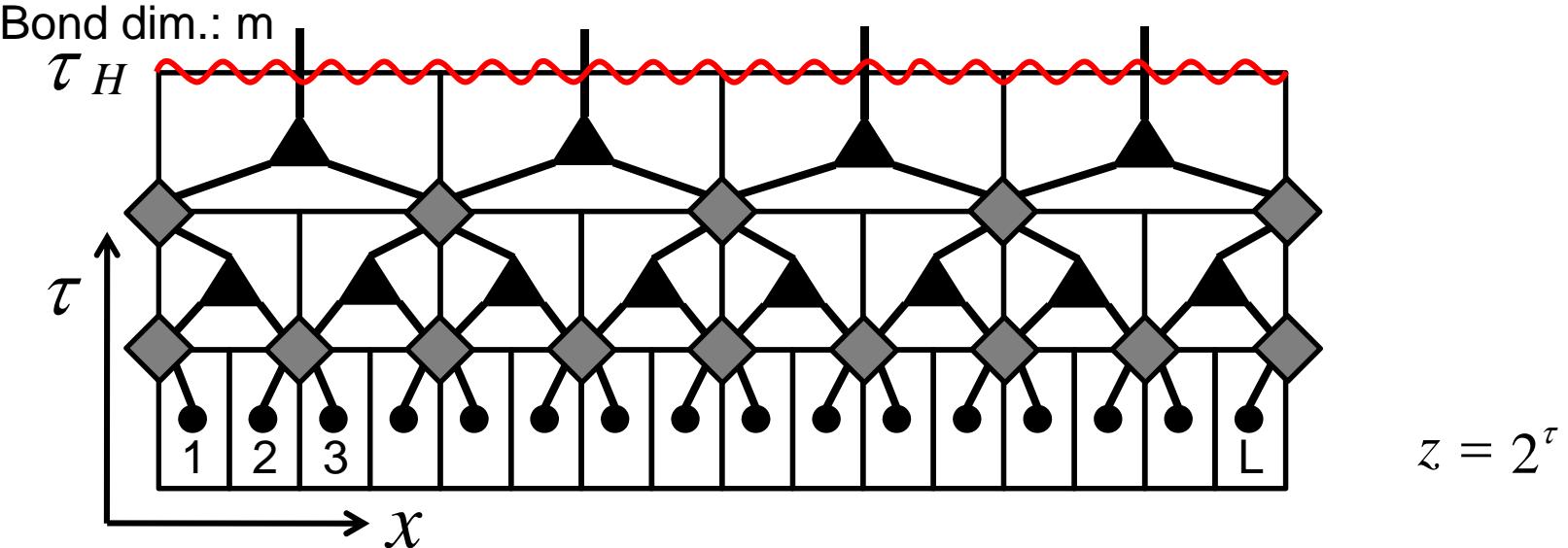


Vertical axis = energy scale, temperature scale

Wave function approach at finite-T  $\rightarrow$  thermofield dynamics  
 $\rightarrow$  Connection between original and tilde spaces

## Temperature of MERA Network

Truncation of upper MERA layers = AdS black hole



$$z = 2^\tau$$

$$\text{Area of interface: } \frac{L}{2^{\tau_H}} = A \quad \text{Total dim. at interface: } \chi = m^A$$

Beckenstein–Hawking entropy & Calabrese–Cardy formula:

$$S_{BH} = A \ln m = \frac{L}{z_H} \ln m$$

$$S_{CFT} = \frac{c}{3} \ln \left( \frac{\beta}{\pi \epsilon} \sinh \left( \frac{\pi L}{\beta} \right) \right)$$

$$k_B T = \left( \frac{3}{c \pi} \ln m \right) \frac{1}{z_H} \propto \frac{1}{z_H}$$

## Summary

### Information geometrical interpretation of MERA

- Bures metric of MERA  $\sim$  radial axis of AdS
- massive case  $\sim$  deformation of AdS

### Finite-T MERA network

- Truncation of the network at the IR region
- layer number of MERA  $\sim 1/T$
- consistency with field-theoretical treatment

# 資料集

## Target of this talk: Fisher metric and AdS/CFT

Information geometry → Fisher metric

$$g_{\mu\nu}(\theta) = \sum_n p_n(\theta) \frac{\partial \ln p_n(\theta)}{\partial \theta^\mu} \frac{\partial \ln p_n(\theta)}{\partial \theta^\nu}$$

$$\sum_n p_n(\theta) = 1$$

$\theta$  : internal parameters properly  
describing the system

- (a) The **universal** metric in information geometry  
(at least in a classical level)
- (b) **Natural source to produce quantum-classical correspondence**
- (c) **Similarity to Lagrangian of free scalar field**
- (d) We may construct examples in condensed matter physics.
- (e) Point (entropy) → Spacetime (metric+connection)

Questions:

What kind of correspondence comes from information geometry ?  
Is the new correspondence related to AdS/CFT ?

# 1D free scalar field

Hamiltonian

$$H = \frac{1}{2} \int dk \left[ \pi(k) \pi(-k) + \varepsilon_k^2 \phi(k) \phi(-k) \right] \quad \varepsilon_k = \sqrt{k^2 + m^2}$$

IR state

$$\left( \sqrt{M} \phi(x) + \frac{i}{\sqrt{M}} \pi(x) \right) |\Omega\rangle = 0 \quad M \approx O(\Lambda), \Lambda : \text{UV cut-off}$$

IR state  $\rightarrow$  scale invariant

$$e^{-iuL} \phi(x) e^{iuL} = e^{u/2} \phi(e^u x)$$

$$e^{-iuL} \phi(k) e^{iuL} = e^{-u/2} \phi(e^{-u} k)$$

$$e^{-iuL} \pi(x) e^{iuL} = e^{u/2} \pi(e^u x)$$

$$e^{-iuL} \pi(k) e^{iuL} = e^{-u/2} \pi(e^{-u} k)$$

## Assumption for functional form of $K(u)$

$$K(u) = \frac{1}{2} \int dk g(k, u) [\phi(k) \pi(-k) + \pi(k) \phi(-k)]$$

Effect of entangler  
on scale dimension

$$U(0, u)^{-1} \phi(k) U(0, u) = e^{-f(k, u)} e^{-u/2} \phi(e^{-u} k)$$

$$U(0, u)^{-1} \pi(k) U(0, u) = e^{f(k, u)} e^{-u/2} \pi(e^{-u} k)$$

$$K_I(u) = e^{iuL} K(u) e^{-iuL} = \frac{1}{2} \int dk g(k e^{-u}, u) [\phi(k) \pi(-k) + \pi(k) \phi(-k)]$$

$$\frac{\partial f(k, u)}{\partial u} = g(k e^{-u}, u)$$

$$E = \langle \Psi | H | \Psi \rangle = \langle \Omega | H(u_{IR}) | \Omega \rangle = \frac{1}{4} \int dk \left[ e^{2f(k, u_{IR})} M + \frac{\epsilon_k^2}{M} e^{-2f(k, u_{IR})} \right]$$

$$\frac{\delta E}{\delta \chi(u)} = \int_{|k| \leq \Lambda e^u} dk \left[ e^{2f(k, u_{IR})} M - \frac{k^2}{M} e^{-2f(k, u_{IR})} \right] \Rightarrow f(k, u_{IR}) = \frac{1}{2} \log \frac{\epsilon_k}{M}$$

$$f(k, u_{IR}) = \int_0^{u_{IR}} g(k e^{-s}, s) ds = \int_0^{-\log \Lambda / |k|} \chi(s) ds \Rightarrow \chi(s) = \frac{1}{2} \frac{e^{2u}}{e^{2u} + m^2 / \Lambda^2}$$